# Efficient measurement of high-dimensional quantum states

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**QIPA 2018, HRI, Prayagraj** 05-Dec-2018

#### Outline

### **Orbital Angular Momentum of light**

- What it means?
- Generation of OAM states of light

#### **Efficient measurement of states of light**

- in the orbital angular momentum (OAM) basis.
- in the transverse position basis.

#### **Orbital angular momentum (OAM) of light**

Laguerre-Gaussian (LG) modes are solutions to paraxial Wave Equation

$$LG_{p}^{l}(\rho,\phi,z) = \frac{C}{\left(1+z^{2}/z_{R}^{2}\right)^{1/2}} \exp\left[i(2p+l+1)\tan^{-1}\left(\frac{z}{z_{R}}\right)\right] \left[\frac{\rho/2}{w(z)}\right]^{l} L_{p}^{l} \left[\frac{2\rho^{2}}{w^{2}(z)}\right]$$

$$\times \exp\left[-\frac{\rho^{2}}{w^{2}(z)}\right] \exp\left[-\frac{ik^{2}\rho^{2}z}{2(z^{2}+z_{R}^{2})}\right] e^{-il\phi}$$
Transverse phase of the LG mode with p=0  
Yao and Padgett, Advances in optics and photonics, **3**, 161 (2011)  

$$p=2$$

$$p=0$$

$$p=0$$

$$l=0$$

$$l=1$$

$$p=0$$

$$l=0$$

- Orbital angular momentum per photon in an LG mode: *ħl* Allen et al., PRA 45, 8185 (1992)
- OAM modes  $\psi_l(\phi) = \frac{1}{\sqrt{2\pi}} e^{-il\phi}$  form a complete basis & provide an infinite dimensional discrete basis

(i) higher allowed error rate in cryptography

(ii) higher transmission bandwidth

(iii) Fundamental tests of quantum mechanics

Phys. Rev. Lett. 88, 127902 (2002).

Phys. Rev. Lett. 90, 167906 (2003).

Phys. Rev. Lett. 85, 4418–4421 (2000).

ARTICLES PUBLISHED ONLINE: 24 JUNE 2012 | DOI: 10.1038/NPHOTON.2012.138 nature photonics

## Terabit free-space data transmission employing orbital angular momentum multiplexing

Jian Wang<sup>1,2</sup>\*, Jeng-Yuan Yang<sup>1</sup>, Irfan M. Fazal<sup>1</sup>, Nisar Ahmed<sup>1</sup>, Yan Yan<sup>1</sup>, Hao Huang<sup>1</sup>, Yongxiong Ren<sup>1</sup>, Yang Yue<sup>1</sup>, Samuel Dolinar<sup>3</sup>, Moshe Tur<sup>4</sup> and Alan E. Willner<sup>1</sup>\*



#### ARTICLE

Received 17 Mar 2014 | Accepted 1 Aug 2014 | Published 16 Sep 2014

DOI: 10.1038/ncomms5876

OPEN

### High-capacity millimetre-wave communications with orbital angular momentum multiplexing

Yan Yan<sup>1,\*</sup>, Guodong Xie<sup>1,\*</sup>, Martin P.J. Lavery<sup>2,\*</sup>, Hao Huang<sup>1,\*</sup>, Nisar Ahmed<sup>1</sup>, Changjing Bao<sup>1</sup>, Yongxiong Ren<sup>1</sup>, Yinwen Cao<sup>1</sup>, Long Li<sup>1</sup>, Zhe Zhao<sup>1</sup>, Andreas F. Molisch<sup>1</sup>, Moshe Tur<sup>3</sup>, Miles J. Padgett<sup>2</sup> & Alan E. Willner<sup>1</sup>

#### 1. Using an spiral phase plate:





Oemrawsingh et al., Applied Optics, 43, 688 (2004) Yao and Padgett, Advances in optics and photonics, **3**, 161 (2011)



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**States in OAM basis :**  $\langle \boldsymbol{\phi} | \boldsymbol{l} \rangle = e^{-il\phi}$ 

When different OAM eigenmodes are uncorrelated.  $\langle \alpha_{l_1} \alpha_{l_2}^* \rangle_e = S_{l_1} \delta_{l_1, l_2}$ 

$$W(\phi_1,\phi_2) = \frac{1}{2\pi} \sum_{l=-\infty}^{\infty} S_l e^{-il(\phi_1 - \phi_2)} \iff \rho = \frac{1}{2\pi} \sum_{l=-\infty}^{\infty} S_l |l\rangle \langle l|$$

Diagonal **Mixed States** 

 $W(\Delta \phi) = \frac{1}{2\pi} \sum_{l=-\infty}^{\infty} S_l e^{-il\Delta \phi} \implies S_l = \int_{-\pi}^{\pi} W(\Delta \phi) e^{il\Delta \phi} d\phi \qquad \text{Angular Wiener-Khintchine theorem}$ Angular  $A \times Jha, G \times Jha, G$ 

A K Jha, G S Agarwal, R W Boyd, PRA 84, 063847 (2011)

Angular correlation spectrum function

Aim: Measure the angular correlation function  $W(\phi_1, \phi_2)$ For diagonal states it yields the OAM spectrum

#### **Existing methods for measuring OAM spectrum of Light**



• Very sensitive to noise

#### Measuring Orbital Angular Momentum of Light (A new scheme)

Input state in the OAM basis (diagonal mixed state)  $\infty$ 

$$\psi_{\rm in}(\phi) = \frac{1}{\sqrt{2\pi}} \sum_{l=-\infty} \alpha_l e^{-il\phi}$$

#### with $\langle \alpha_l \, \alpha_{l'}^* \rangle_e = S_l \, \delta_{l,l'}$

A reflection flips the wave-front along the reflection axis





#### Measuring Orbital Angular Momentum of Light (A new scheme)

Input state in the OAM basis (diagonal mixed state)

$$\psi_{\text{in}}(\phi) = \frac{1}{\sqrt{2\pi}} \sum_{l=-\infty}^{\infty} \alpha_l e^{-il\phi} \quad \text{with} \quad \langle \alpha_l \, \alpha_{l'}^* \rangle_e = S_l \, \delta_{l,l'}$$

$$\psi_{\text{out}}(\phi) = \sqrt{\frac{k_1}{2\pi}} \sum_{l=-\infty}^{\infty} \alpha_l e^{-il\phi+i\omega t_1} + \sqrt{\frac{k_2}{2\pi}} \sum_{l=-\infty}^{\infty} \alpha_l e^{+il\phi+i\omega t_2}$$

$$I_{\text{out}}(\phi) = \langle \psi_{\text{out}}(\phi)\psi^*_{\text{out}}(\phi)\rangle_e$$

$$= \frac{k_1}{2\pi} + \frac{k_2}{2\pi} + 2\sqrt{k_1k_2}W(2\phi)\cos\delta$$
where  $\delta \equiv \omega(t_1 - t_2)$ 

$$W(2\phi) \equiv \frac{1}{2\pi} \sum_{l=-\infty}^{\infty} S_l \, e^{-il2\phi}$$

- $W(2\phi)$  is encoded in the interferogram. A single-shot measurement of  $I_{out}(\phi)$  yields  $W(2\phi)$ From  $W(\Delta\phi)$ ,  $S_l$  can be computed, in a single shot manner.  $S_l = \int_{-\infty}^{\pi} W(2\phi)e^{il2\phi}d\phi$
- Still sensitive to background noise and other experimental parameters

#### Measuring Orbital Angular Momentum of Light (A new scheme)

Input state in the OAM basis (diagonal mixed state)

$$\psi_{\text{in}}(\phi) = \frac{1}{\sqrt{2\pi}} \sum_{l=-\infty}^{\infty} \alpha_l e^{-il\phi} \quad \text{with} \quad \langle \alpha_l \, \alpha_l^* \rangle_e = S_l \, \delta_{l,l'}$$
No noise:  $I_{\text{out}}(\phi) = \frac{k_1 + k_2}{2\pi} + 2\sqrt{k_1 k_2} W(2\phi) \cos \delta$ 
With noise:  $I^{\delta}_{\text{out}}(\phi) = I^{\delta}_n(\phi) + \frac{k_1 + k_2}{2\pi} + 2\sqrt{k_1 k_2} W(2\phi) \cos \delta$ 
Difference intensity:  $\Delta I_{\text{out}}(\phi) \equiv I^{\delta c}_{\text{out}}(\phi) - I^{\delta d}_{\text{out}}(\phi)$ 
 $\Delta I_{\text{out}}(\phi) = \Delta I_n(\phi) + 2\sqrt{k_1 k_2} W(2\phi) (\cos \delta_c - \cos \delta_d)$ 
If shot-to-shot noise is the same:  $\Delta I_n(\phi) = 0$  Then:  
 $\Delta I_{\text{out}}(\phi) = 2\sqrt{k_1 k_2} (\cos \delta_c - \cos \delta_d) W(2\phi) \propto W(2\phi)$ 

•  $\Delta I_{out}(\phi)$  has the same functional form as  $W(2\phi)$ .

• So by measuring  $\Delta I_{out}(\phi)$  the spectrum  $S_l$  can be obtained in a single-shot as well as in a noise-insensitive manner  $S_l = \int_{-\pi}^{\pi} W(2\phi) e^{il2\phi} d\phi$ 11

G. Kulkarni, R. Sahu, O. S. Magana-Loaiza, R. W. Boyd, and A. K. Jha; Nature Communications, 8, 1054 (2017)

#### **Experimental measurement of OAM spectrum of Light**



G. Kulkarni, R. Sahu, O. S. Magana-Loaiza, R. W. Boyd, and A. K. Jha; Nature Communications, 8, 1054 (2017)

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G. Kulkarni, R. Sahu, O. S. Magana-Loaiza, R. W. Boyd, and A. K. Jha; Nature Communications, 8, 1054 (2017)

#### **Measuring Orbital Angular Momentum of Light (Quantum)**



The state of the signal photon is

$$\rho_{S} = \operatorname{tr}_{i}(|\psi\rangle\langle\psi|) = \sum_{l=-\infty}^{\infty} S_{l} |l\rangle_{ss}\langle l|$$

Angular coherence function of the signal photon is

$$W_s(\phi_1, \phi_2) \to W_s(\Delta \phi) = \frac{1}{2\pi} \sum_{l=-\infty}^{\infty} S_l e^{-il\Delta \phi}$$

 $\sim$ 

Pump Signal  $|\psi\rangle = \sum_{l=-\infty}^{\infty} \sqrt{S_l} |l\rangle_s |-l\rangle_i$ Idler OAM-Entangled State:

- $S_l$  is called the angular Schmidt spectrum
- Accurate measurement of S<sub>l</sub> is very important
   for quantification of OAM entanglement.
  - The current methods involve coincidence measurements, which is very difficult.

Nature 412 **313** (2001) Phys Rev A **76**, 042302 (2007) Phys Rev Lett **104**, 020505 (2010) New J Phys **14**, 073046 (2012)

• The OAM spectrum of the signal photon is same as the angular Schmidt spectrum of the entangled two-photon state.

14

G. Kulkarni, R. Sahu, O. S. Magana-Loaiza, R. W. Boyd, and A. K. Jha; Nature Communications, 8, 1054 (2017)

#### **Measuring Orbital Angular Momentum of Light (Quantum)**



G. Kulkarni, R. Sahu, O. S. Magana-Loaiza, R. W. Boyd, and A. K. Jha; Nature Communications, 8, 1054 (2017)

#### **Measuring Orbital Angular Momentum of Light (Quantum)**



G. Kulkarni, S. Aarav, L. Taneja, and A. K. Jha, **Phys. Rev. A 97, 063846 (2018)** 

#### **Angular correlation function: Application**



G. Kulkarni, S. Aarav, L. Taneja, and A. K. Jha, Phys. Rev. A 97, 063846 (2018)

#### **Spatial partially correlated fields**

• Partially coherent fields are extremely important for imaging through scattering, etc.



Nature Photonics 6, 355 (2012).

#### uncorrelated transverse wavevectors

 $\langle a^*(\boldsymbol{q}_1)a(\boldsymbol{q}_2)\rangle_e = I(\boldsymbol{q}_1)\delta(\boldsymbol{q}_1 - \boldsymbol{q}_2)$ **Diagonal mixed state**  $W(\boldsymbol{\Delta}\boldsymbol{\rho}, z) = \int I(\boldsymbol{q})e^{-i\boldsymbol{q}\cdot\boldsymbol{\Delta}\boldsymbol{\rho}} d\boldsymbol{q}$ 

Propagation-invariant spatially partially coherent field



S. Aarav, A. Bhattacharjee, H. Wanare, and A. K. Jha, PRA 96, 033815 (2017)

#### **Spatial partially correlated fields: efficient measurement**



A. Bhattacharjee, S. Aarav, and A. K. Jha, APL 113, 051102 2018)

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#### **Spatial partially correlated fields: efficient measurement**



A. Bhattacharjee, S. Aarav, and A. K. Jha, APL 113, 051102 (2018)

#### **Spatial partially correlated fields: Application**



A. Bhattacharjee, S. Aarav, and A. K. Jha; arXiv: 1810.00686

#### Conclusions

- Demonstrated a single-shot technique for measuring the angular correlation function.
- For diagonal mixed states, the angular correlation function yields the OAM spectrum through a Fourier transform.
- The technique can be used for measuring the angular Schmidt spectrum of OAM-entangled states in a single-shot manner without requiring coincidence detection.
- The technique can be extended for measuring the spatial correlation functions.

#### Acknowledgements

- IIT Kanpur & SERB DST
- Quantum Optics and Entanglement Group





Lavanya



Shaurya





Swati

- Opening for PhD Students (Theory + Experiment)
- Opening for Postdocs (60K p/m + ; Deadline 15<sup>th</sup> Jan 2019)
   Thank you for your attention