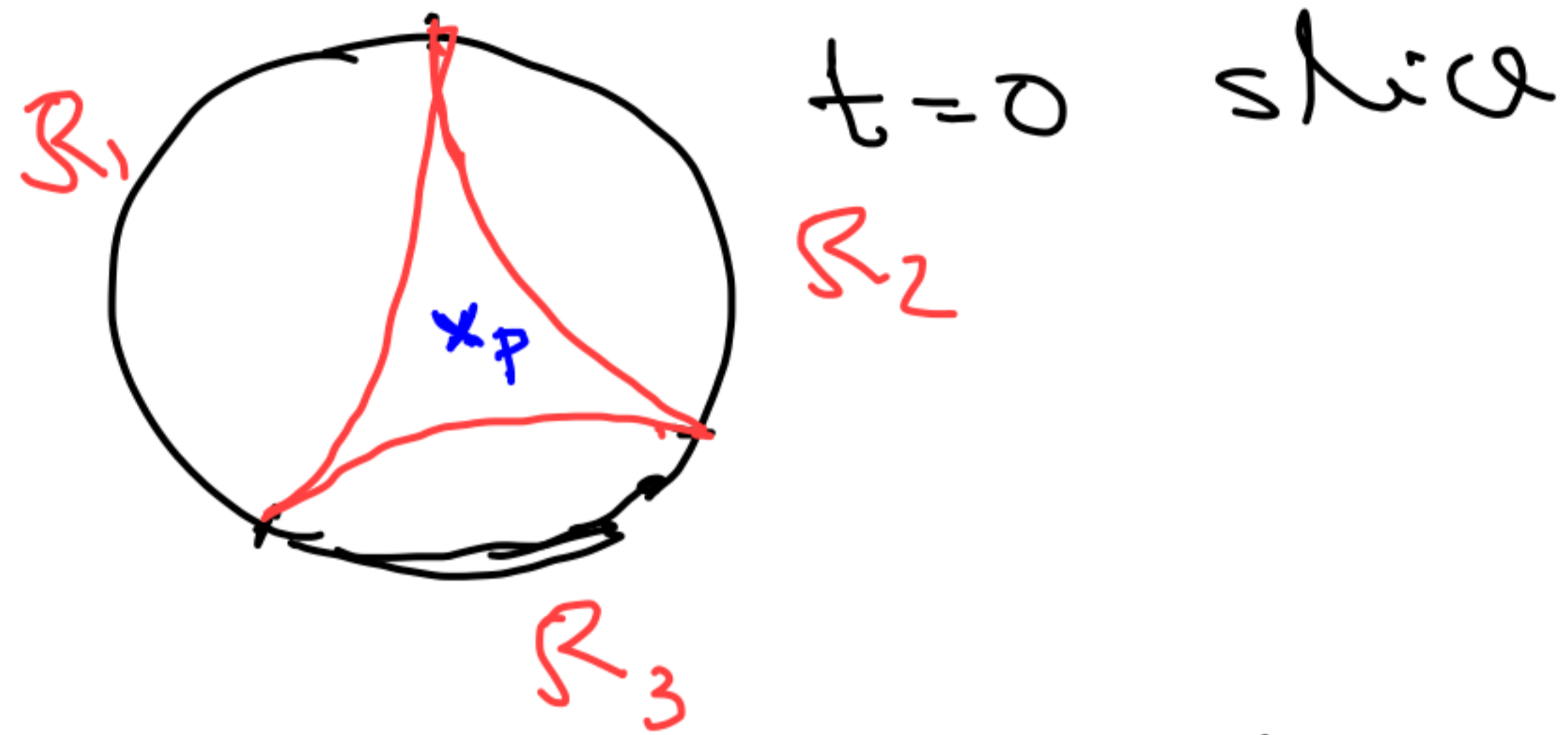


# The situation in AdS/CFT



None of  $R_1$ ,  $R_2$  or  $R_3$  contains information about  $P$  but  $R_1 \cup R_2$ ,  $R_2 \cup R_3$  or  $R_3 \cup R_1$  do.

Q. How is the information encoded?

We'll model the system using a quantum error correction code.

Encode a state  $|\psi\rangle \in \mathcal{H}$  in a bigger Hilbert space such that even if we lose some part of the information on the bigger Hilbert space we can recover  $|\psi\rangle$ .

Note: We cannot copy  $|\psi\rangle$  (no-cloning theorem)  
 $|\psi\rangle \rightarrow |\psi\rangle \otimes |\psi\rangle \otimes |\psi\rangle$  . X. not possible.



We'll work with qudits  $\rightarrow$  a  $d$ -state system ( $d$   $\neq$  dimension of space).  
Hilbert space in which the original state  $|\psi\rangle$  lives is taken to be a product of  $k$  qudits.  $\Rightarrow d^k$  dimensional space.  
 $\Rightarrow$  logical qudits (bulk of AdS).  
Hilbert space in which we encode the state  $|\psi\rangle$  as  $|\tilde{\psi}\rangle$  is taken to be a product of  $n$  qudits ( $n > k$ ). <sup>CFT</sup>  
 $\Rightarrow$  physical qudits (boundary)

$|4\rangle \rightarrow |\tilde{4}\rangle$  a linear map.

Code subspace: Subspace of states of the physical qudits spanned by the  $|\tilde{\psi}\rangle$ 's. ( $d^k$  dimensional subspace of  $d^n$  dimensional space).

We'll illustrate this using the 2-qubit code.  $|0\rangle, |1\rangle, |2\rangle \rightarrow 3$  states of a qubit. Encode one logical qubit into 3 physical qubits.

$$|\psi\rangle = \sum_{i=0}^2 c_i |i\rangle \quad \rightarrow \quad |\tilde{\psi}\rangle = \sum_{i=0}^2 c_i |\tilde{i}\rangle$$

$|\tilde{0}\rangle, |\tilde{1}\rangle, |\tilde{2}\rangle$  span the code subspace.

Hilbert space of physical qubits

is  $3^3 = 27$  dimensional.

Prescriptions:

$$|\tilde{0}\rangle = \frac{1}{\sqrt{3}} (|1000\rangle + |1111\rangle + |2222\rangle)$$

$$|\tilde{1}\rangle = \frac{1}{\sqrt{3}} (|1012\rangle + |1201\rangle + |2011\rangle)$$

$$|\tilde{2}\rangle = \frac{1}{\sqrt{3}} (|1021\rangle + |1102\rangle + |2101\rangle)$$

Invariant  
under  
cyclic  
perm of  
the phys.  
qubits.



$$|\tilde{\alpha}\rangle = U_{12} (|i\rangle_1 \otimes |X\rangle_{23})$$

$$|X\rangle_{23} = \frac{1}{\sqrt{3}} (|00\rangle_{23} + |11\rangle_{23} + |22\rangle_{23})$$

$U_{12}$ : Unitary operator acting on the

12 Hilbert space:

$$U: |00\rangle \rightarrow |00\rangle$$

$$|11\rangle \rightarrow |20\rangle \rightarrow |21\rangle \rightarrow |02\rangle \rightarrow |22\rangle \rightarrow |10\rangle$$

$$\rightarrow |12\rangle \rightarrow |01\rangle \rightarrow |11\rangle$$

$$|\tilde{\alpha}\rangle = U_{12} (|0\rangle_1 \otimes \frac{1}{\sqrt{3}} (|00\rangle_{23} + |11\rangle_{23} + |22\rangle_{23}))$$

$$|\tilde{i}\rangle = U_{12} (|i\rangle_1 \otimes |X\rangle_{23})$$

$$\begin{aligned} |\tilde{i}\rangle &= U_{23} (|i\rangle_2 \otimes |X\rangle_{31}) \\ |\tilde{i}\rangle &= U_{31} (|i\rangle_3 \otimes |X\rangle_{12}) \end{aligned} \quad \left. \vphantom{\begin{aligned} |\tilde{i}\rangle &= U_{23} (|i\rangle_2 \otimes |X\rangle_{31}) \\ |\tilde{i}\rangle &= U_{31} (|i\rangle_3 \otimes |X\rangle_{12}) \end{aligned}} \right\} \begin{array}{l} \text{follow} \\ \text{from cyclicity.} \end{array}$$

$$U_{23} \frac{1}{\sqrt{3}} \sum_{k=0}^2 |k\rangle_1 |i\rangle_2 |k\rangle_3$$

Recovery

$$U_{12}^\dagger |\tilde{i}\rangle = |i\rangle_1 \otimes |X\rangle_{23}$$

Apply  $U_{12}^\dagger$  and read out the state in the first qubit.

$$U_{12}^\dagger |\psi\rangle = \sum_i c_i |i\rangle_1 \otimes |X\rangle_{23}$$

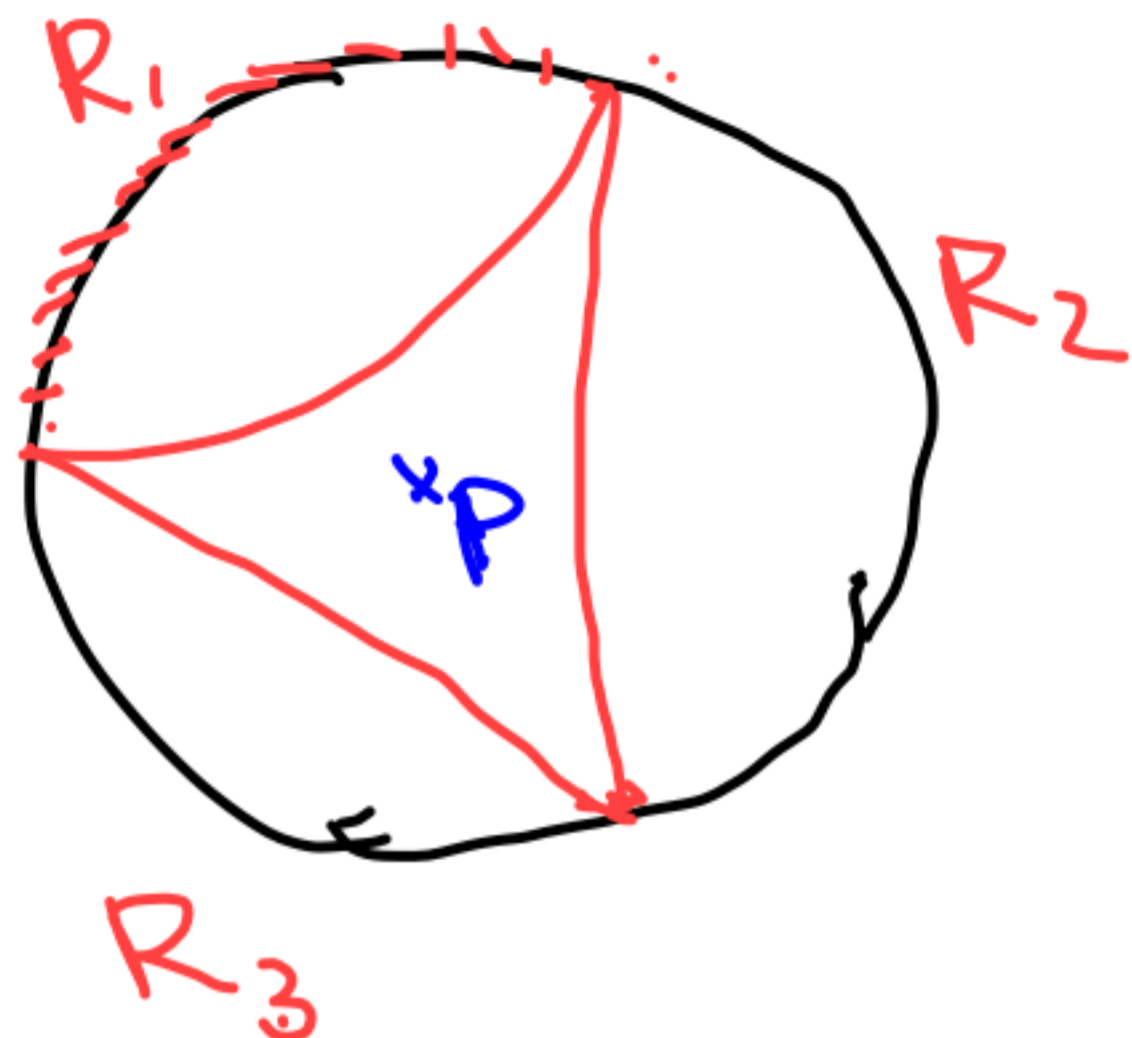
$$\sum_i c_i |i\rangle_2$$

$$U_{23}^\dagger |\psi\rangle = \sum_i c_i |i\rangle_2 \otimes |X\rangle_{31}$$

$$U_{31}^\dagger |\psi\rangle = \sum_i c_i |i\rangle_3 \otimes |X\rangle_{12}$$

$\Rightarrow$  If we lose any one of the 3 qubits we can still recover the information.





Logical operator:  $G|i\rangle = \sum_{j=0}^2 G_{ji}|j\rangle$   
 & numbers.

Represent  $G$  on physical qubits.

→ Want to construct  $\tilde{G}$  such that  
 $\tilde{G}|\tilde{i}\rangle = \sum_{j=0}^2 G_{ji}|\tilde{j}\rangle$

Can choose

$$\tilde{G}_{12} = U_{12} G_1 U_{12}^\dagger$$

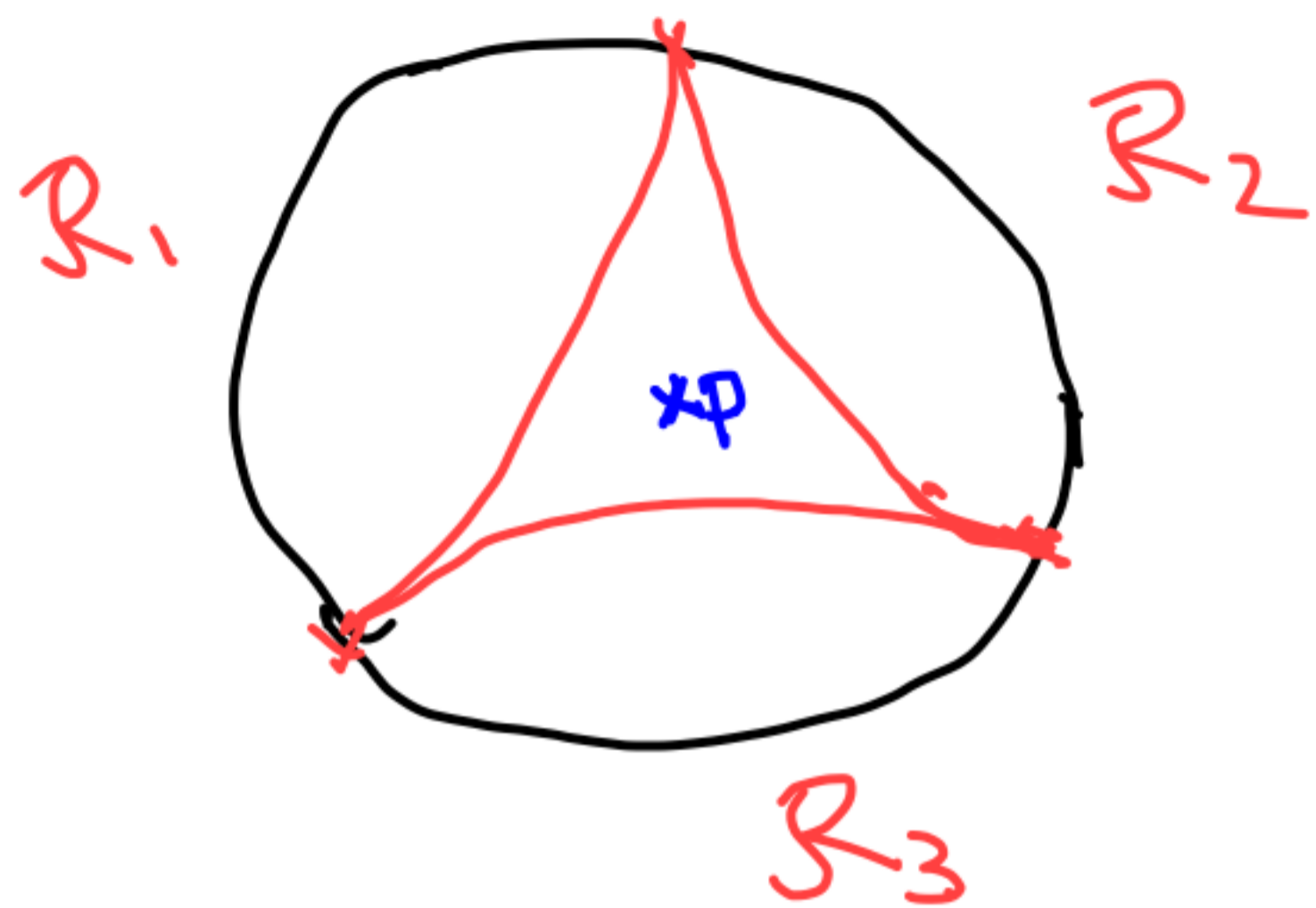
$$\tilde{G}_{12} |i\rangle = U_{12} G_1 U_{12}^\dagger U_{12} (|i\rangle_1 \otimes |X\rangle_{23})$$

$$= U_{12} \sum_j G_{ji} |j\rangle_1 \otimes |X\rangle_{23}$$

$$= \sum_j G_{ji} |j\rangle$$

$$\tilde{G}_{23} = U_{23} G_2 U_{23}^\dagger, \quad \tilde{G}_{31} = U_{31} G_3 U_{31}^\dagger$$

- Different operators but they act on the code subspace in the same way.



Take a logical operator  $G$ .

$X_3$ : some operator on  $R_3$  (3rd qubit)

In what sense is  $[X_3, G] = 0$ ?

$\langle \vec{v} | [X_3, G] | \vec{\phi} \rangle = 0$  if  $|\vec{v}\rangle$  and  $|\vec{\phi}\rangle$  are in the code subspace.



$$\langle \vec{f} | [X_3, G] | \vec{\phi} \rangle$$

$$= \langle \vec{f} | X_3 G | \vec{\phi} \rangle - \langle \vec{f} | G X_3 | \vec{\phi} \rangle$$

$$= \langle \vec{f} | X_3 G_{12} | \vec{\phi} \rangle - \langle \vec{f} | X_3 G_{21} | \vec{\phi} \rangle$$

$$= \langle \vec{f} | X_3 G_{12} | \vec{\phi} \rangle - \langle \vec{f} | X_3 G_{12} | \vec{\phi} \rangle$$

$$= \langle \vec{f} | [X_3, G_{12}] | \vec{\phi} \rangle = 0$$

Define:

$$T_{i_1 i_2 i_3} = \sqrt{3} \left\langle i_1 | \otimes \left\langle i_2 | \otimes \left\langle i_3 | \right\rangle \right.$$

basis in code subspace.

$T_{i_1 i_2 i_3}$  is a "perfect tensor."

$T_{i_1 \dots i_n}$  is called a perfect tensor

if any balanced bipartition of indices

give a unitary matrix.

$M_{i_1 i_2 i_3} = T_{i_1 i_2 i_3}$  should be unitary

Proof.

$$|X\rangle_{23} = \sum_k \frac{1}{\sqrt{3}} |k\rangle_2 \otimes |k\rangle_3$$

$$T_{i_1 i_2 i_3 j} = \sum_{k=0}^2 \langle i_1 | \otimes \langle i_2 | \otimes \underbrace{\langle i_3 | U_{12} | j \rangle}_{|j\rangle} \otimes |k\rangle_2 \otimes |k\rangle_3$$

$$\langle i_3 | k \rangle_3 = \delta_{i_3 k}$$

$$= \langle i_1 | \otimes \langle i_2 | U_{12} | j \rangle \otimes |i_3\rangle_2$$

$$= (U_{12})_{i_1 i_2, j i_3} \rightarrow \text{Unitary matrix.}$$

→ Similar analysis holds for other partitions.



