

**Detection of genuine multipartite  
entanglement and its applications in secure  
communication**

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# OUTLINE

## ○ INTRODUCTION

- Tasks/problems & possible applications.
- Non-locality test (without any inequality).

## ○ OUR WORKS:

- Detection of true multipartite entanglement.
- DI-Quantum Key Distribution with measurement inputs.
- Quantum Digital Signatures.
- DI Quantum Liar Detection & Byzantine Agreement.
- DI Quantum Random Number Generator, etc.

# Tasks/Problems

Let  $H = H_1 \otimes H_2$ ;  $n \times m$  dim. Hilber space.

$$\begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \otimes \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} a_1 b_1 \\ a_1 b_2 \\ a_2 b_1 \\ a_2 b_2 \end{pmatrix}$$

Hilbert Space:  $H_1$

$H_2$

$H = H_1 \otimes H_2$

Orthonormal basis:  $\{|\eta_i\rangle_1\}_{i=1}^n$   $\{|\chi_j\rangle_2\}_{j=1}^m$   $\{|\eta_i\rangle_1 \otimes |\chi_j\rangle_2\}_{i=1,2,\dots,n}^{j=1,2,\dots,m}$

Any  $|\psi\rangle_{12} \in H$  can be expressed as 
$$\sum_{i=1,2,\dots,n} \sum_{j=1,2,\dots,m} \alpha_{ij} |\eta_i\rangle_1 |\chi_j\rangle_2$$

Product state:  $|\psi\rangle_{12} = |\eta\rangle_1 \otimes |\chi\rangle_2$

Entangled state:  $|\psi\rangle_{12} \neq |\eta\rangle_1 \otimes |\chi\rangle_2$  [A key feature, exists in quantum correlations.]

## Tasks/Problems cont...

Consider a system  $H = H_1 \otimes H_2 \otimes \dots \otimes H_n$ ;  $\text{Dim.}(H) = d_1 \cdot d_2 \dots d_n$

$|\Psi\rangle \in H$  is called:

- i) **fully product** if,  $|\Psi\rangle = |\eta\rangle_1 |\chi\rangle_2 \dots |e\rangle_n$
- ii) **bi-separable/product** if,  $|\Psi\rangle = |\phi\rangle_K |\xi\rangle_{\bar{K}}$ ;  $K \subset \{1, 2, \dots, n\}$
- iii) **genuine entangled** if,  $|\Psi\rangle \neq |\phi\rangle_K |\xi\rangle_{\bar{K}}$

E. g.:  $\alpha |0\rangle_1 |0\rangle_2 \dots |0\rangle_n + \beta |1\rangle_1 |1\rangle_2 \dots |1\rangle_n$  is genuinely entangled.

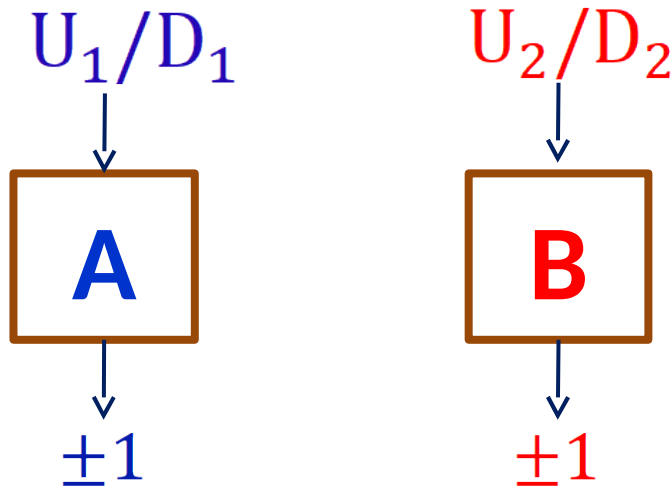
## Problems:

- General witness for **bi-separable** & **genuine entangled** states.
- Non-classicality & monogamous characteristics of  $|\Psi\rangle \in H$

# Applications

- Provides secure quantum protocols for various cryptographic & communication tasks. E.g.,
  - Key distribution
  - Digital signatures
  - Secret sharing
  - Byzantine agreement
  - Random number generator
  - Oblivious transfer
  - Dining cryptographers
  - Anonymous veto etc.
- Quantum algorithms & computation.
- Quantum simulation & metrology, etc.

# Hardy's Paradox [L. Hardy PRL 1992]



$$P(+, + | U_1, U_2) = q > 0$$

$$P(+, + | U_1, D_2) = 0$$

$$P(+, + | D_1, U_2) = 0$$

$$P(-, - | D_1, D_2) = 0$$

$P(a,b|X,Y)$  is the joint probability of getting the outcome  $(a,b)$  for the given input  $(X,Y)$ .

This set of conditions cannot be satisfied by any **Local-Realistic (LR) Theory** (Classical Theory).

# HARDY'S PARADOX & QM

$P(a, b|X, Y) = |\langle \psi | (|X = a\rangle |Y = b\rangle)|^2$  for the quantum state  $|\psi\rangle$ .  
 $|X = a\rangle$  is the eigenstate corresponding to the eigenvalue  $\mathbf{a}$ .

$$\begin{array}{ll} P(+, +|U_1, U_2) = q > 0 & |\phi_4\rangle = |U_1 = +1\rangle |U_2 = +1\rangle \\ P(+, +|U_1, D_2) = 0 & |\phi_3\rangle = |U_1 = +1\rangle |D_2 = +1\rangle \\ P(+, +|D_1, U_2) = 0 & |\phi_2\rangle = |D_1 = +1\rangle |U_2 = +1\rangle \\ P(-, -|D_1, D_2) = 0 & |\phi_1\rangle = |D_1 = -1\rangle |D_2 = -1\rangle \end{array}$$

Let  $|D_j = +1\rangle = \mathbf{a}_j |U_j = +1\rangle + \mathbf{b}_j |U_j = -1\rangle$ ,  $j = 1, 2$ ; with  $|\mathbf{a}_j|^2 + |\mathbf{b}_j|^2 = 1$  &  $0 < |\mathbf{a}_j| < 1$ .

$|\phi_1\rangle, |\phi_2\rangle, |\phi_3\rangle, |\phi_4\rangle$  are linearly independent.

If  $\mathbf{S} = \{|\phi_1\rangle, |\phi_2\rangle, |\phi_3\rangle\}$ , then  $\dim(\mathbf{S}) = 3$ .

Hardy state  $|\Psi\rangle \perp \mathbf{S}$  &  $\dim(H_A \otimes H_B) = 2 \times 2$ .

$\therefore |\Psi\rangle$  is unique [Ref. G Kar, PLA 97].

# HARDY STATE

$$|\phi_1\rangle = |D_1 = -1\rangle|D_2 = -1\rangle \quad |\phi_3\rangle = |U_1 = +1\rangle|D_2 = +1\rangle$$

$$|\phi_2\rangle = |D_1 = +1\rangle|U_2 = +1\rangle \quad |\phi_4\rangle = |U_1 = +1\rangle|U_2 = +1\rangle$$

By Gram-Schmidt orthogonalization procedure

$$|\phi'_1\rangle = |\phi_1\rangle;$$

$$|\phi'_i\rangle = \frac{|\phi_i\rangle - \sum_{j=1}^{i-1} \langle \phi'_j | \phi_i \rangle |\phi'_j\rangle}{\sqrt{1 - \sum_{j=1}^{i-1} |\langle \phi'_j | \phi_i \rangle|^2}}; i = 2, 3, 4$$

$$\therefore \text{Hardy state } |\Psi\rangle = |\phi'_4\rangle$$



## PROBABILITY OF SUCCESS

Probability of Success  $q = |\langle \Psi | \phi_4 \rangle|^2 = \frac{|a_1 a_2|^2 |b_1 b_2|^2}{1 - |a_1 a_2|^2}$ .

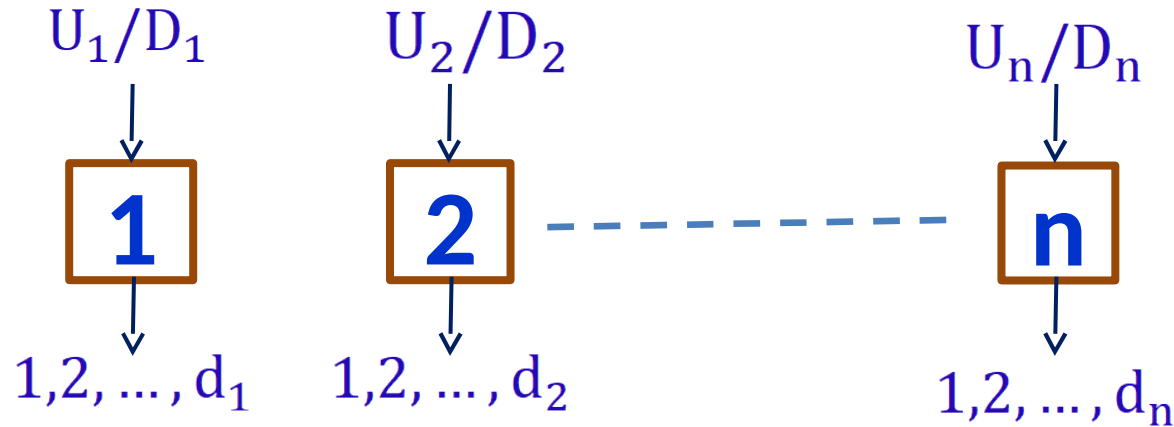
Its maximum is  $\frac{5\sqrt{5} - 11}{2} = 0.09$ , where  $|a_1| = |a_2| = \sqrt{\frac{\sqrt{5}-1}{2}}$ .  
 $\therefore U_1 \equiv U_2$  &  $D_1 \equiv D_2$ .

For  $q=0.09$  (max.), associated Hardy state  $|\Psi\rangle$  is **device-independent** [Rabelo et al. PRL 2012].

For  $q_{\max} = \frac{5\sqrt{5} - 11}{2}$  the state is equivalent to  $|\psi_{\max}\rangle_{12} \otimes |\eta_{1,2}\rangle$ .

# NON-LOCALITY TEST FOR GENUINE ENTANGLEMENT.

Consider a system  $H = H_1 \otimes H_2 \otimes \dots \otimes H_n$ ;  $\text{Dim.}(H) = d_1 \cdot d_2 \dots d_n$



$$P(11 \dots 1 | D_1 D_2 \dots D_n) = 0$$

$$P(a_r 1 | D_r U_{r+1}) = 0 \quad a_r \neq 1$$

$$P(11 \dots 1 | U_1 U_2 \dots U_n) = q > 0$$

Again cannot be satisfied by any LR theory

Also, states of the form  $|\Phi\rangle_K \otimes |\xi\rangle_{\bar{K}}$ ;  $K \subset \{1, 2, \dots, n\}$  cannot.

That is, **only genuine entangled states** can satisfy.

[Rahaman et al., Phys. Rev. A 2014]

**Proof:** If  $\rho$  is not a genuine entangled state then all joint probabilities can be expressed as,

$$P_{\rho}(x_1 x_2 \dots x_n | X_1 X_2 \dots X_n) = \sum_m p_m Q(x_1 x_2 \dots x_m | X_1 X_2 \dots X_m) R(x_{m+1} \dots x_n | X_{m+1} \dots X_n)$$

$P_{\rho}(11 \dots 1 | U_1 U_2 \dots U_n) > 0$  implies, for some  $m \in \{m\}$ ,

$$Q(11 \dots 1 | U_1 U_2 \dots U_m) R(11 \dots 1 | U_{m+1} U_{m+2} \dots U_n) = q' > 0,$$

**Middle conditions:**  $Q(a_m \neq 1 | D_m) R(1 | U_{m+1}) = 0$  &  $R(a_n \neq 1 | D_n) Q(1 | U_1) = 0$ .

$\therefore R(1 | U_{m+1}) \neq 0$ , so  $Q(a_m \neq 1 | D_m) = 0$ . Thus,  $Q(1 | D_m) = 1$ .

$\therefore Q(x_1 x_2 \dots x_m | X_1 X_2 \dots X_m) = Q'(x_1 x_2 \dots x_{m-1} | X_1 X_2 \dots X_{m-1}) Q(x_m | X_m)$ .

Also middle conditions give us,  $Q'(a_{m-1} \neq 1 | D_{m-1}) Q_m(1 | U_m) = 0$ .

Again,  $Q_m(1 | U_m) > 0$ , so  $Q'(a_{m-1} \neq 1 | D_{m-1})$  i. e.,  $Q'(1 | D_{m-1}) = 1$ .

Proceeding like this we show that  $Q$  fully factorizes.

Following the same steps, we can show that R is fully factorized with the help of  $\mathbf{R}(\mathbf{a}_n \neq \mathbf{1} | \mathbf{D}_n) \mathbf{Q}(\mathbf{1} | \mathbf{U}_1) = \mathbf{0}$

Thus the state representing the considered term is fully factorizable.

Such states admit local hidden variable models, and as such **cannot** satisfy the mentioned set of joint probability conditions for  $\mathbf{q}' > \mathbf{0}$ .

# For qubits system: $\text{Dim. (H)} = 2 \times 2 \times \dots \times 2$

$$P(11 \dots 1 | D_1 D_2 \dots D_n) = 0; |\phi_-\rangle = |D_1 = 1\rangle |D_2 = 2\rangle \dots |D_n = 1\rangle$$

$$P(a_r 1 | D_r U_{r+1}) = 0; |\phi_{k_r}\rangle = |\dots\rangle \dots |D_r = 2\rangle |U_{r+1} = 1\rangle \dots |\dots\rangle$$

$$P(11 \dots 1 | U_1 U_2 \dots U_n) = q; |\phi_+\rangle = |U_1 = 1\rangle |U_2 = 1\rangle \dots |U_n = 1\rangle$$

**Define a new basis:**  $|00 \dots 0 \dots 0\rangle = |\phi_+\rangle$ ,

$$|00 \dots 01_l 0 \dots 0\rangle = \frac{1}{\beta_l} [|\phi_k(0, \dots, 0, +_l, 0, \dots, 0)\rangle - \alpha_l |\phi_+\rangle], \forall l,$$

$$|0 \dots 01_l 0 \dots 01_m 0 \dots 0\rangle = \frac{1}{\beta_l \beta_m} [|\phi_k(0, \dots, 0, +_l, 0, \dots, 0, +_m, 0, \dots, 0)\rangle - \alpha_l \alpha_m |\phi_+\rangle - \beta_l \alpha_m |00 \dots 01_l 0 \dots 0\rangle - \alpha_l \beta_m |00 \dots 01_m 0 \dots 0\rangle], \forall l \neq m,$$

$$|0 \dots 01_l 0 \dots 01_m 0 \dots 01_k 0 \dots 0\rangle = \frac{1}{\beta_l \beta_m \beta_k} [|\phi_k(0, \dots, 0, +_l, 0, \dots, 0, +_m, 0, \dots, 0, +_k, 0, \dots, 0)\rangle - \alpha_l \alpha_m \alpha_k |\phi_+\rangle - \alpha_l \alpha_m \beta_k |00 \dots 01_k 0 \dots 0\rangle - \alpha_l \beta_m \alpha_k |00 \dots 01_m 0 \dots 0\rangle - \beta_l \alpha_m \alpha_k |00 \dots 01_l 0 \dots 0\rangle - \alpha_l \beta_m \beta_k |00 \dots 01_m 0 \dots 01_k 0 \dots 0\rangle - \beta_l \alpha_m \beta_k |00 \dots 01_l 0 \dots 01_k 0 \dots 0\rangle - \beta_l \beta_m \alpha_k |00 \dots 01_l 0 \dots 01_m 0 \dots 0\rangle],$$

$\forall l \neq m \neq k \neq l, \dots$

$$|11 \dots 1 \dots 1\rangle = \frac{(-1)^N}{\prod_{i=1}^N \alpha_i^*} \left[ |\phi_0\rangle - \left\{ \left( \prod_{j=1}^N \beta_j^* \right) |\phi_+\rangle + (-1)^1 \sum_{i=1}^N \alpha_i^* \left( \prod_{j=1, j \neq i}^N \beta_j^* \right) |00 \dots 01_i 0 \dots 0\rangle + (-1)^2 \sum_{i, l=1, i \neq l}^N \alpha_i^* \alpha_l^* \left( \prod_{j=1, j \neq i, l}^N \beta_j^* \right) |00 \dots 01_i 0 \dots 01_l 0 \dots 0\rangle + \dots + (-1)^{N-1} \sum_{j=1}^N \beta_j^* \left( \prod_{i=1, i \neq j}^N \alpha_i^* \right) |11 \dots 10_j 1 \dots 1\rangle \right\} \right],$$

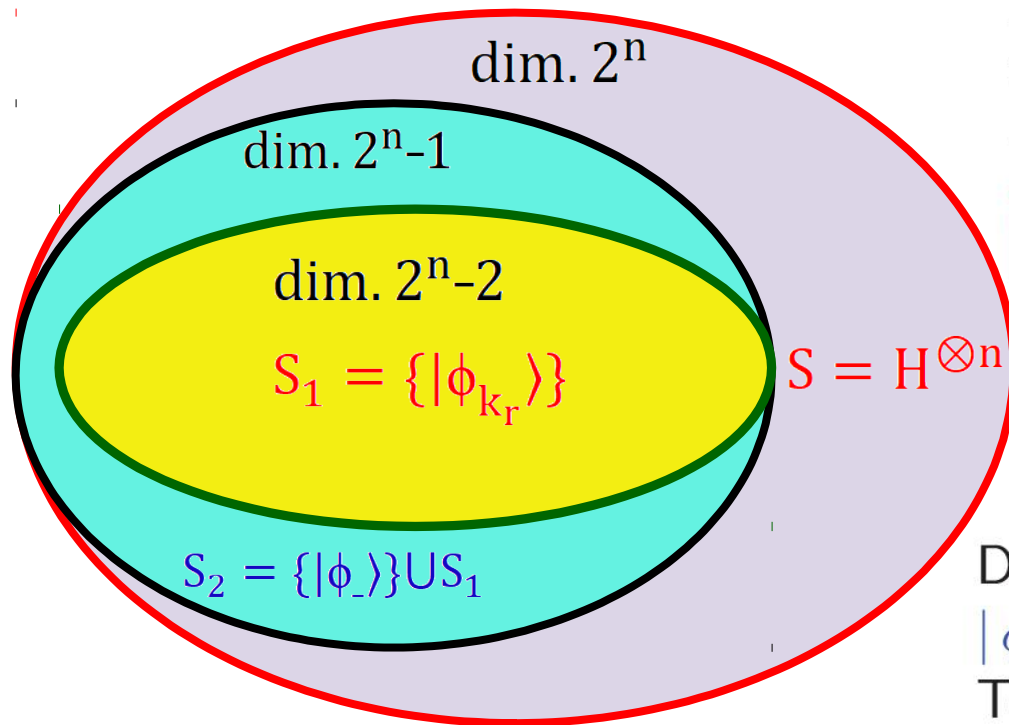
where  $|+\rangle_j = \alpha_j |0\rangle_j + \beta_j |1\rangle_j$ , and  $|-\rangle_j = \beta_j^* |0\rangle_j - \alpha_j^* |1\rangle_j$

**For qubits system:**  $\text{Dim. (H)} = 2 \times 2 \times \dots \times 2$

$$P(11 \dots 1 | D_1 D_2 \dots D_n) = 0; |\phi_{-}\rangle = |D_1 = 1\rangle |D_2 = 2\rangle \dots |D_n = 1\rangle$$

$$P(a_r 1 | D_r U_{r+1}) = 0; |\phi_{k_r}\rangle = |\dots\rangle \dots |D_r = 2\rangle |U_{r+1} = 1\rangle \dots |\dots\rangle$$

$$P(11 \dots 1 | U_1 U_2 \dots U_n) = q; |\phi_{+}\rangle = |U_1 = 1\rangle |U_2 = 1\rangle \dots |U_n = 1\rangle$$



Let  $\mathcal{S}_1 = \{|\phi_{k_r}\rangle\}$

$|\phi_{-}\rangle \perp \{\mathcal{S}_1\}$  and  $|\phi_{-}\rangle \not\perp |\phi_{+}\rangle$ .

$\therefore |\phi_{+}\rangle \notin \{\mathcal{S}_1\}$

$|\phi_{+}\rangle \neq |\phi_{-}\rangle, \therefore \text{dim.}(\mathcal{S}_1) = 2^n - 2$

Define:  $\mathcal{S}_2 = \{|\phi_{-}\rangle\} \cup \mathcal{S}_1$ .

$|\phi_{+}\rangle \notin \mathcal{S}_2, \therefore \text{dim.}(\mathcal{S}_2) = 2^n - 1$ .

To satisfy Hardy conditions:  $|\psi\rangle \perp \mathcal{S}_2$ .

Hardy state  $|\Psi\rangle$  is unique & genuinely entangled

[Rahaman et al., Phys. Rev. A 2014]

## Relaxed Hardy type test for genuine multiparty entangled states

$$P(11 \dots 1 | D_1 D_2 \dots D_n) = 0$$

$$P(1 \dots \neg 1 \dots 1 | U_1 \dots D_r \dots U_n) = 0$$

$$P(1 \dots 1 \dots 1 \dots 1 | U_1 \dots D_i \dots D_j \dots U_n) = q$$

Only **genuine** multiparty entangled states can satisfy

[S. S. Bhattacharya, A. Roy, A. Mukherjee & R. Rahaman, *Phys. Rev. A*, 92, 012111 (2015)]

# **KEY DISTRIBUTION PROTOCOL**



# Private key cryptography

$M=10110010$



$E=11010001$

$C=01100011$



EVE

M



Communication Channel

$C=01100011$



$D=11010001$

$M=10110010$

## Difficulties in Private Key

- The key bits cannot be reused for any future protocol.
- Key bits must be delivered in advance, guarded assiduously until used.

## Public key cryptosystems-

- W. Diffie and M. Hellman (1976).
- R. Rivest, A. Shamir and L. Adleman (1978) [RSA].

### Possible attacks:

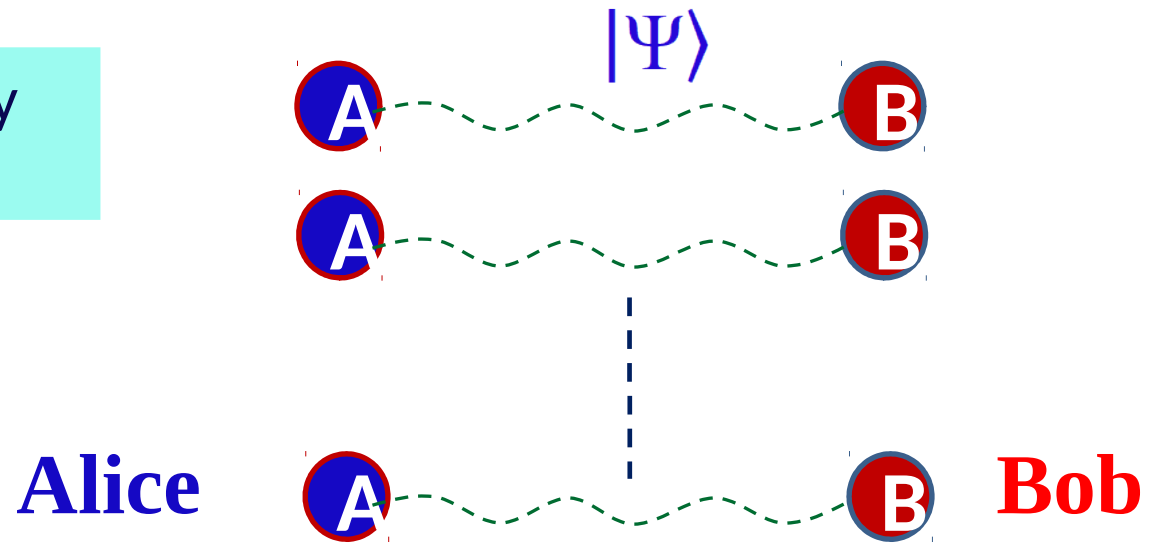
- If you able to factor  $n$ .
  - Security based on computational hardness
  - Can be broken by quantum computers!

### Possible Solution:-

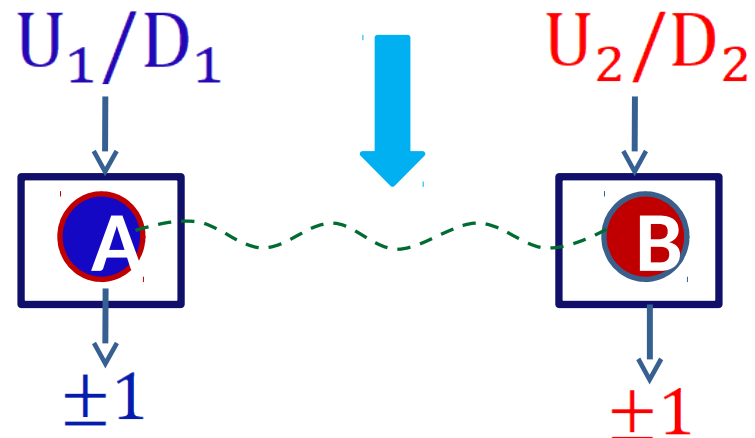
- Quantum cryptography
  - Security based on Laws of Physics

# DI-QKD PROTOCOL [Rahaman et al. PRA 92, 062304 (2015)]

Alice & Bob share many copies of Hardy State.



Measures own qubits in the basis chosen randomly from  $\{U, D\}$ .



## DI-QKD [*Rahaman et al. PRA 2015*] **CONT...**

Since,  $\langle \Psi | (|D_1 = +1\rangle |D_2 = +1\rangle) \neq 0$

For  $|\Psi\rangle$ ;  $P(+, + | D_1, D_2) > 0$ .

Thus,  $P(+, + | U_1, U_2) > 0$   
&  $P(+, + | D_1, D_2) > 0$ .

Remember Hardy's Paradox

$$P(+, + | U_1, U_2) = q > 0$$

$$P(+, + | U_1, D_2) = 0$$

$$P(+, + | D_1, U_2) = 0$$

$$P(-, - | D_1, D_2) = 0$$

\ For (+,+) outcome local inputs are correlated  
[(U,U) or (D,D)].

**Key assign: U  $\rightarrow$  0 & D  $\rightarrow$  1.**

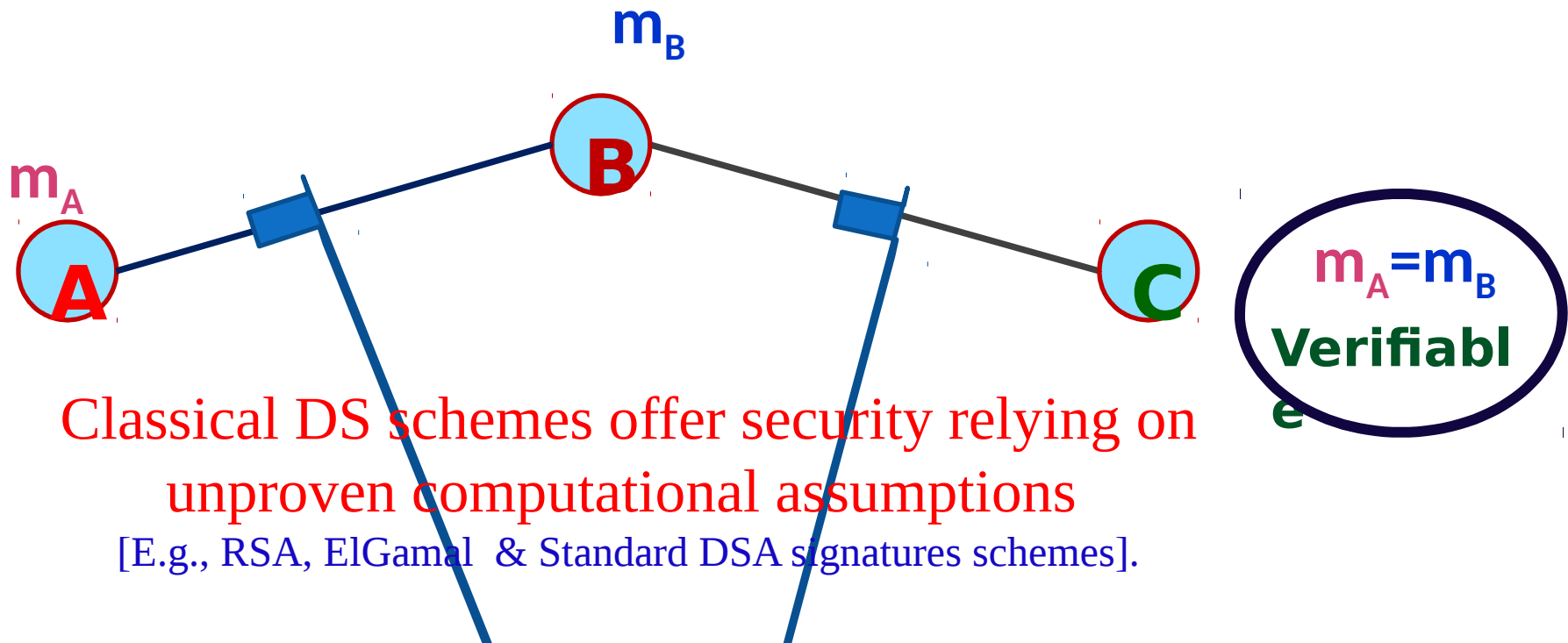
For  $q=0.09$  (max.), associated Hardy state  $|\Psi\rangle$  is device-independent [*Rabelo et al. PRL 2012*].

**⊗ Our QKD also DI in this case.**

# **DIGITAL SIGNATURES (DS) PROBLEM**

## Digital signature (DS)

- DS allows to send authentic message(s) from one sender to multiple recipients.
- In a DS the known sender cannot deny having sent the message.
- Also, the message was not altered in transit.



Classical DS schemes offer security relying on unproven computational assumptions

[E.g., RSA, ElGamal & Standard DSA signatures schemes].

**QUANTUM DIGITAL  
SIGNATURES(QDS) PROTOCOL**

## Existing QDS schemes

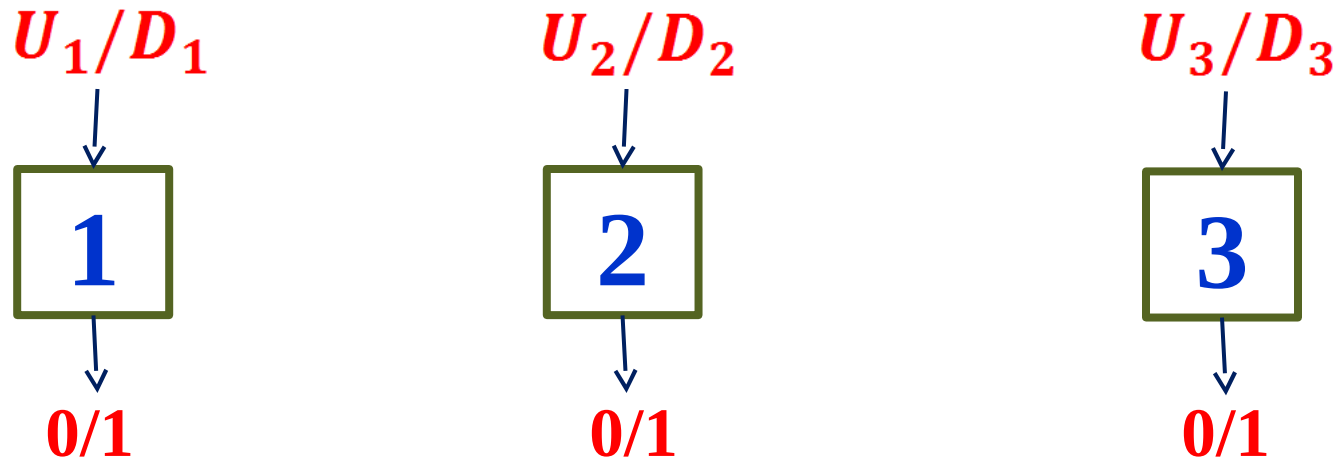
- In 2001, Gottesman & Chuang, arXiv:quant-ph/0105032.
- Experimental demonstration with coherent states:
  - (i) E. Andersson et. al., PRA 2006.
  - (ii) P. J. Clarke et. al., Nature Com. 2012.
- QDS schemes without Quantum Memory:
  - (i) V. Dunjko et. al., PRL 2014.
  - (ii) R. J. Collins et. al., PRL 2014.



# THREE QUBITS HARDY PARADOX

Consider a system  $H=H_1 \otimes H_2 \otimes H_3$

$$\text{Dim.}(H)=2 \times 2 \times 2=8$$



$$P(000|U_1U_2U_3) = q > 0$$

$$P(00|U_iD_j) = 0 \quad i \neq j$$

$$P(111|D_1D_2D_3) = 0$$

**Again cannot be satisfied by any LR theory**

For 3-qubits system:  $\text{Dim.}(\mathbf{H})=2 \times 2 \times 2=8.$

**Let us assign:**

$$P(000|U_1U_2U_3) = q \quad |\phi_+\rangle = |0\rangle|0\rangle|0\rangle$$

$$P(00|U_iD_j) = 0 \quad |\phi_{k_r}\rangle = |\dots\rangle|0_i\rangle|0'_j\rangle$$

$$P(111|D_1D_2D_3) = 0 \quad |\phi_-\rangle = |1'\rangle|1'\rangle|1'\rangle$$

Let  $S_1 = \{|\phi_{k_r}\rangle\} \cup \{|\phi_-\rangle\}.$

Then  $\text{dim.}(S_1) = 2^3 - 1$

**Hardy state**  $|\Psi\rangle \perp S_1$

**Hardy state  $|\Psi\rangle$  is unique & genuinely entangled**

[Rahaman et al., Phys. Rev. A 2014].

## 3-QUBIT HARDY STATE

Probability of success  $q = 0.0181938$ .

In this case also,  $U_1 \cong U_2 \cong U_3 \cong U$  &  
 $D_1 \cong D_2 \cong D_3 \cong D$

$$P(000|U_1U_2U_3) = q > 0$$

$$P(00|U_iD_j) = 0 \quad i \neq j$$

$$P(111|D_1D_2D_3) = 0$$

**Hardy State:**  $|\Psi\rangle = c_0|000\rangle + c_1P[|001\rangle] + c_2P[|011\rangle] + c_3|111\rangle$

$$c_0 = \frac{|\alpha|^3|\beta|^3}{\sqrt{1-|\alpha|^6}}, c_1 = \frac{-\beta|\alpha|^4|\beta|}{\sqrt{1-|\alpha|^6}}, c_2 = \frac{\beta^2|\alpha|^5}{|\beta|\sqrt{1-|\alpha|^6}}, c_3 = \frac{\beta^3\sqrt{1-|\alpha|^6}}{|\beta|^3}$$

$$|0'\rangle = \alpha|0\rangle + \beta|1\rangle \text{ \& \ } |1'\rangle = \beta^*|0\rangle - \alpha^*|1\rangle$$

Device Independent Hardy Test (3-qubit):-

For  $q = 0.0181938$ , state is equivalent to  $|\Psi\rangle_{123} = |\psi_{\max}\rangle_{123} \otimes |\eta_{1'2'3'}\rangle$ .

# Quantum digital signatures protocol

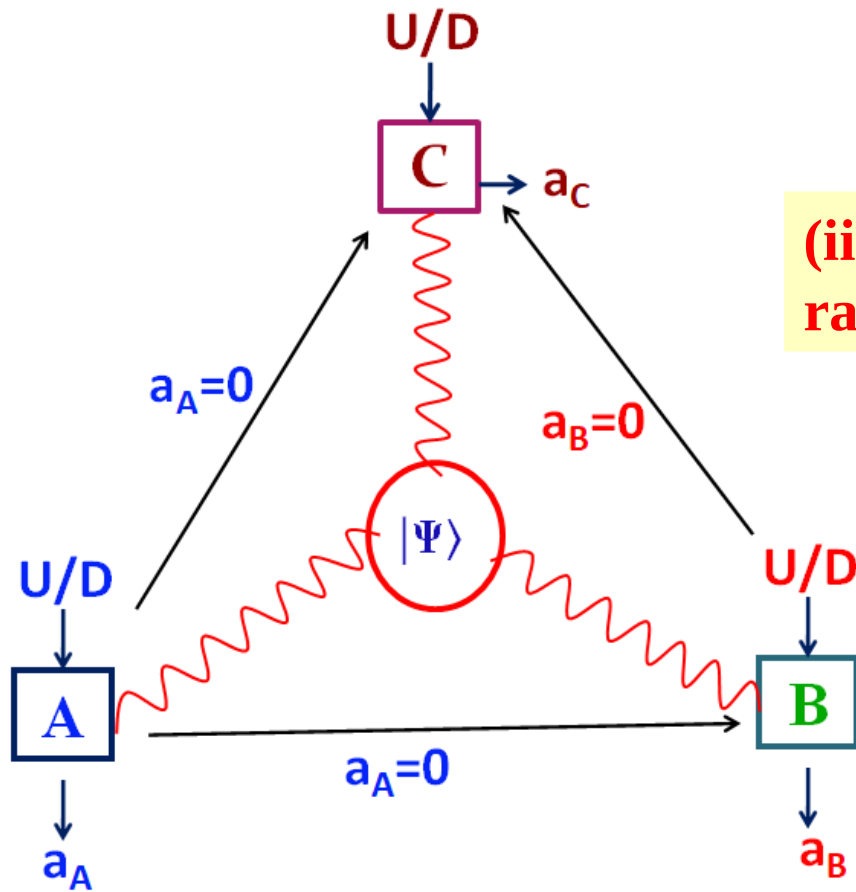
- **S1. Distribution of resources:** 'A' prepares and shares a large number of 3-qubits Hardy state  $|\Psi\rangle$  with B and C.

$$|\Psi\rangle = c_0|000\rangle + c_1P[|001\rangle] + c_2P[|011\rangle] + c_3|111\rangle$$

$$c_0 = \frac{|\alpha|^3|\beta|^3}{\sqrt{1-|\alpha|^6}}, c_1 = \frac{-\beta|\alpha|^4|\beta|}{\sqrt{1-|\alpha|^6}}, c_2 = \frac{\beta^2|\alpha|^5}{|\beta|\sqrt{1-|\alpha|^6}}, c_3 = \frac{\beta^3\sqrt{1-|\alpha|^6}}{|\beta|^3}$$

$$|0'\rangle = \alpha|0\rangle + \beta|1\rangle \text{ \& \ } |1'\rangle = \beta^*|0\rangle - \alpha^*|1\rangle$$

**S2. Actions:** (i) **A** measures all his qubits in the message basis he want to convey. [ $U$  for  $m=0$  and  $D$  for  $m=1$ ].



$$P(000|UUU)=q$$

$$P(00|UD)=0$$

$$P(111|DDD)=0$$

(ii) **B(C)** measures all his qubits in random basis  $U/D$ .

(iii) **A** sends the list of runs to **B** & **C** when  $a_A=0$ .

(iv) **B** sends the list of runs to **C** when  $a_B=0$ .

They discards the runs when **A** gets outcome 1.

$$P(000|UUU)=q$$

$$P(00|UD)=0$$

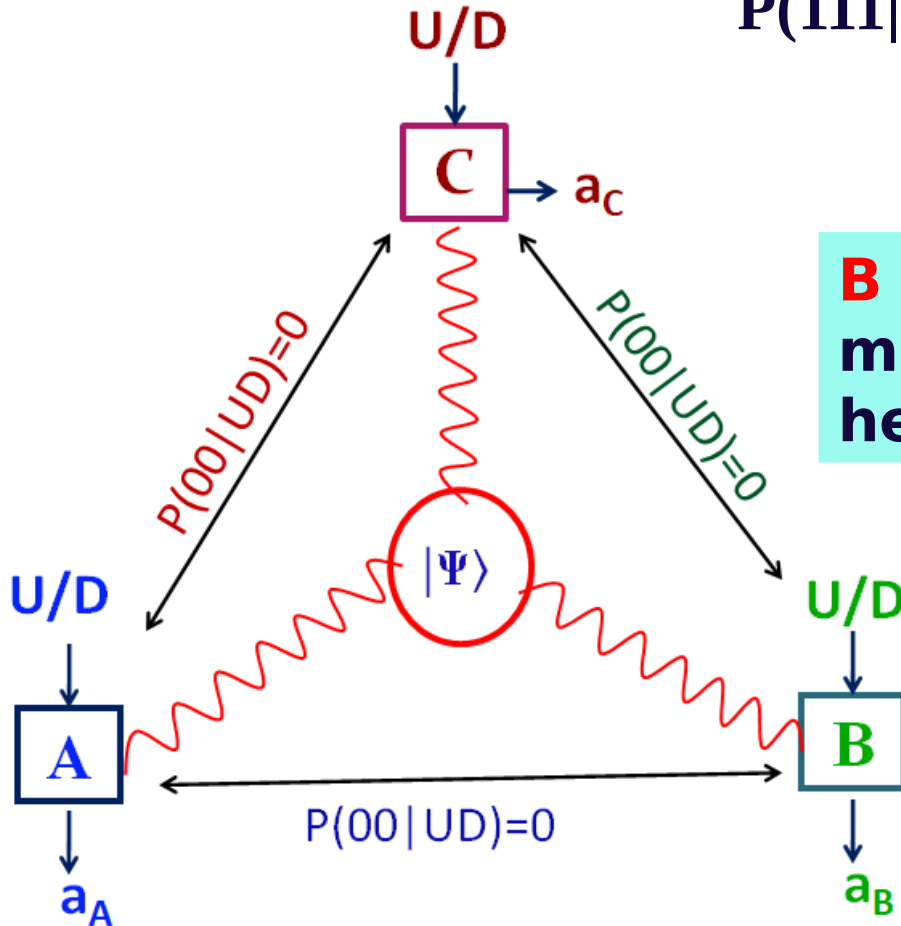
$$P(111|DDD)=0$$

$$P(00|DD)>0$$

$$P(00|UU)>0$$

$$P(00|DU)=0$$

$$P(00|UD)=0$$



**B** can easily figure out the message basis  $[m]$  of **A** with help of **Hardy's conditions**.

**C** can also easily verify **B's** claim with help of **Hardy's conditions**.

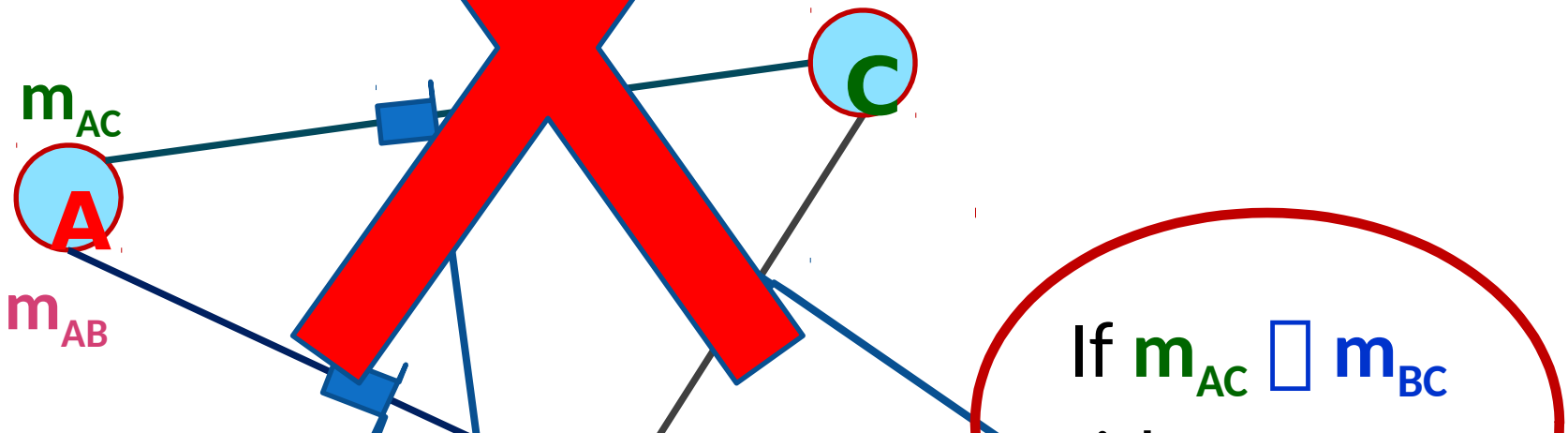
# **LIAR DETECTION (LD) PROBLEM**

## Liar Detection

'C' receives a message from A in two different paths:

(i) Directly from A [message  $m_{AC}$ ]

(ii) Via B [message  $m_{BC}$ ]



Quantum solution exists: A. Cabello, PRL  
2002, PRA 2003.



## Quantum solution for Liar Detection (QLD) Problem

- S1. Distribution of resources:**

(i) 'C' shares a large number ( $\approx 6N$ ) of maximally entangled states

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}}[|uu\rangle + |u^\perp u^\perp\rangle]$$

(ii) Conversion from  $|\Phi^+\rangle$  to  $|\psi^H\rangle$  between 'C and A' and 'C and B'.

$$(\mathbb{U}^{c1} \otimes \mathbb{I}^2)|u\rangle_c |\Phi^+\rangle_{12} = \frac{1}{\sqrt{2}}|u\rangle_c |\psi^H\rangle_{12} + \frac{1}{\sqrt{2}}|u^\perp\rangle_c |\psi'\rangle_{12},$$

$$|\psi^H\rangle = x_{00}|u\rangle_1|u\rangle_2 + x_{01}(|u\rangle_1|u^\perp\rangle_2 + |u^\perp\rangle_1|u\rangle_2) + x_{11}|u^\perp\rangle_1|u^\perp\rangle_2, \quad |\psi'\rangle = x_{01}^*|uu\rangle - x_{00}^*|uu^\perp\rangle + x_{11}^*|u^\perp u\rangle - x_{01}^*|u^\perp u^\perp\rangle.$$

$$\begin{aligned} \mathbb{U}|uu\rangle &= x_{00}|uu\rangle + x_{01}|uu^\perp\rangle + x_{01}^*|u^\perp u\rangle + x_{11}^*|u^\perp u^\perp\rangle, \\ \mathbb{U}|uu^\perp\rangle &= x_{01}|uu\rangle + x_{11}|uu^\perp\rangle - x_{00}^*|u^\perp u\rangle - x_{01}^*|u^\perp u^\perp\rangle \end{aligned}$$

## Quantum solution for Liar Detection (QLD)

- **S1. Distribution of resources:**

(iii) Conversion from  $|\Phi^+\rangle$  to  $|\psi^H\rangle$  between 'A and B'.

C applies a two outcome joint measurement  $\{M, I-M\}$  on her two qubits.

$$M = |\psi^{H*}\rangle\langle\psi^{H*}|$$

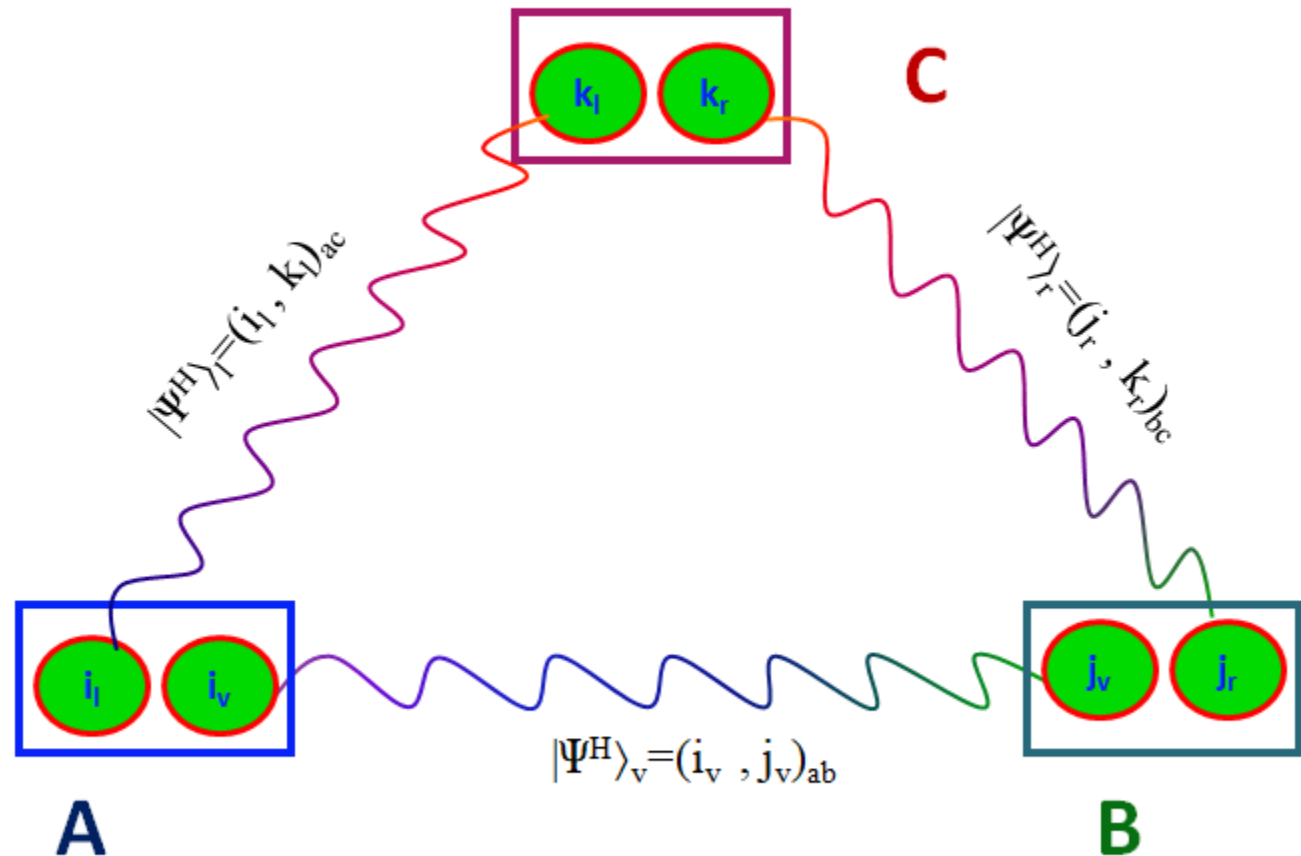
$$|\psi^{H*}\rangle = x_{00}^*|uu\rangle + x_{01}^* (|uu^\perp\rangle + |u^\perp u\rangle) + x_{11}^*|u^\perp u^\perp\rangle.$$

$$(M^{13} \otimes \mathbb{I}^{24})|\Phi^+\rangle_{12}|\Phi^+\rangle_{34} = \frac{1}{2}|\psi^{H*}\rangle_{13} \otimes |\psi^H\rangle_{24}.$$

(iv) 'C' prepares a list  $L_B = \{(j_v, k_v)_{AB}\}_{v=1}^t$  for 'B' and sends to him.

Neither B nor C reveal the inform of  $L_B$  to A.

After a successful distribution of qubits



each party shares  $\approx N$  copies of Hardy state with others.

## QLD Protocol.....

### **S2. Actions on qubits distributed by C :**

(i) **Action by 'A'**: **A** measures all his qubits in the message basis he want to convey.

*[Measurement  $U$  for the message  $m=0$  and  $D$  for  $m=1$ ].*

(ii) **Action by 'B'**: (a) **B** measures qubits of  $L_B$  in random bases  $U/D$ .

**B** can easily figure out the message basis [ $m_{AB}$ ] of **A** by comparing his measurement data and the results of **A** with **Hardy's conditions**.

(b) **B** measures rest of his qubits in the message basis and sends the results to **C**.

## QLD Protocol .....

### [S2. Actions on qubits distributed by C]

(iii) **Action by 'C'**: C measures all his qubits in the random bases **U/D**.

(a) 'C' can easily figure out the message basis [ $\mathbf{m}_{AC}$ ] of **A** by comparing his measurement data and the results of **A** with **Hardy's conditions**.

(b) Similarly, 'C' can find out the message basis [ $\mathbf{m}_{BC}$ ] of **B** with comparing the data and Hardy's conditions.

$$P(11 | D_1 D_2) = 0$$

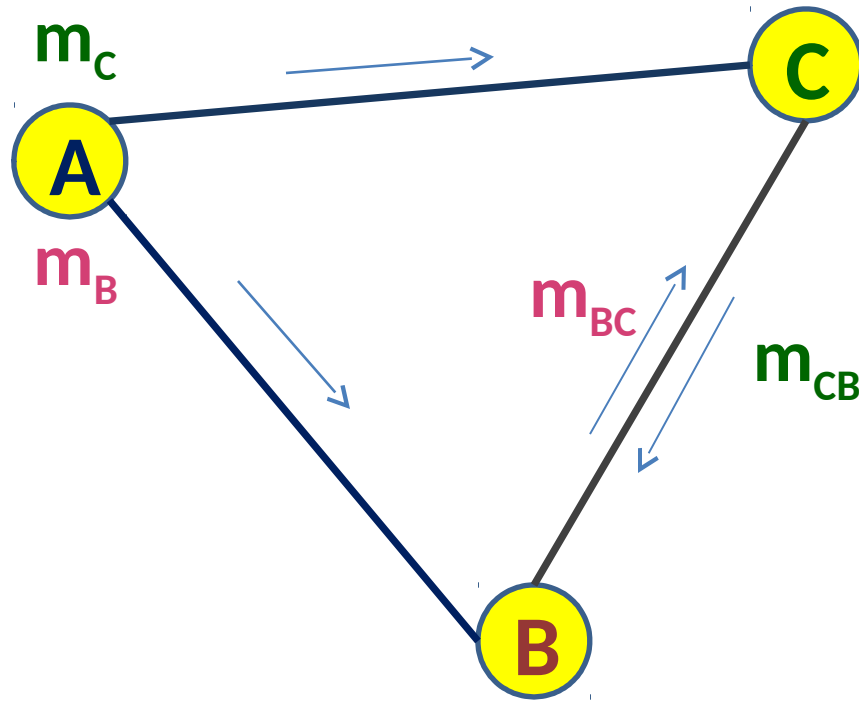
$$P(00 | D_1 U_2) = 0$$

$$P(00 | U_1 D_2) = 0$$

$$P(00 | U_1 U_2) = q > 0$$

# Byzantine Agreement (BA)

Generals of the **Byzantine Army** communicating with each other



If  $m_{CB} = m_{BC}$ , all are loyal.

The generals must reach a consensus among themselves whether to attack or retreat based on the messages exchanged.

But some generals can be **traitors**; they may send **conflicting messages** to the other generals.

# Byzantine Agreement

The solution to the problem must allow

- (i) all the loyal generals to agree upon a common plan of action.
- (ii) if the commanding general (A) is loyal then all the loyal generals must obey the order (s)he sends.
- No Classical solution [Fitzi et. al. CRYPTO 2001]



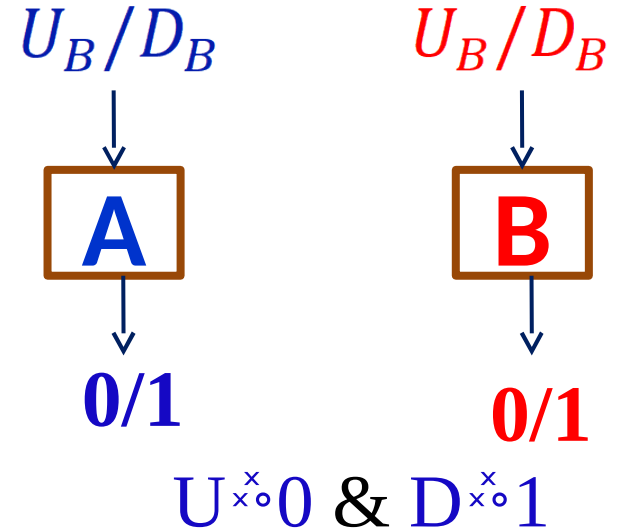
# **QUANTUM RANDOM NUMBER GENERATION**

# Randomness of a measurement's outcomes:

Randomness of the measurement outcomes (a,b) for the inputs (x,y) estimated by the **min-entropy function**

[R. Koenig et al., IEEE Trans. Inf. Theory, 09]

$$H_\infty(a, b|x, y) \equiv -\log_2[\text{Max}_{\{a,b\}} P(a, b|x, y)]$$



**Device-independent Case:** Used semidefinite programming (SDP)

**Minimize:**  $\text{Max}_{\{a,b\}} P(a, b|U_A, U_B)$

**Subject to:**  $\Delta_{\text{Hardy}} = [\Delta_{ij}] \geq 0$ ,    where  $\Delta_{ij} = \text{Tr}(E_i^\dagger E_j \rho)$

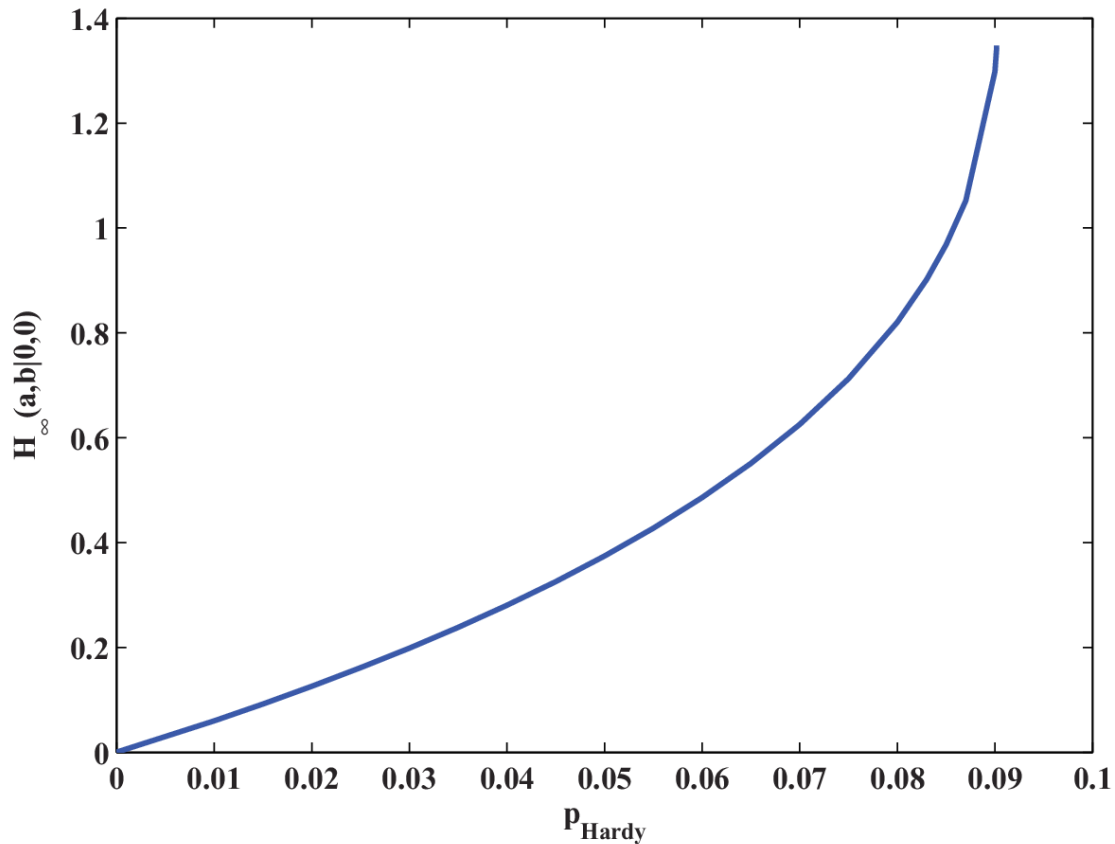
$P(0,0|U_A, U_B) = 0$ ,                       $E_i, E_j \in \{I, A_{a|x}, B_{b|y}, A_{a|x} B_{b|y}\}$

$P(0,1|D_A, U_B) = 0$ ,                      POVM  $\{A_{0|x}, A_{1|x}\} \rightarrow X \in \{U_A, D_A\}$

$P(1,0|U_A, U_B) = 0$ ,

$P(0,0|D_A, D_B) = p_{\text{Hardy}}$ .

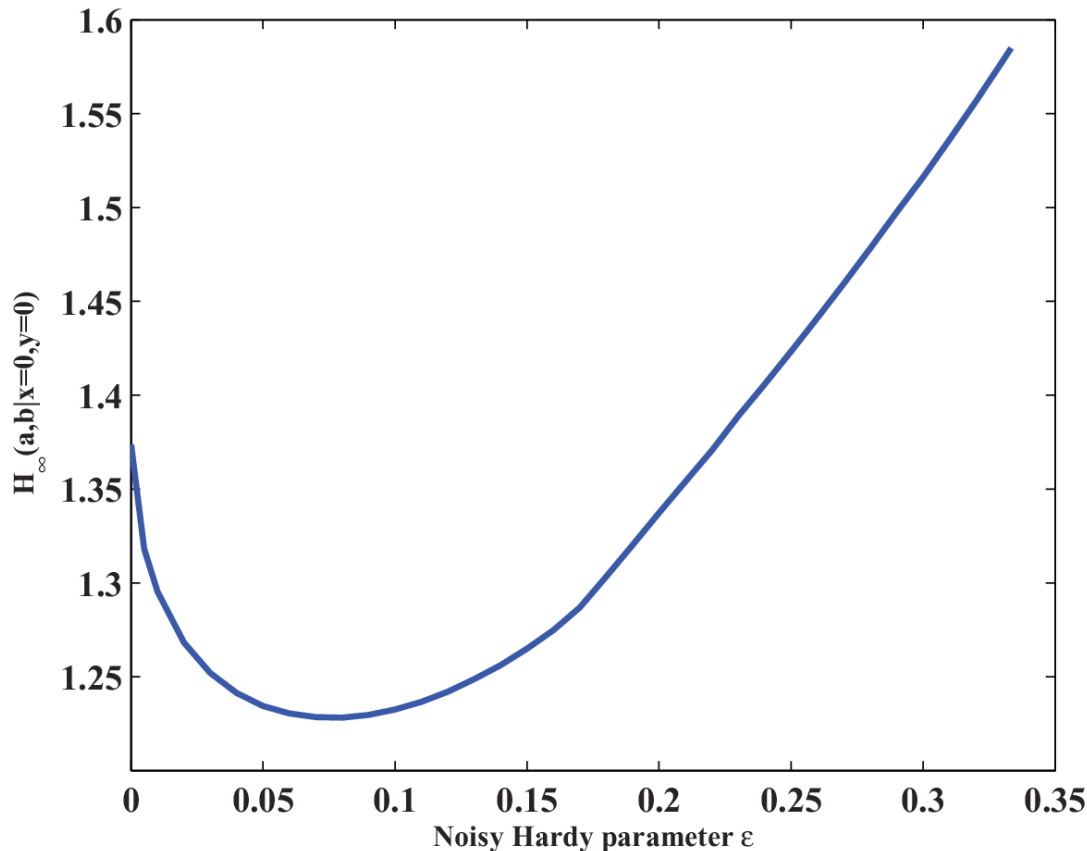
A lower bound on min-entropy  $H_\infty(a, b|0,0)$   
as a function of Hardy's parameter  $q = p_{\text{Hardy}}$ .



$$\begin{aligned} P(0,0|U_A, U_B) &= 0, \\ P(0,1|D_A, U_B) &= 0, \\ P(1,0|U_A, U_B) &= 0, \\ P(0,0|D_A, D_B) &= p_{\text{Hardy}}. \end{aligned}$$

Maximal randomness can reach up to 1.35 if the corresponding Hardy probability obtains its maximal value.

**Noisy Case:** A lower bound on min-entropy  $H_\infty(a, b|0,0)$  as a function noise parameter  $\epsilon$ .



## SDP Scheme

Minimize:  $\text{Max}_{\{a,b\}} P(a, b|U_A, U_B)$

Subject to:  $\Delta_{\text{noisyHardy}} \geq 0,$

$P(0,0|U_A, U_B) \leq \epsilon,$

$P(0,1|D_A, U_B) \leq \epsilon,$

$P(1,0|U_A, U_B) \leq \epsilon,$

$P(0,0|D_A, D_B) \geq q \geq 3\epsilon.$

Maximal randomness can reach up to **1.58** for  $\epsilon = 0.333$ .

# Cabello scenario: A lower bound on min-entropy

## SDP Scheme

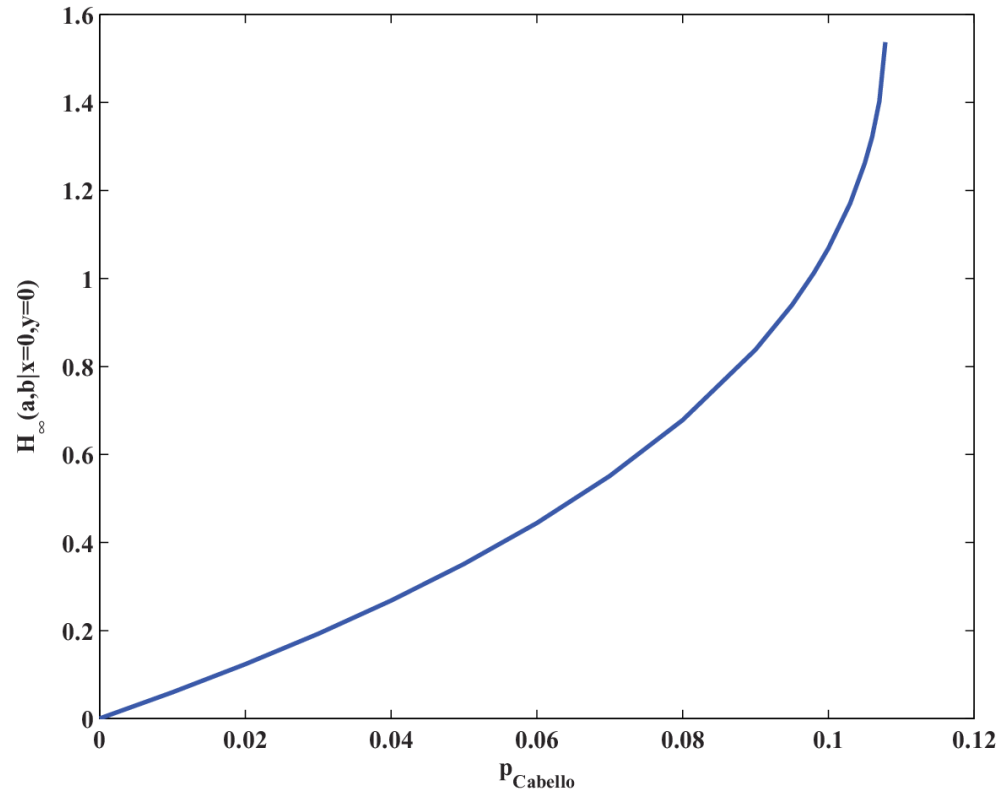
Minimize:  $\text{Max}_{\{a,b\}} P(a, b|U_A, U_B)$

Subject to:  $\Delta_{\text{Cabello}} = [\Delta_{ij}] \geq 0,$

$P(0,1|D_A, U_B) = 0,$

$P(1,0|U_A, U_B) = 0,$

$P(0,0|D_A, D_B) - P(0,0|U_A, U_B) = p.$



Maximal randomness can reach up to **1.56** when Cabello parameter ( $p$ ) reaches its maximal value **0.10784**.

# Semi-Device Independent scenario

## SDP Scheme

Minimize:  $\text{Max}_{\{a,b\}} P(b|U_A, a, U_B)$

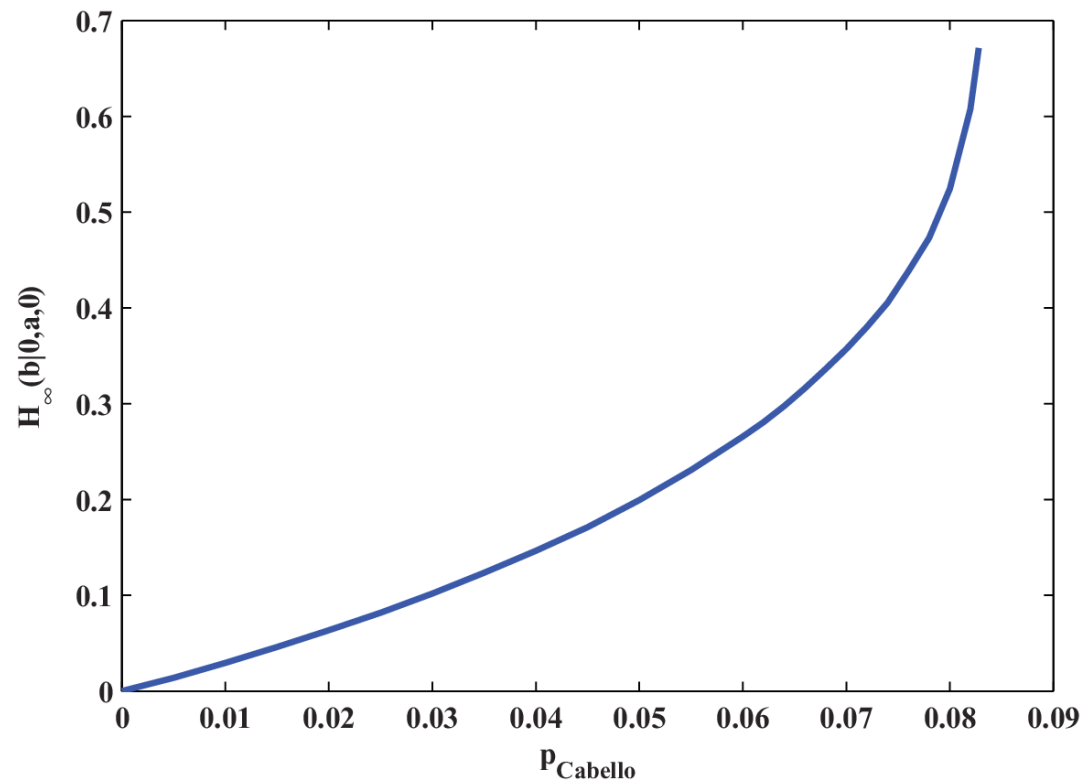
Subject to:  $\Delta_{\text{Cabello}} = [\Delta_{ij}] \geq 0,$

$P(0,1|D_A, U_B) = 0,$

$P(1,0|U_A, U_B) = 0,$

$P(a|x) = P(a|x, y) = \frac{1}{2}; x \in \{U, D\},$

$P(0,0|D_A, D_B) - P(0,0|U_A, U_B) = p.$



Maximal randomness can reach up to **0.68**. Other existing protocols can generate a maximum **0.23** bit randomness.

# CONCLUSIONS

- Proposed generalized Hardy type test for detection of multiparty entanglement.
- Based on Hardy correlations, we have proposed Device-Independent quantum protocols for
  - Key Distribution.
  - Liar Detection & Byzantine Agreement
  - Random Number Generator
  - Quantum Digital Signatures etc.

## References:-

- **R. Rahaman**, M. G. Parker, P. Mironowicz & M. Pawłowski, '*Device-independent quantum key distribution based on measurement inputs*', Phys. Rev. A, 92, 062304(2015).
- **R. Rahaman**, M. Wiesniak & M. Zukowski, '*Quantum Byzantine agreement via Hardy correlations and entanglement swapping*', Phys. Rev. A, 92, 042302(2015).
- H.-W. Li, M. Pawłowski, **R. Rahaman**, G. Guo & Z.-F. Han, '*Device- and semi-device-independent random numbers based on noninequality paradox*', Phys. Rev. A 92, 022327(2015).
- S. S. Bhattacharya, A. Roy, A. Mukherjee & **R. Rahaman**, '*All-versus-nothing violation of local realism from the Hardy paradox under no-signaling*', Phys. Rev. A 92, 012111(2015).
- **R. Rahaman**, M. Wiesniak & M. Zukowski, '*True Multipartite Entanglement Hardy Test*', Phys. Rev. A, 90, 062338(2014).
- **R. Rahaman**, '*Quantum digital signatures based non-locality test*' (Under Preparation).



**Thank You**