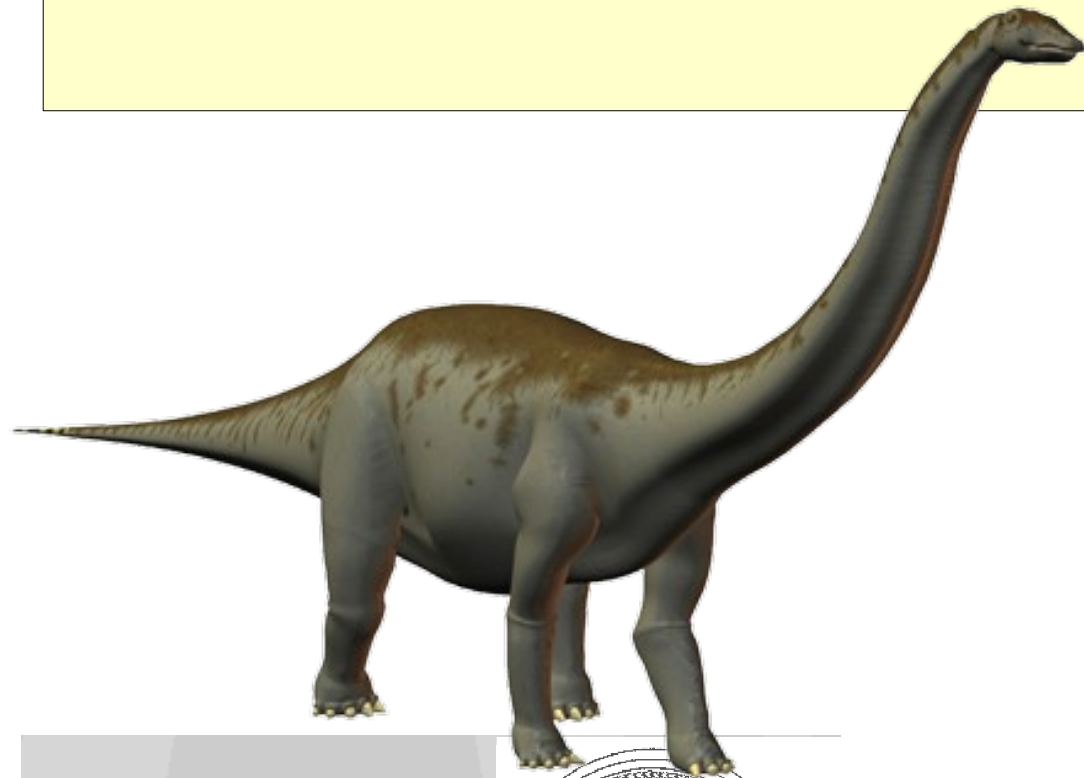


Stronger uncertainty relations



Lorenzo Maccone

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Arun K. Pati

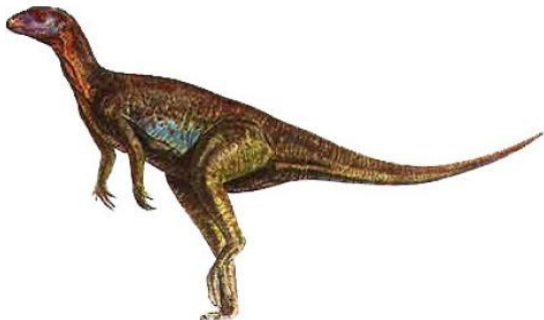
Quantum Information and
Computation Group, Harish-
Chandra Research Institute,
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What I'm going to talk about

Heisenberg uncertainty relations do NOT convey fully the notion of complementarity

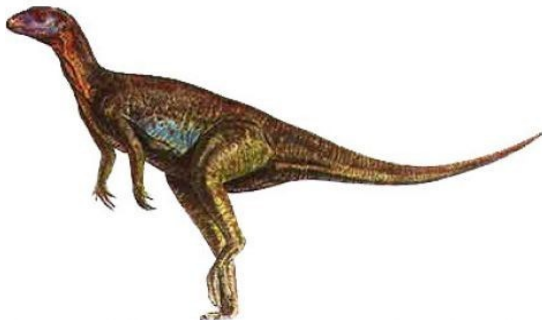


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Heisenberg uncertainty relations do NOT convey fully the notion of complementarity

$$\Delta A \Delta B \geq \frac{1}{2} |\langle [A, B] \rangle|$$

both terms can be zero, even if
A and B are incompatible



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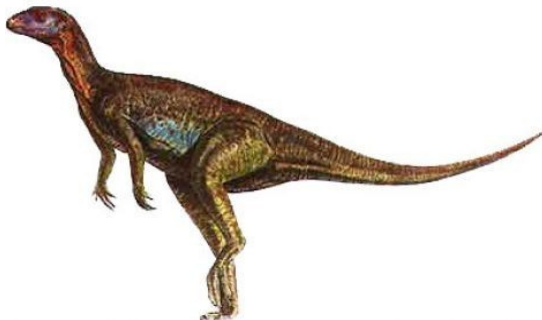
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$$0 \geq 0$$



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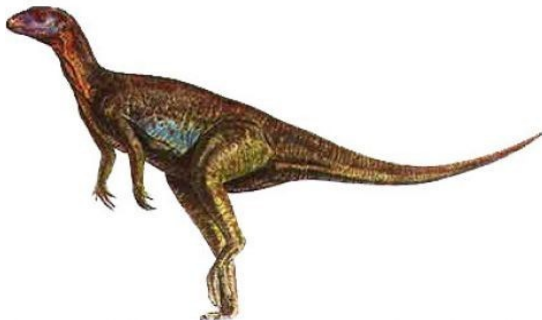
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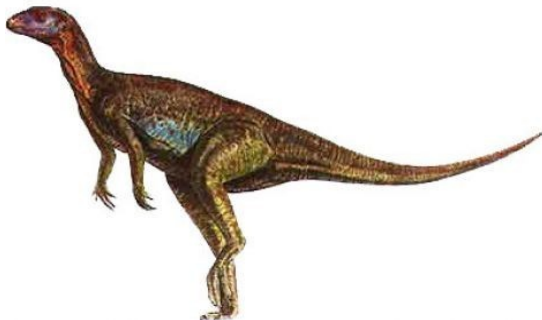
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$$0 \geq 0$$

Can we do better?  **YES!**



Reminder



Reminder



Heisenberg-Robertson UR

$$\Delta A^2 \Delta B^2 \geq \left| \frac{1}{2} \langle [A, B] \rangle \right|^2$$

variances and expectations calculated on system state $|\psi\rangle$

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Schroedinger UR

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both trivial if $|\psi\rangle$ = eigenstate of A **OR** B

we need to give a bound to the *sum*

$$\Delta A^2 + \Delta B^2 \geq \dots$$

(the product $\Delta A^2 \Delta B^2$ is trivially zero if one of the two is)



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$$\Delta A^2 + \Delta B^2 \geq 2\Delta A\Delta B$$



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nice, but...

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Heisenberg

Main result:
two lower bounds



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nonzero except when $|\psi\rangle$
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$$\Delta A^2 + \Delta B^2 \geq \mathcal{M}$$

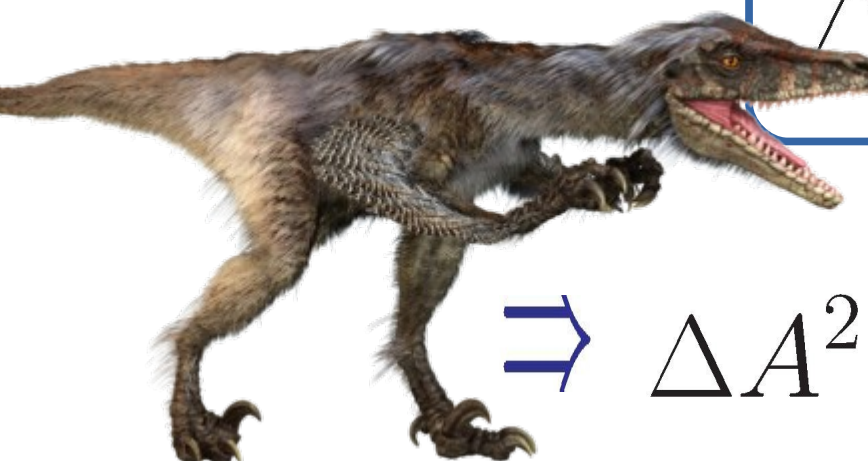


Main result: two lower bounds

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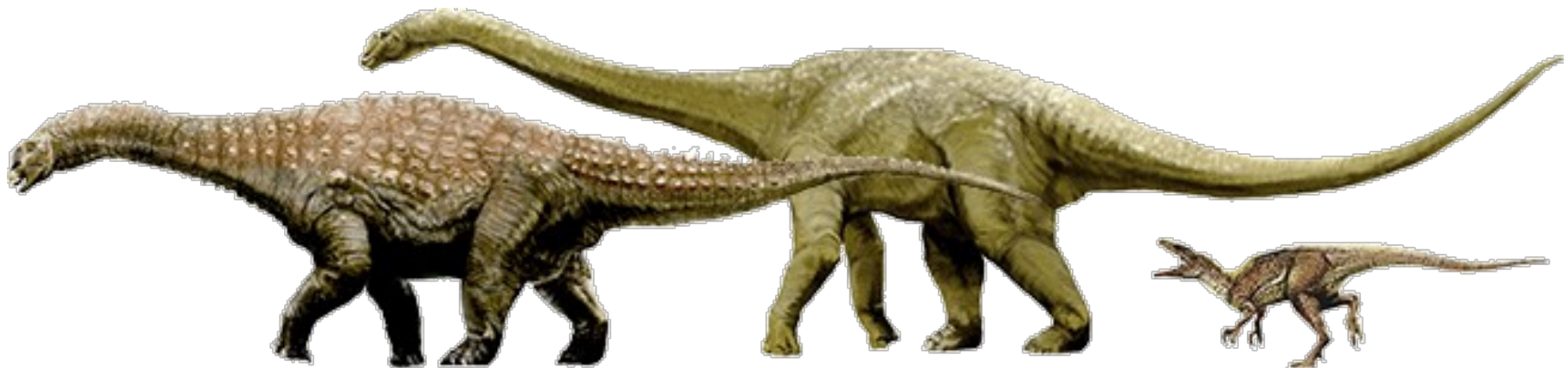
$$\Delta A^2 + \Delta B^2 \geq \mathcal{M}$$


$$\Rightarrow \Delta A^2 + \Delta B^2 \geq \max(\mathcal{L}, \mathcal{M})$$

SUM of variances

$$\Delta A^2 + \Delta B^2 \geq \dots \Rightarrow$$

dimensional regularization needed
if A and B have different units

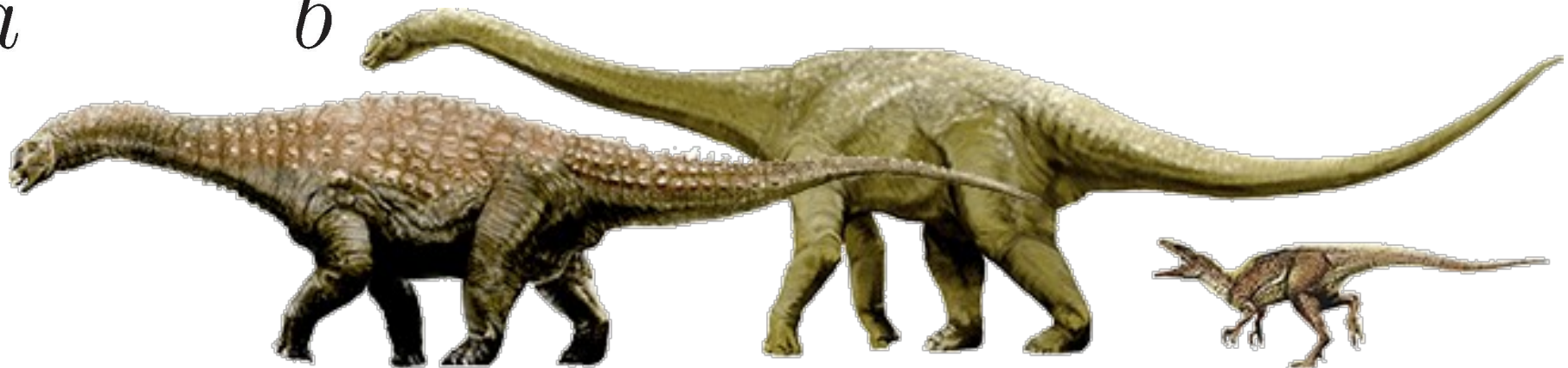


SUM of variances

$$\Delta A^2 + \Delta B^2 \geq \dots \Rightarrow$$

dimensional regularization needed
if A and B have different units

$$\frac{\Delta A^2}{a} + \frac{\Delta B^2}{b} \geq \dots$$



First bound:

$$\Delta A^2 + \Delta B^2 \geq \pm i \langle [A, B] \rangle + |\langle \psi | A \pm iB | \psi^\perp \rangle|^2$$



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$|\psi^\perp\rangle$ arbitrary state orthogonal to the system state $|\psi\rangle$

constructive procedure to choose $|\psi^\perp\rangle$, or just choose a random one

First bound:

$$\Delta A^2 + \Delta B^2 \geq \pm i \langle [A, B] \rangle + |\langle \psi | A \pm iB | \psi^\perp \rangle|^2$$



Green: Heisenb-Robertson UR

First bound:

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Green: Heisenberg-Robertson UR

Red: new part, always $\neq 0$ except if $|\psi\rangle$ is joint eigenstate of A and B

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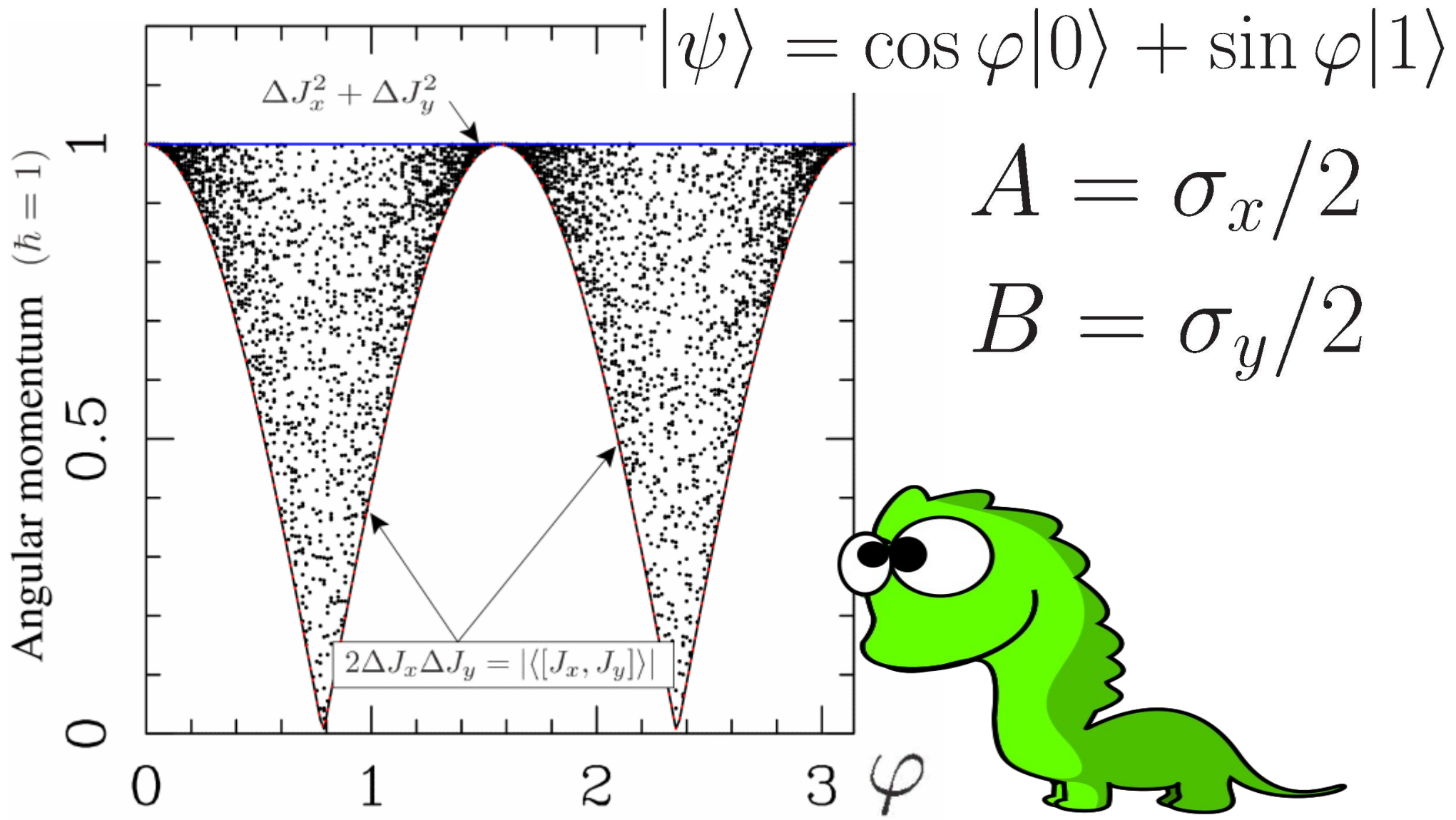
Red: new part, always $\neq 0$ except if $|\psi\rangle$ is joint eigenstate of A and B

Minimizing over $|\psi^\perp\rangle \rightarrow$ only HR (red)

Maximizing over $|\psi^\perp\rangle \rightarrow$ ineq becomes eq

First bound: numerical test for random $|\psi^\perp\rangle$

$$\Delta A^2 + \Delta B^2 \geq \pm i \langle [A, B] \rangle + |\langle \psi | A \pm iB | \psi^\perp \rangle|^2$$



First bound: proof

$$\Delta A^2 + \Delta B^2 \geq \pm i \langle [A, B] \rangle + |\langle \psi | A \pm iB | \psi^\perp \rangle|^2$$

application of Schwartz inequality
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$$C \equiv A - \langle A \rangle; \quad D \equiv B - \langle B \rangle$$

$$\|(C \mp iD)|\psi\rangle\|^2 = \Delta A^2 + \Delta B^2 \mp i \langle [A, B] \rangle$$

$$\begin{aligned} |\langle \psi | (A \pm iB) | \psi^\perp \rangle|^2 &= |\langle \psi | A \pm iB - \langle A \pm iB \rangle | \psi^\perp \rangle|^2 \\ &= |\langle \psi | C \pm iD | \psi^\perp \rangle|^2 \leq \|(C \mp iD)|\psi\rangle\|^2 \end{aligned}$$

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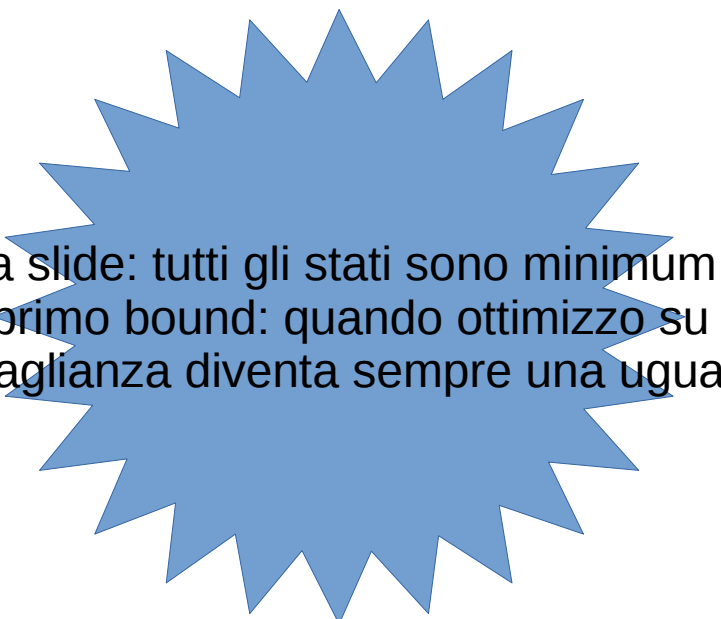
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proof by an anonymous referee

→ Masanao Ozawa
(original proof more complicated)

First bound: minimum uncert. st.



togliere questa slide: tutti gli stati sono minimum uncertainty states per il primo bound: quando ottimizzo su $|\psi^{\text{perp}}\rangle$, la disuguaglianza diventa sempre una uguaglianza!

First bound: minimum uncert. st.

Harmonic osc. $X = \frac{1}{\sqrt{2}}(a + a^\dagger)$, $P = \frac{i}{\sqrt{2}}(a^\dagger - a)$

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• Fock states $|n\rangle$

$$\Delta X^2 + \Delta P^2 = (2n + 1) \equiv \mathcal{L}$$

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MUS 

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• Fock states $|n\rangle$

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• Coherent states $|\alpha\rangle$

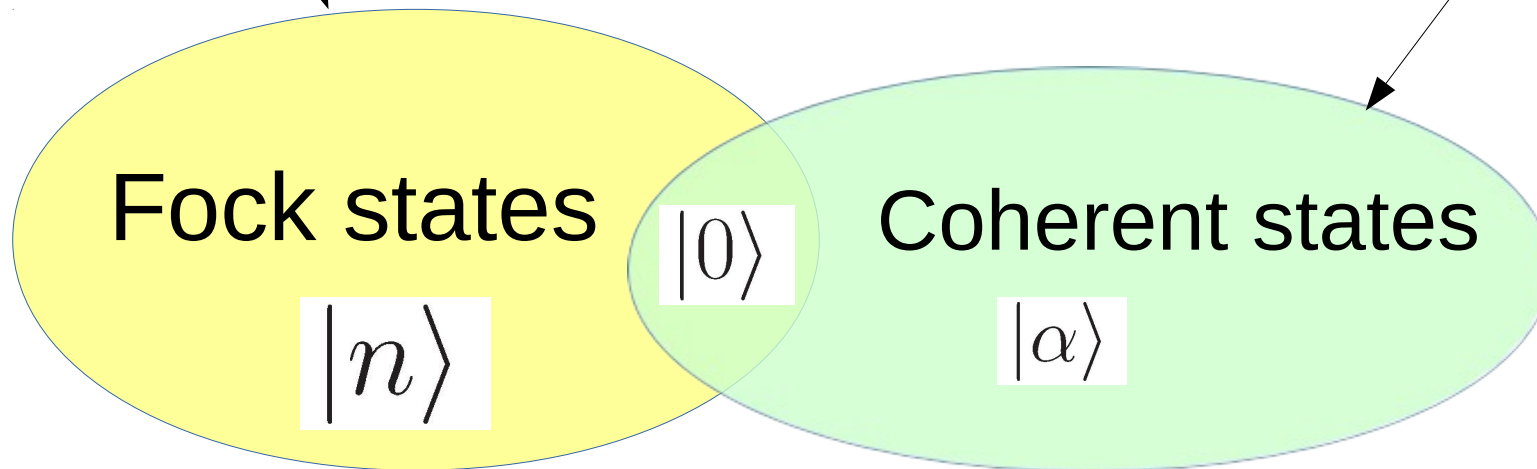
$$\Delta X^2 + \Delta P^2 = 1 \equiv \mathcal{L}$$

MUS



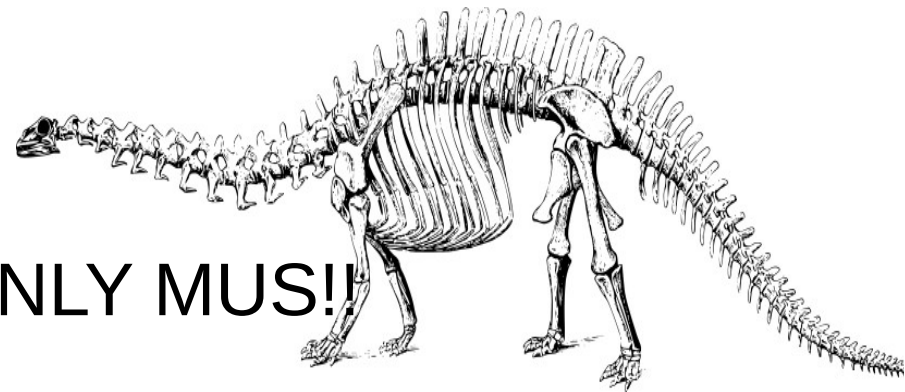
HR Non-MUS

HR MUS



They're all MUS (for **first**, but not **second** bound)

These are just examples, not the **ONLY** MUS!!



Second bound:

$$\Delta A^2 + \Delta B^2 \geq \frac{1}{2} \Delta S^2$$

$$S = A + B$$



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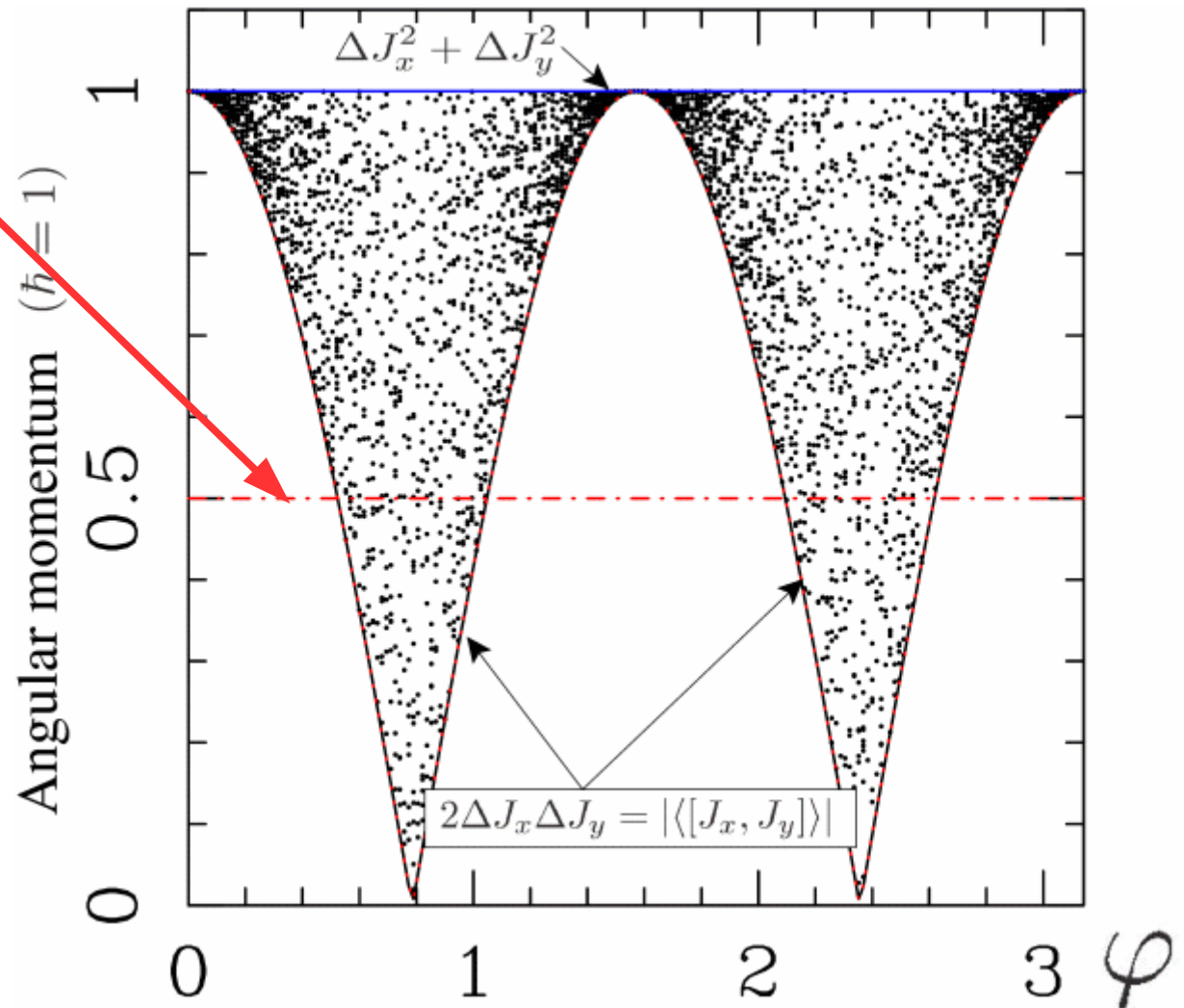
$$\Delta S^2 = \langle (A + B)^2 \rangle - \langle A + B \rangle^2$$

this bound is $\neq 0$ if $|\psi\rangle$
is not an eigenstate of
 $A+B$



Second bound: example

$$\Delta A^2 + \Delta B^2 \geq \frac{1}{2} \Delta S^2$$



Second bound: proof

$$\Delta A^2 + \Delta B^2 \geq \frac{1}{2} \Delta S^2$$

application of the parallelogram ineq.

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application of the parallelogram ineq.

$$C \equiv A - \langle A \rangle; \quad D \equiv B - \langle B \rangle$$

$$2\Delta A^2 + 2\Delta B^2 = \|(C + D)|\psi\rangle\|^2 + \|(C - D)|\psi\rangle\|^2$$

$$\Delta(A + B) = \|(C + D)|\psi\rangle\|, \quad \Delta(A - B) = \|(C - D)|\psi\rangle\|$$

$$\begin{aligned} \Delta A^2 + \Delta B^2 &= \frac{1}{2} [\Delta(A + B)^2 + \Delta(A - B)^2] \\ &\geq \frac{1}{2} \Delta(A + B)^2, \end{aligned}$$

State dependence.

$$\Delta A^2 + \Delta B^2 \geq \max(\mathcal{L}, \mathcal{M})$$

both our bounds are state-dependent



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$$\mathcal{L} = \mathcal{L}(|\psi\rangle) \quad \mathcal{M} = \mathcal{M}(|\psi\rangle)$$



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? the second becomes trivial
minimizing over $|\psi\rangle$ (just take an eigenstate of $A+B$)

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$$\Delta A^2 + \Delta B^2 \geq \frac{1}{2} \Delta S^2$$

$$\min_{|\psi\rangle} \max_{|\psi^\perp(\psi)\rangle} \pm i \langle [A, B] \rangle + |\langle \psi | A \pm iB | \psi^\perp \rangle|^2$$



Using the same techniques, we can
tighten the Heis-Rob bound!

$$\Delta A \Delta B \geq \pm \frac{i}{2} \langle [A, B] \rangle$$





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tighten the Heis-Rob bound!

$$\Delta A \Delta B \geq \pm \frac{i}{2} \langle [A, B] \rangle / \left(1 - \frac{1}{2} \left| \langle \psi | \frac{A}{\Delta A} \pm i \frac{B}{\Delta B} | \psi^\perp \rangle \right|^2 \right)$$



Heisenberg uncertainty
relation vs. principle

two very different notions!!!!!!!!!!!!!!!!!!!!

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→ **WRONG** interpretation!

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u relation: refers only to the **preparation** of the state. Can't **prepare** a state with sharp values for incompatible observables.

Heisenberg uncertainty

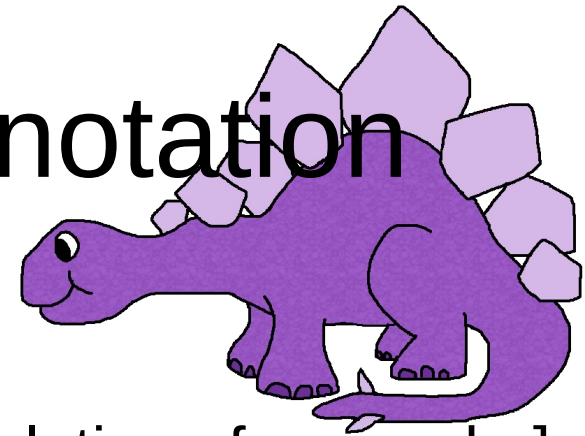
In this talk: unc
RELATIONS!

WRONG interpretation!

u relation: refers only to the **preparation** of the state. Can't **prepare** a state with sharp values for incompatible observables.

Asher Peres introduced this notation

uncertainty relation:



“The only correct interpretation of [the uncertainty relations for x and p] is the following: If the **same preparation** procedure is repeated many times, and is followed either by a measurement of x , or by a measurement of p , the various results obtained for x and for p have standard deviations, Δx and Δp , whose product cannot be less than $\hbar/2$. **There never is any question here that a measurement of x ‘disturbs’ the value of p and vice versa, as sometimes claimed.** These measurements are indeed incompatible, but they are performed on different particles (all of which were identically prepared) and therefore these measurements cannot disturb each other in any way. The uncertainty relation [...] only reflects the intrinsic randomness of the outcomes of quantum tests.”

(he didn't want to talk about unc principle)

try looking up “uncertain principle” in his book.

stronger uncertainty relations for
variances with nontrivial lower bound.

$$\Delta A^2 + \Delta B^2 \geq \pm i \langle [A, B] \rangle + |\langle \psi | A \pm iB | \psi^\perp \rangle|^2$$

$$\Delta A^2 + \Delta B^2 \geq \frac{1}{2} \Delta S^2$$

$$S = A + B$$



PRL 113 260401

Entanglement and Complementarity



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Chiara

Macchiavello

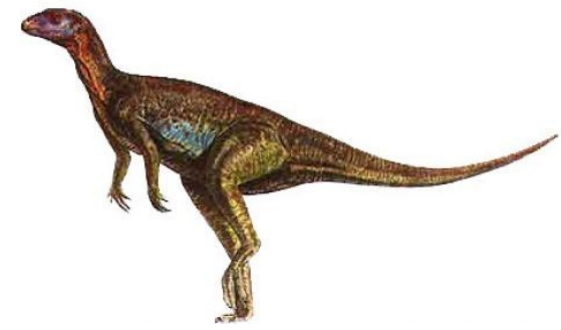
Dagmar Bruss

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What I'm going to talk about

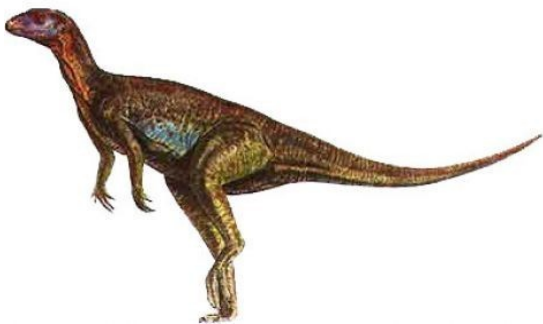
We always say that entangled states are more correlated... WHAT DOES IT MEAN exactly?



What I'm going to talk about

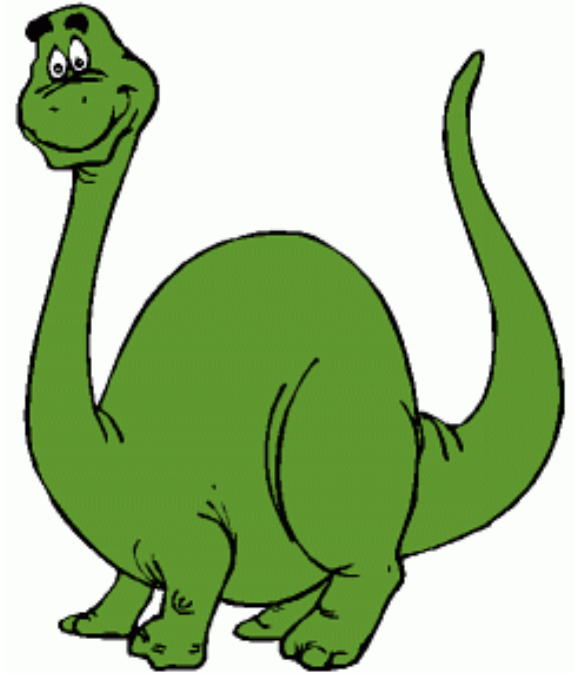
We always say that entangled states are more correlated... WHAT DOES IT MEAN exactly?

they have more correlations
among complementary
observables than separable ones



Usual approaches to study entanglement

- Non locality
- Negative partial transpose
- Bell inequality violations
- Enhanced precision in measurements
- etc.



Here: we use correlations
among two (or more)
**COMPLEMENTARY
PROPERTIES**

different way to think about
entanglement, as
correlations among
complementary properties



Remember: Complementary properties.



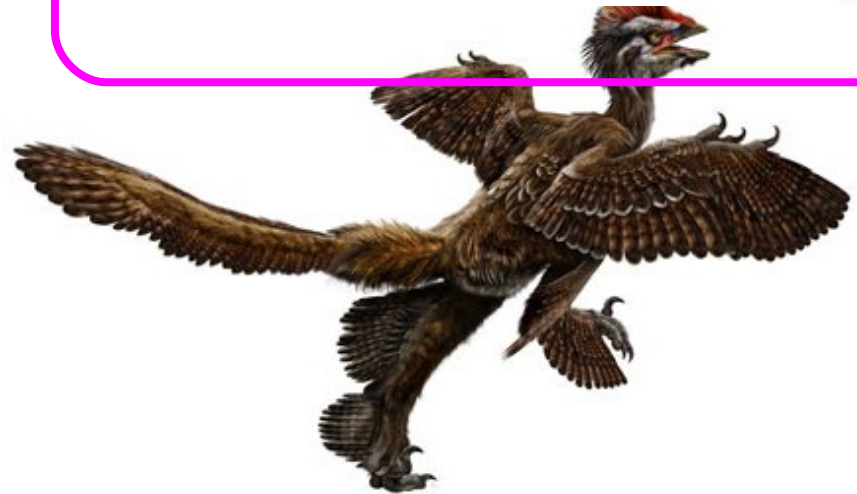
Remember: Complementary properties.

Two observables: the knowledge of one gives no knowledge of the other

$$A = \sum_a f(a) |a\rangle \langle a|$$

$$C = \sum_c g(c) |c\rangle \langle c|$$

$$|\langle a|c\rangle|^2 = \frac{1}{d}$$



simplest example:



simplest example:



$$\frac{|00\rangle + |11\rangle}{\sqrt{2}} = \frac{|++\rangle + |--\rangle}{\sqrt{2}}$$

Maximally entangled state: perfect correlation BOTH on 0/1 and on +/-

$$|\pm\rangle \equiv \frac{|0\rangle \pm |1\rangle}{\sqrt{2}}$$

simplest example:



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Maximally entangled state: perfect correlation BOTH on 0/1 and on +/-

$$(|00\rangle\langle 00| + |11\rangle\langle 11|)/2$$

simplest example:



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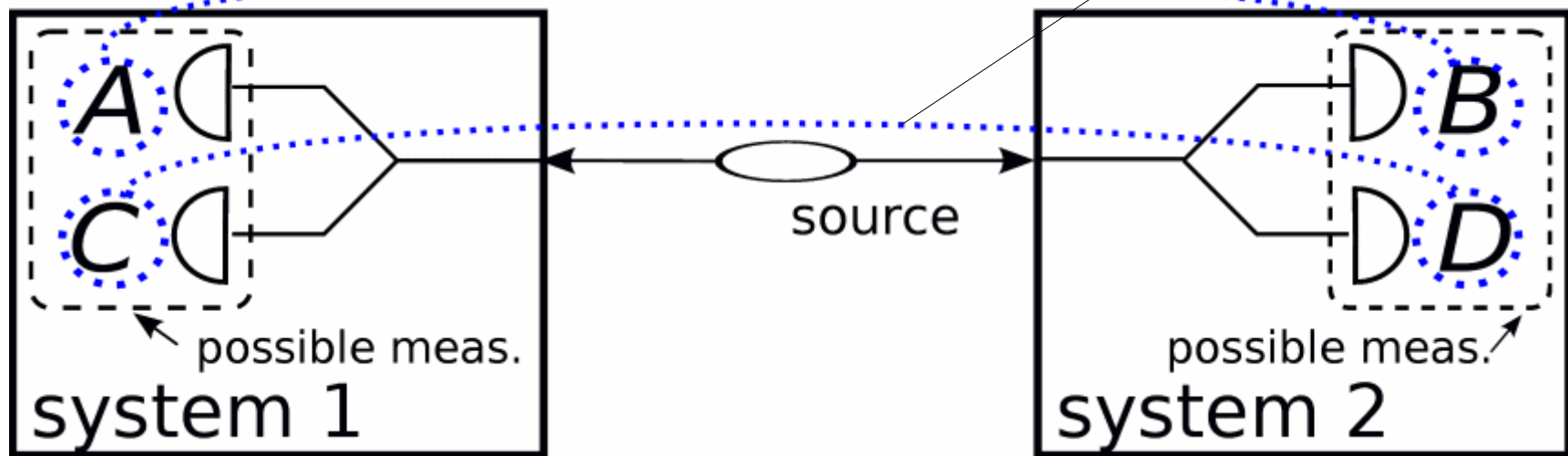
Maximally entangled state: perfect correlation BOTH on 0/1 and on +/-

$$(|00\rangle\langle 00| + |11\rangle\langle 11|)/2 = (|+\rangle\langle +| + |-\rangle\langle -|)/2 \otimes (|+\rangle\langle +| + |-\rangle\langle -|)/2$$

separable state: perfect correlation for 0/1, no correlation for +/-

Simple experiment

- On system 1 measure either A or C
- On system 2 measure either B or D
- Calculate correlations $A-B$ and $C-D$



How to measure correlation?



How to measure correlation?

- Mutual information

$$I_{AB} = H(A) + H(B) - H(A, B)$$



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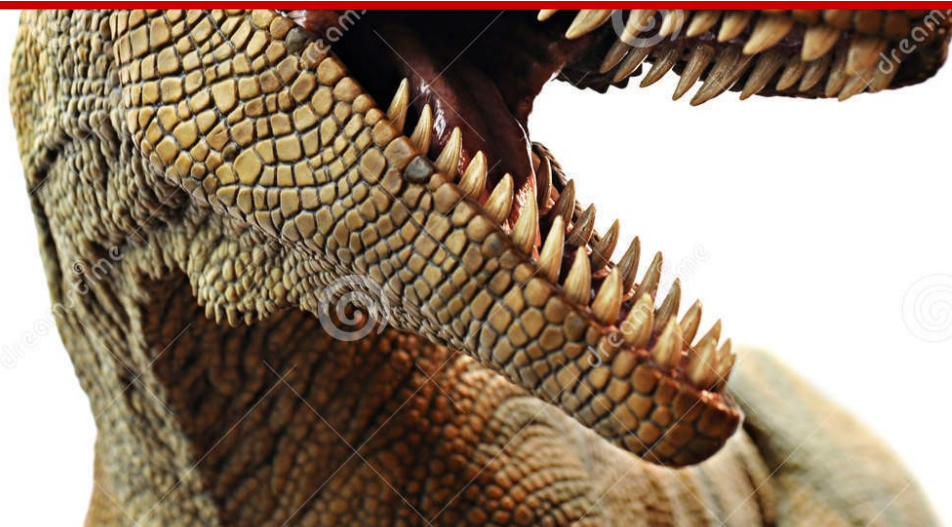
- Mutual information

$$I_{AB} = H(A) + H(B) - H(A, B)$$

- Pearson correlation coefficient

$$C_{AB} \equiv \frac{\langle AB \rangle - \langle A \rangle \langle B \rangle}{\sigma_A \sigma_B}$$

$|C_{AB}| = 1 \Rightarrow$
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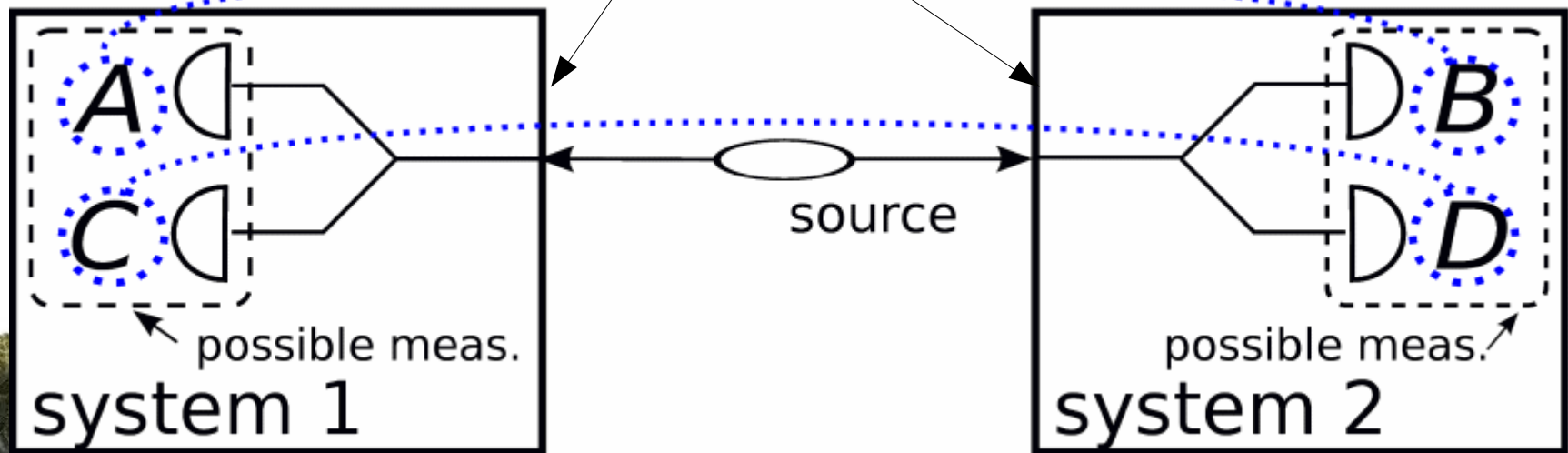


Use these to measure correlations among

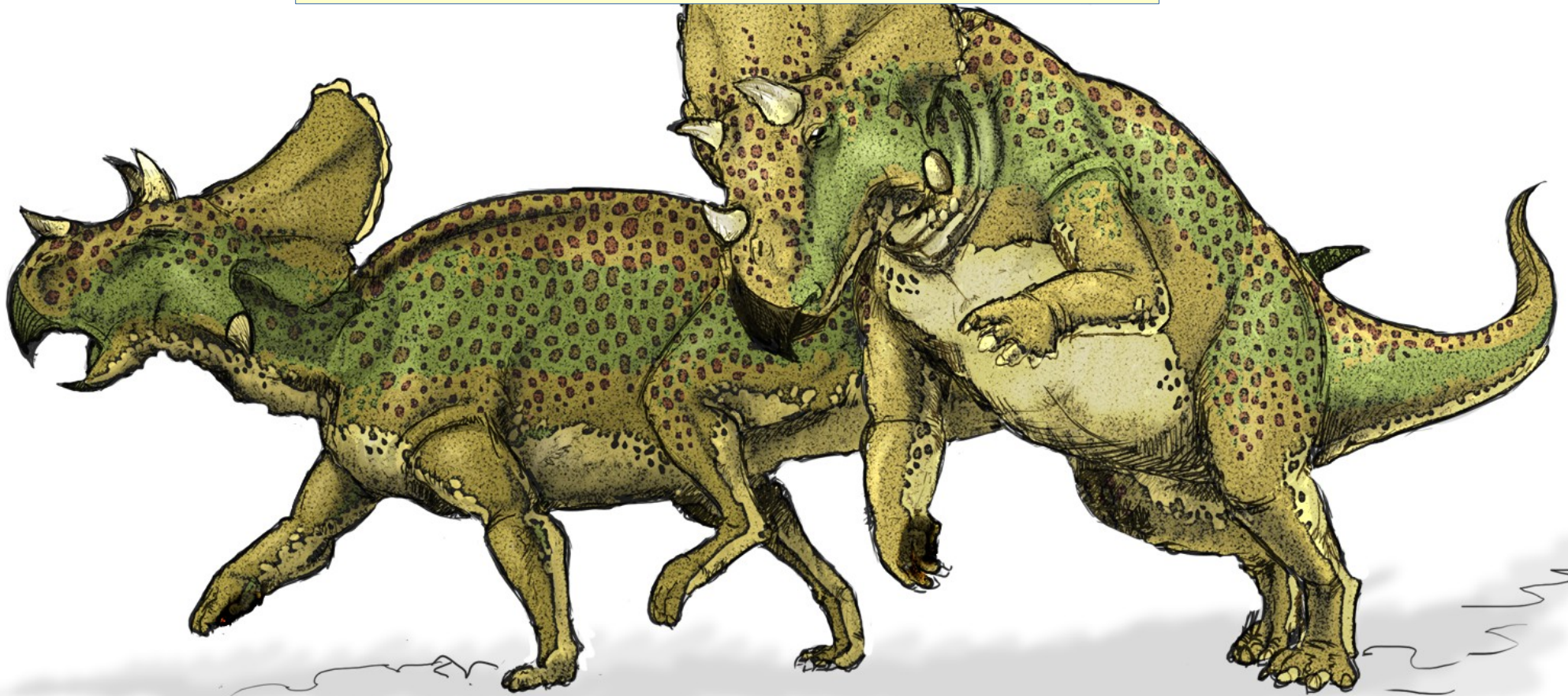
2 complementary properties

$$A \otimes B \longleftrightarrow \text{complement to} \quad C \otimes D$$

of 2 systems



Some results...



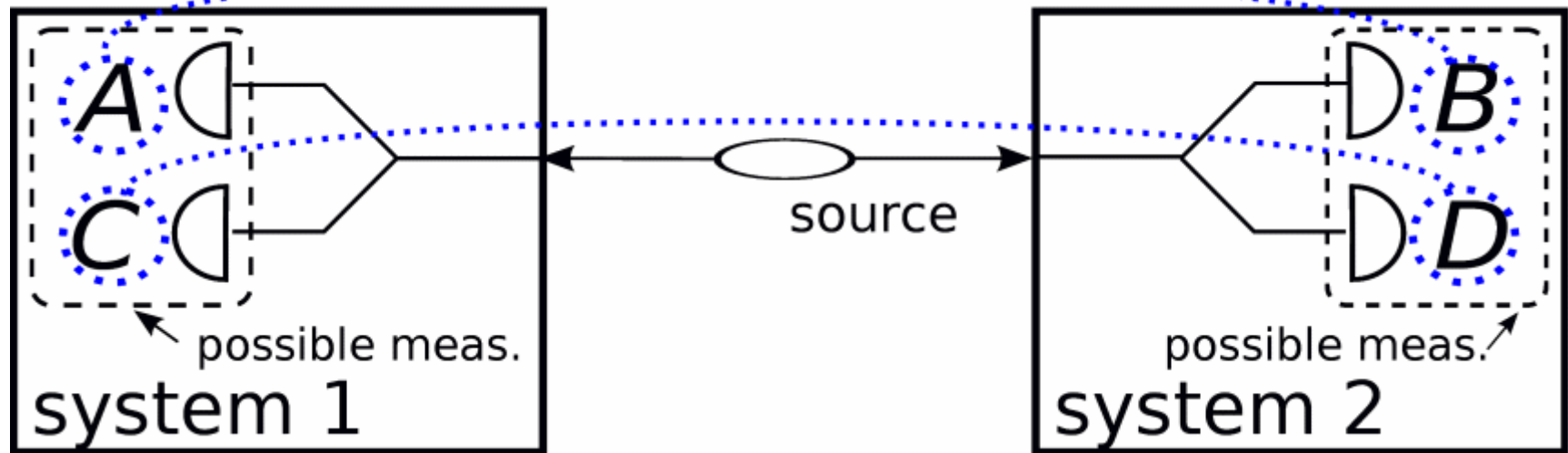
Start with mutual information

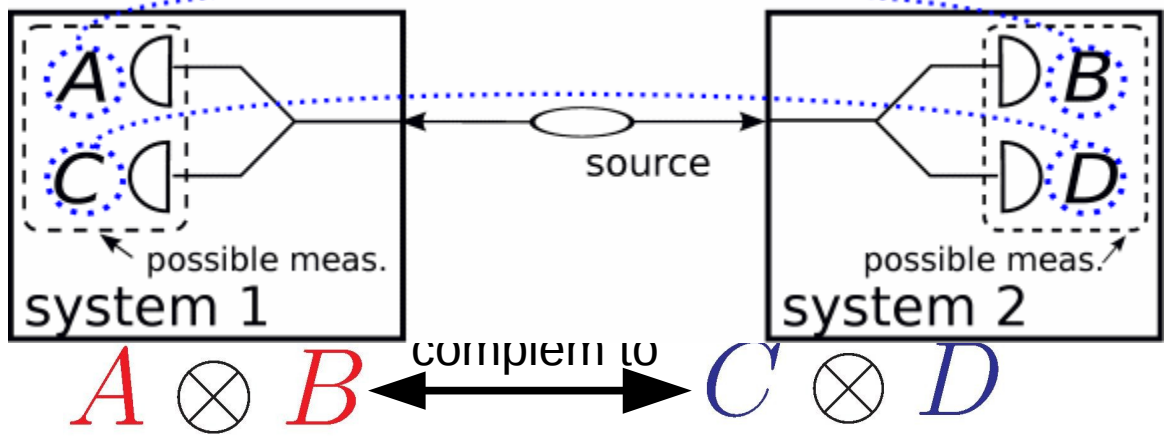
$$I_{AB} = H(A) + H(B) - H(A, B)$$



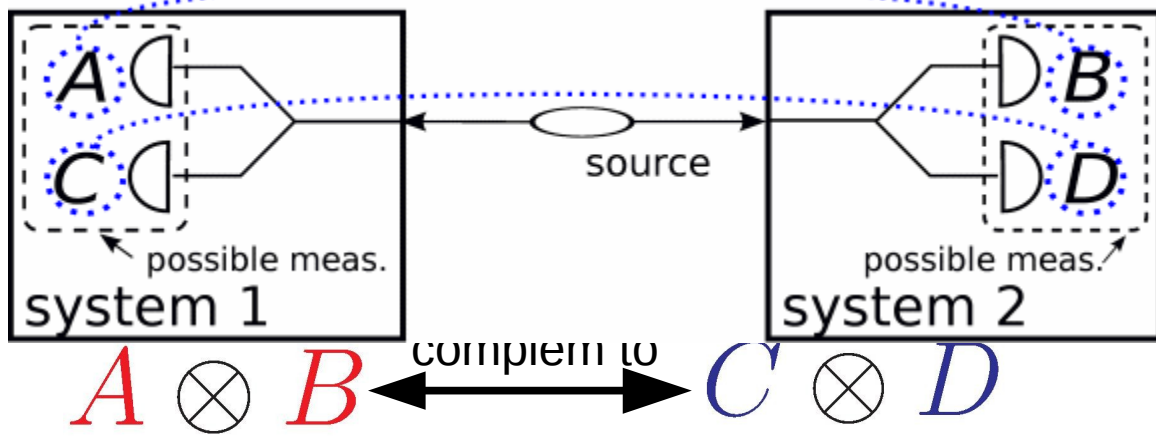
“total” correlation given by the sum

$$I_{AB} + I_{CD}$$





The system state is **maximally entangled** iff perfect correlation on **both $A-B$ and $C-D$**

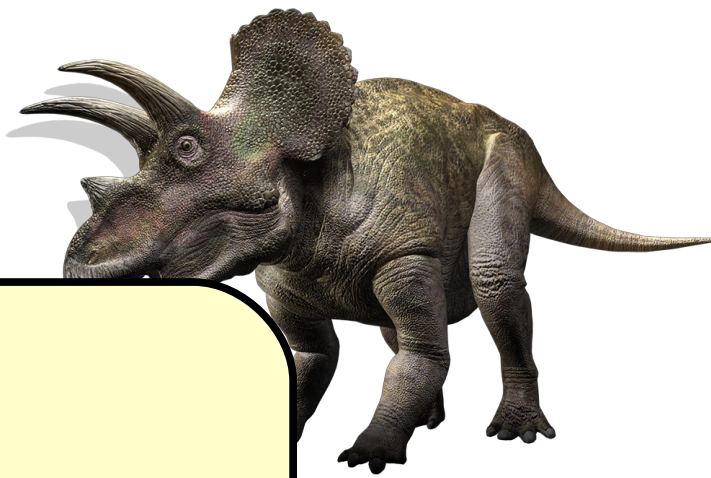
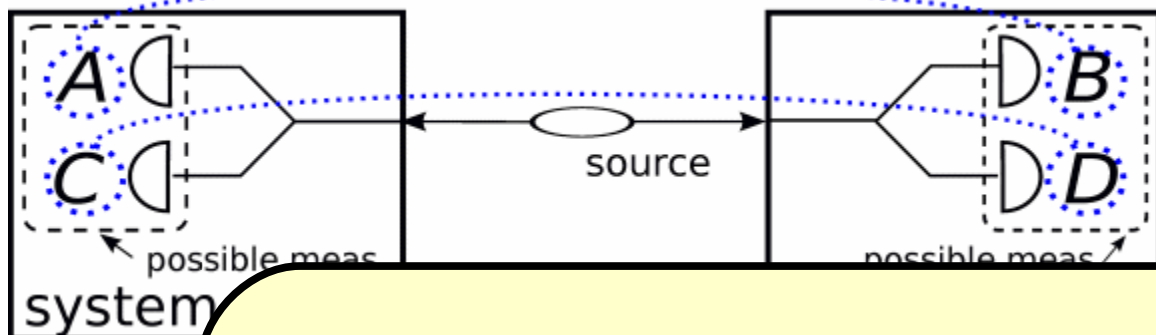


The system state is **maximally entangled** iff perfect correlation on **both $A-B$ and $C-D$**

$$I_{AB} + I_{CD} = 2 \log d$$

(for some observ $ABCD$)

$$\Leftrightarrow |\Psi_{12}\rangle \text{ maximally entangled}$$



$$I_{AB} \leq \log_2 d$$

$$I_{CD} \leq \log_2 d$$

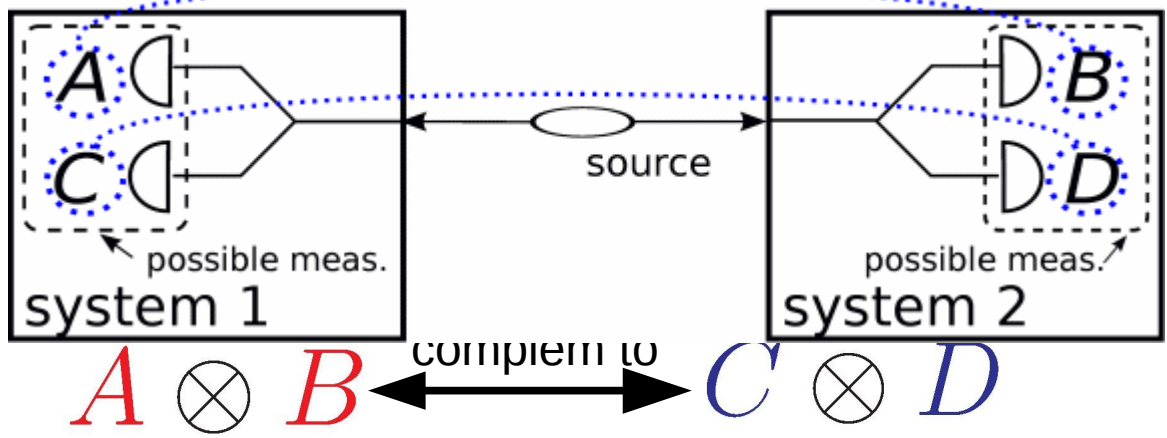
on both A-B and C-D

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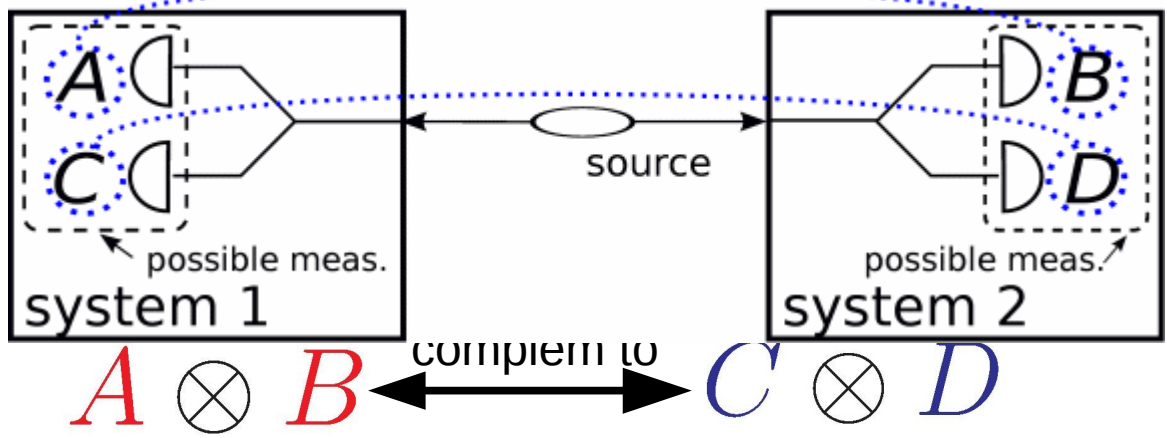
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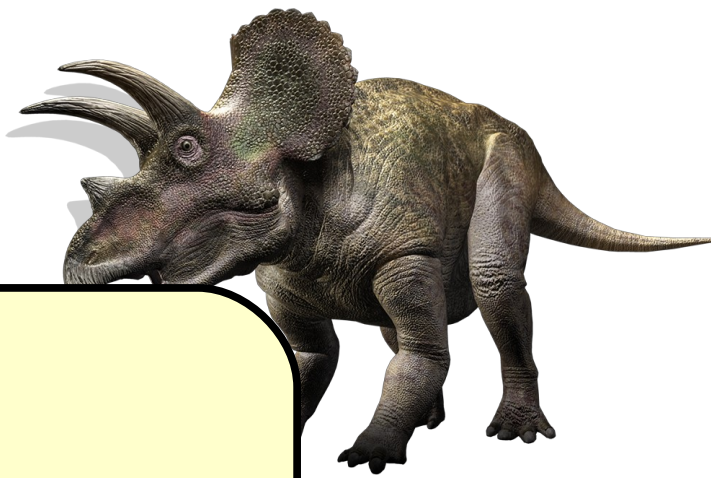
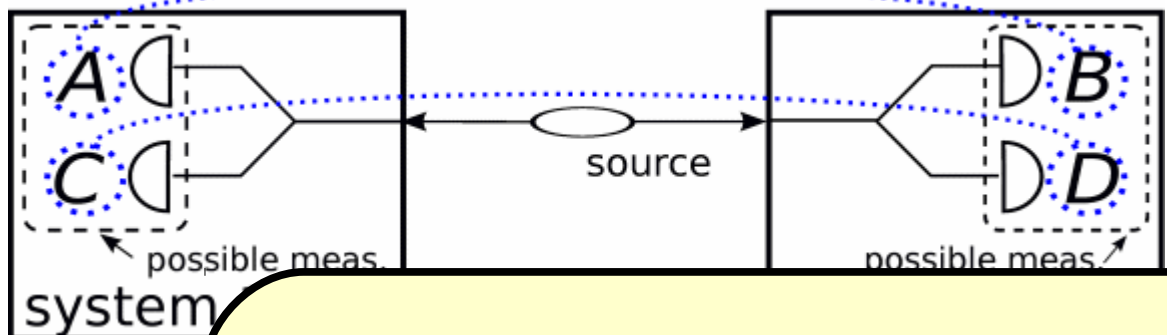


The system state is **entangled** if correlations on **both** $A-B$ and $C-D$ are large enough



The system state is **entangled** if correlations on **both** A - B and C - D are large enough

$$I_{AB} + I_{CD} > \log d \implies \rho_{12} \text{ ent}$$



A

$$I_{AB} \leq \log_2 d$$

$$I_{CD} \leq \log_2 d$$

are large enough

led if
d C-D

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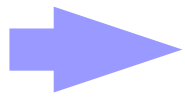
Can the bound be made **tighter**?



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NO!!

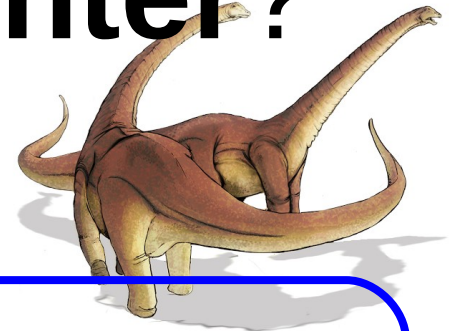


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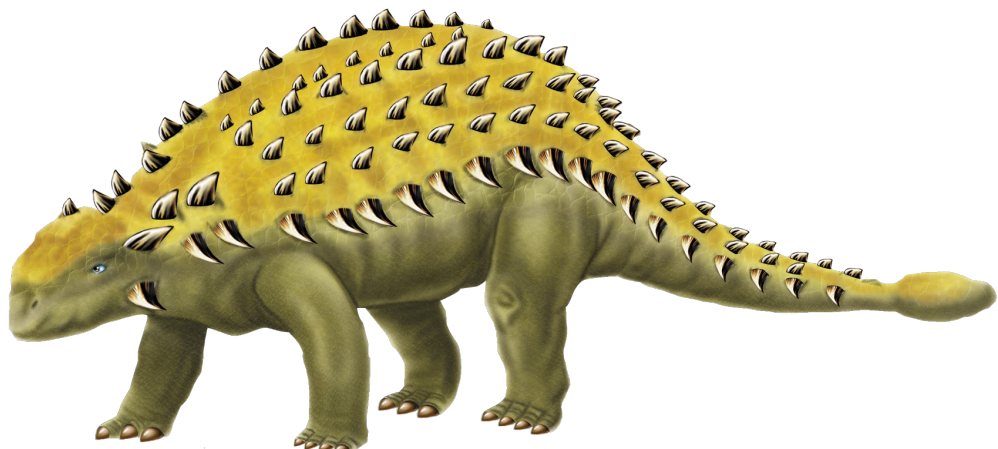
the **separable** state

$$\frac{1}{2} (|00\rangle\langle 00| + |11\rangle\langle 11|)$$

saturates it: $I_{AB} + I_{CD} = \log d$

The system state is **entangled** if correlations on **both** A - B and C - D are large enough

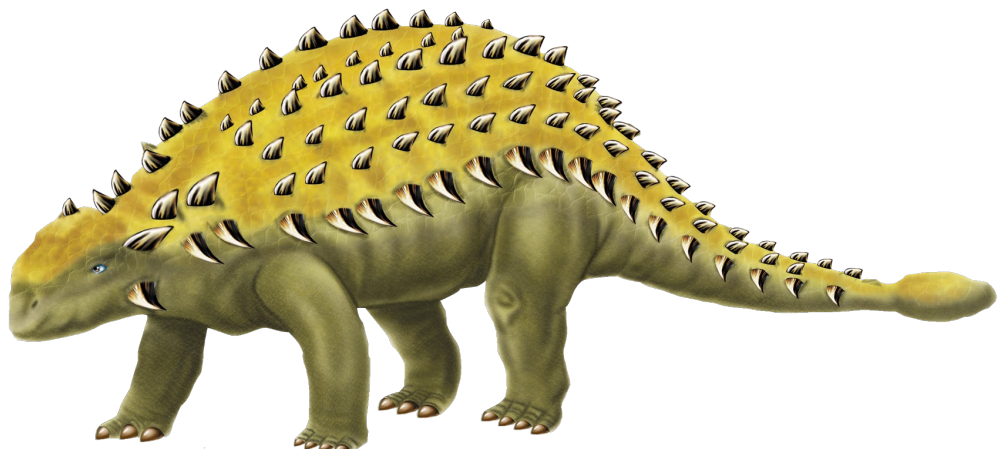
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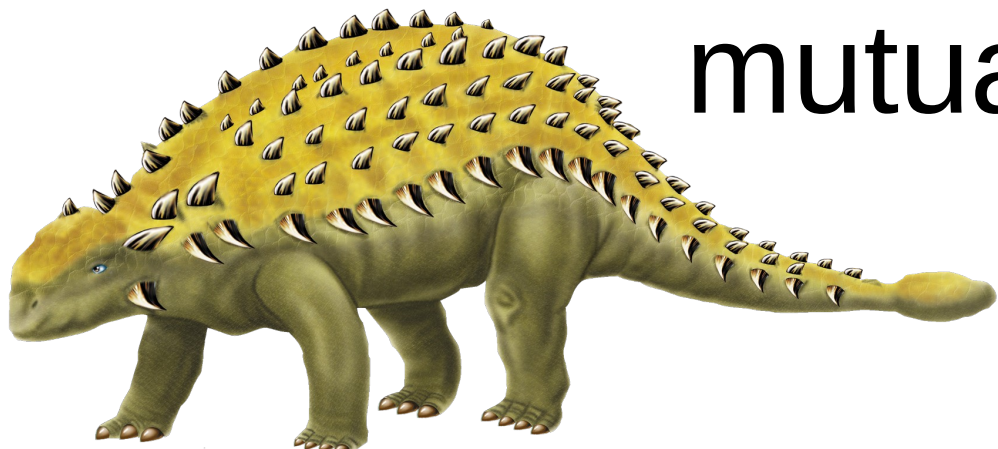
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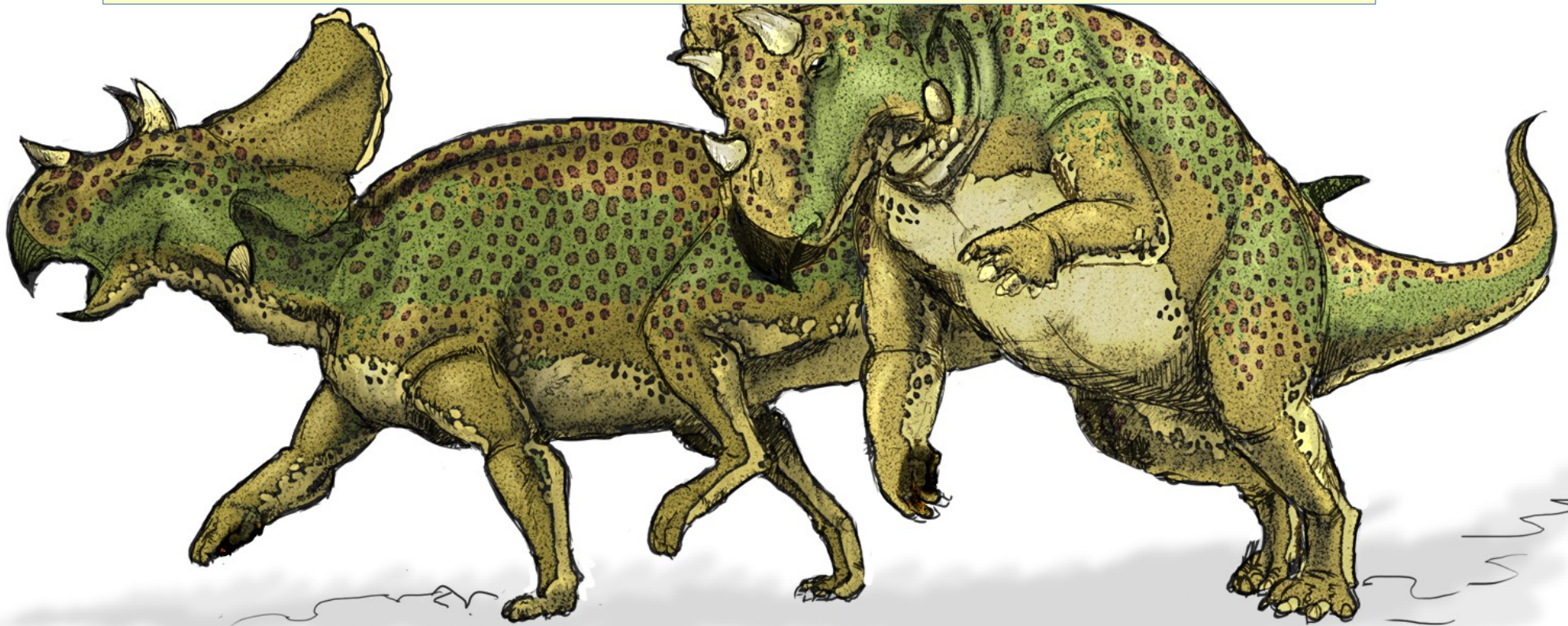
➔ **NO!!**

$$|\psi_\epsilon\rangle = \epsilon|00\rangle + \sqrt{1 - \epsilon^2}|11\rangle$$

is entangled but has negligible mutual info for $\epsilon \rightarrow 0$



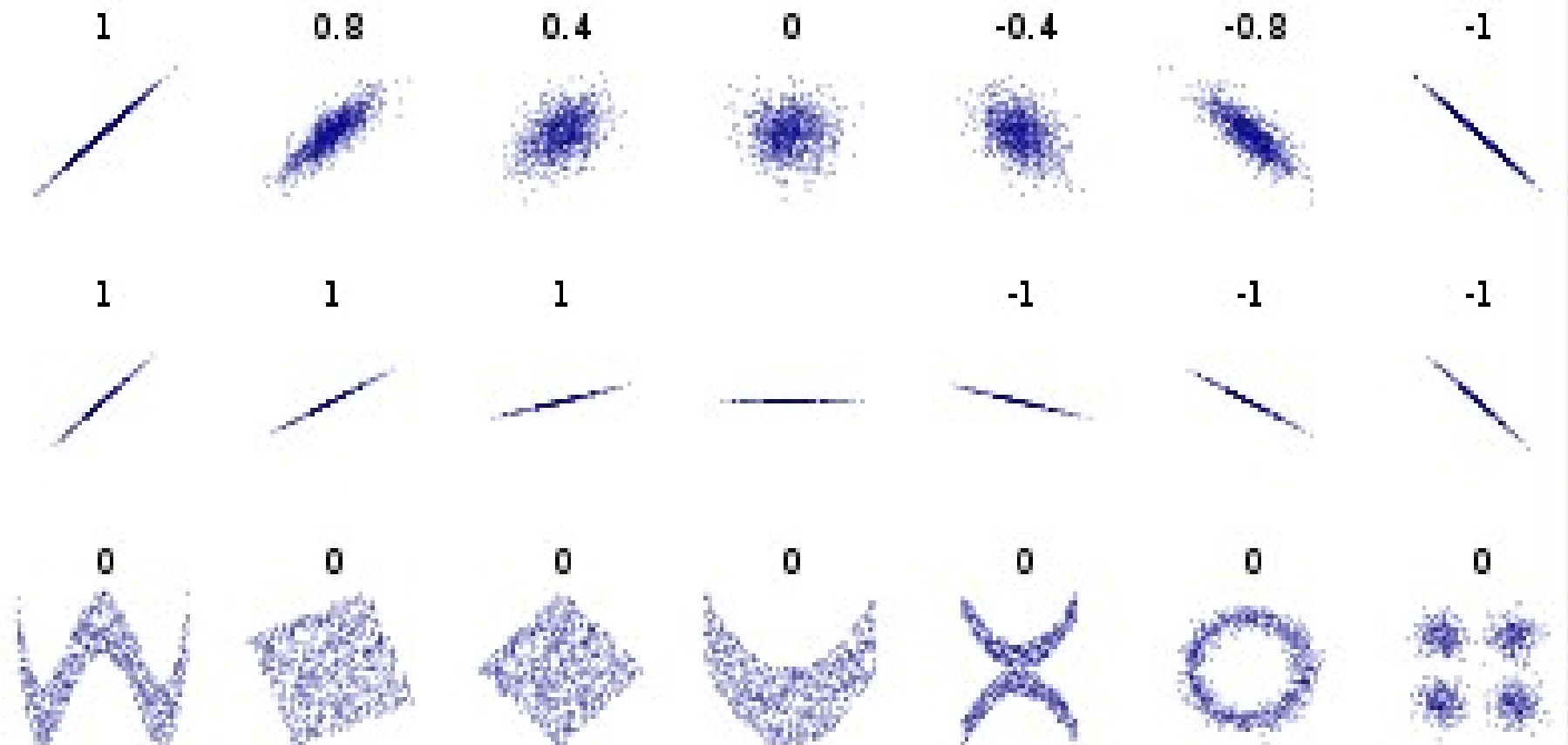
Another measure of
correlation...



Pearson correlation coefficient

$$\mathcal{C}_{AB} \equiv \frac{\langle AB \rangle - \langle A \rangle \langle B \rangle}{\sigma_A \sigma_B}$$

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... but its modulus is still $\leq |1|$:

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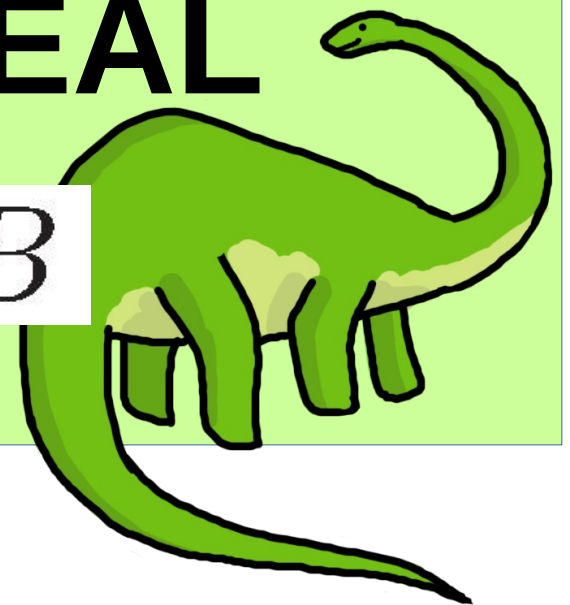
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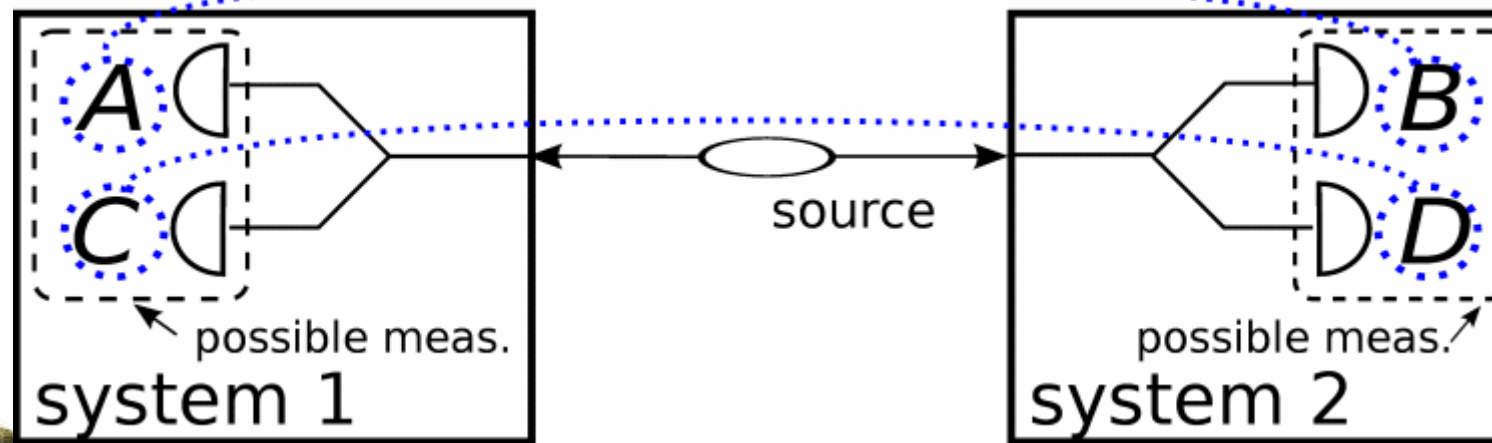
not a problem for us: A and
 B commute, so it's **REAL**

$$A \otimes B = A \otimes \mathbb{1} + \mathbb{1} \otimes B$$



Total correlation: again use the sum

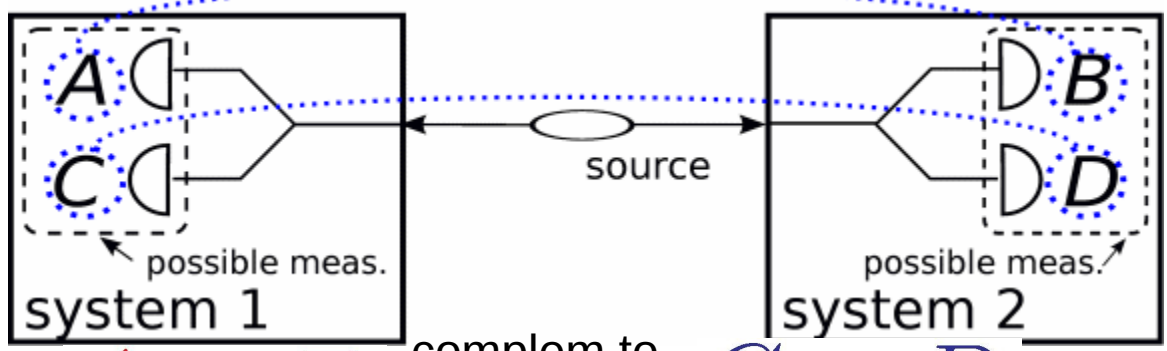
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complem to

$$A \otimes B \longleftrightarrow C \otimes D$$



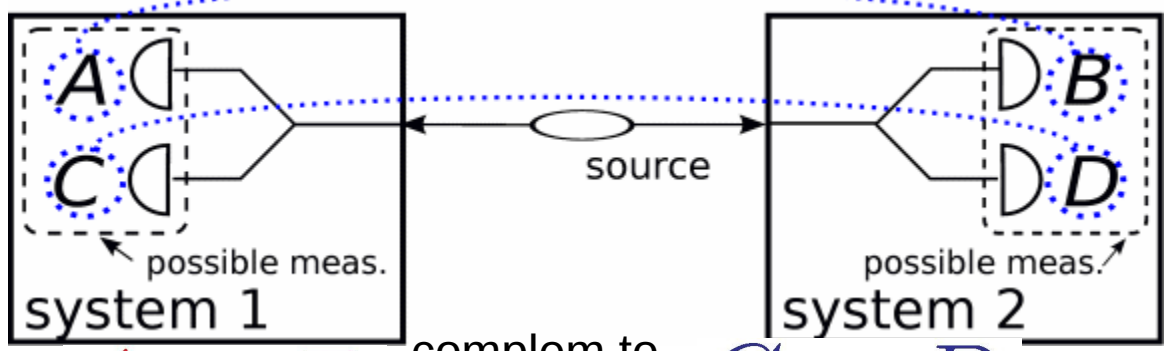


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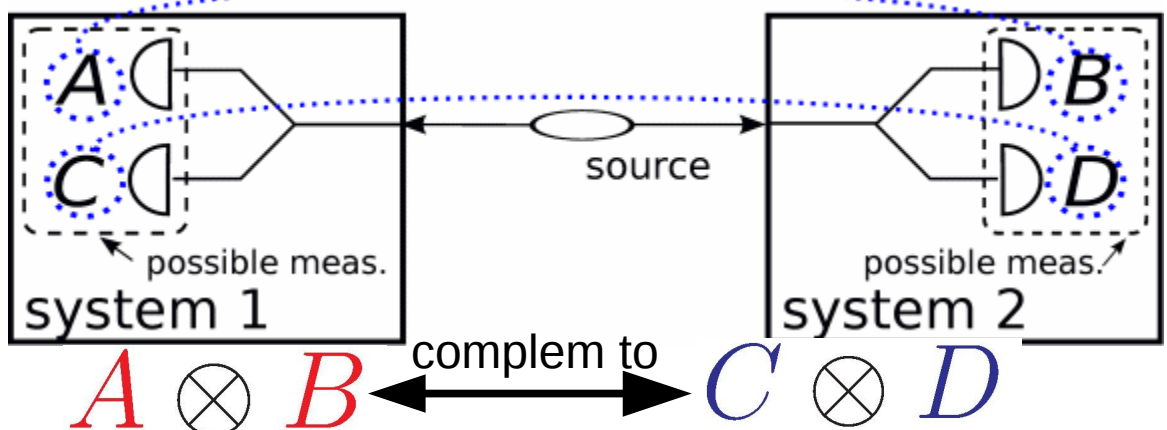


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True also using Pearson! (for linear observables: Pearson measures only linear correl)

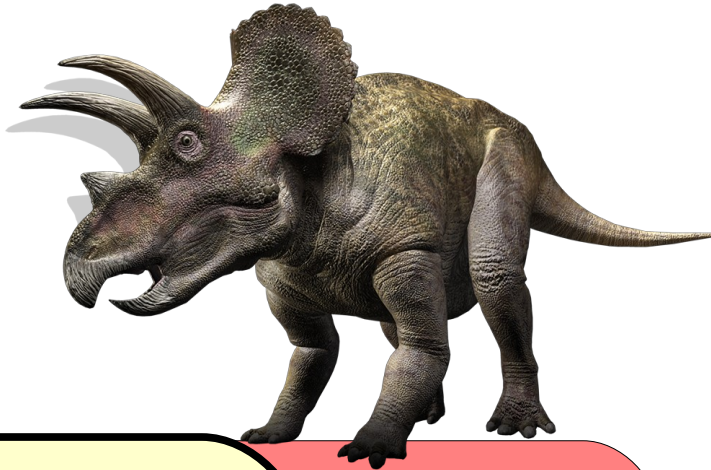
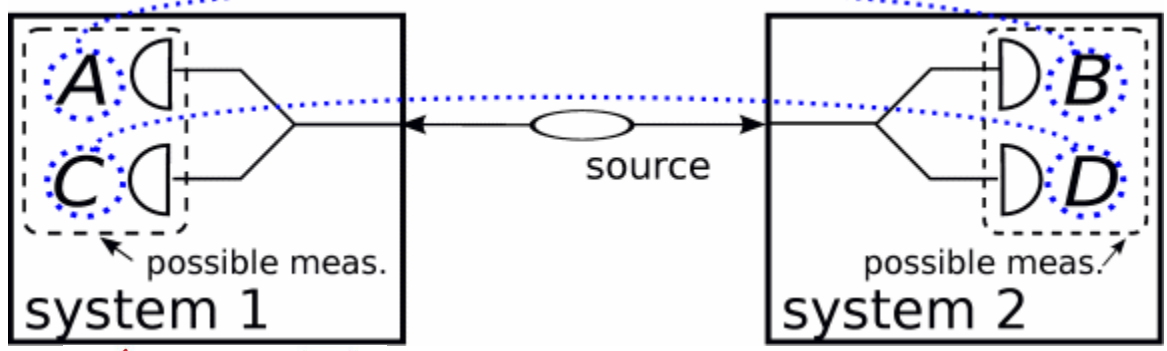


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$$\Leftrightarrow |\Psi_{12}\rangle \quad \text{maximally entangled}$$



$A \otimes B$

$|\mathcal{C}_{AB}| \leq 1$

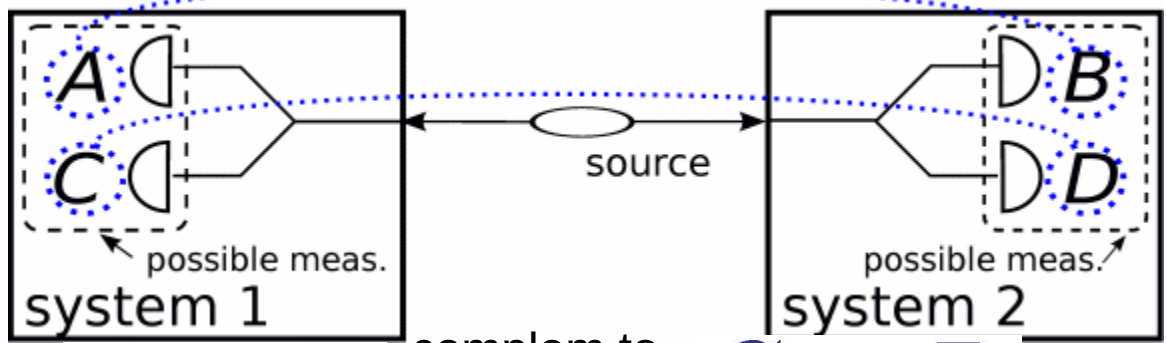
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True linear observables: Pearson measures only linear correlation

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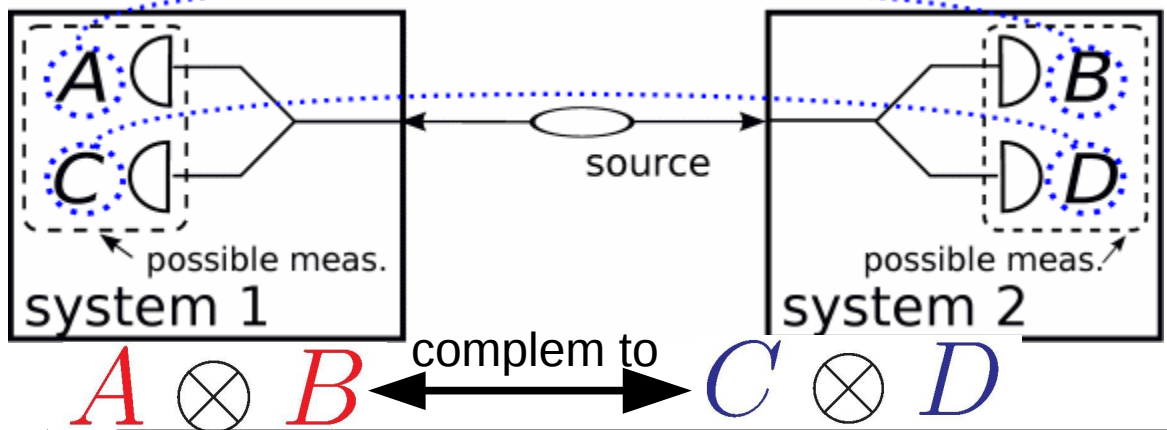
$\Leftrightarrow |\Psi_{12}\rangle$ maximally entangled



$A \otimes B$ \longleftrightarrow complem to $C \otimes D$

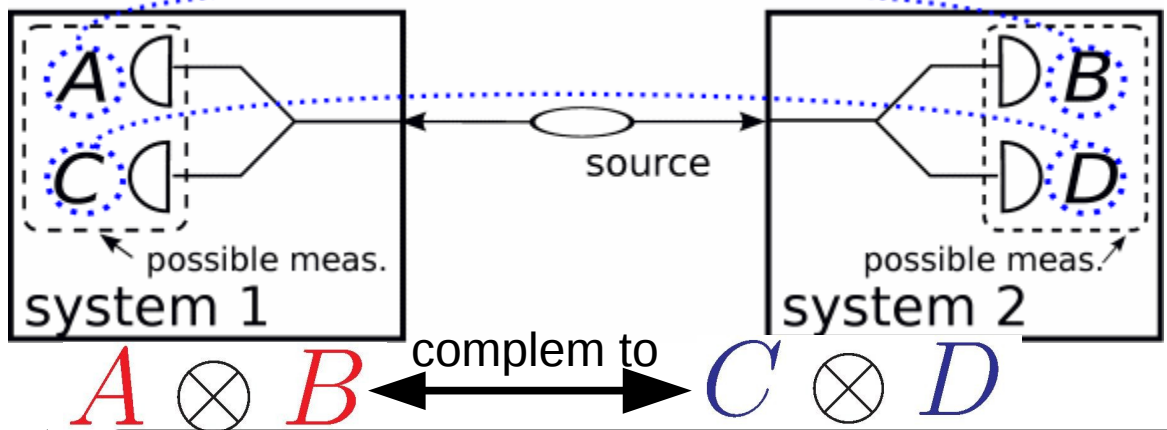


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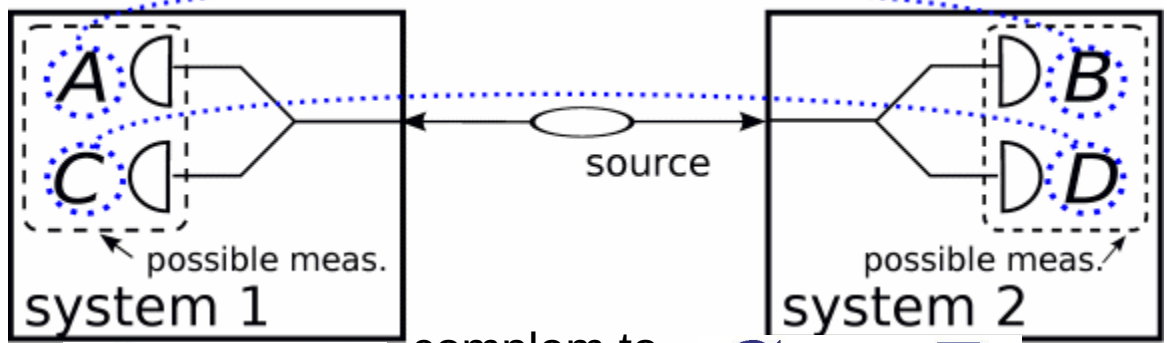


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T
led if

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CO
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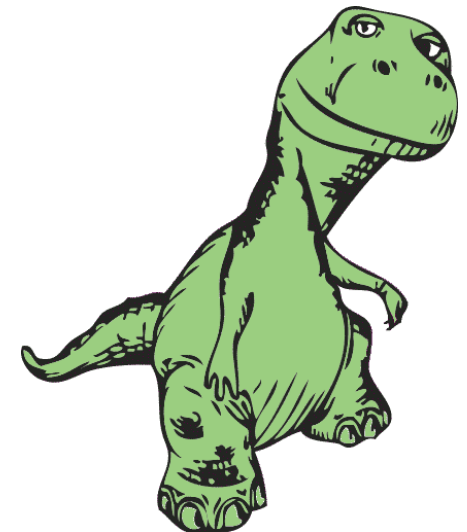
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$$|\mathcal{C}_{AB}| + |\mathcal{C}_{CD}| = 1$$

(perfect correl on one basis,
no correl on the complem)

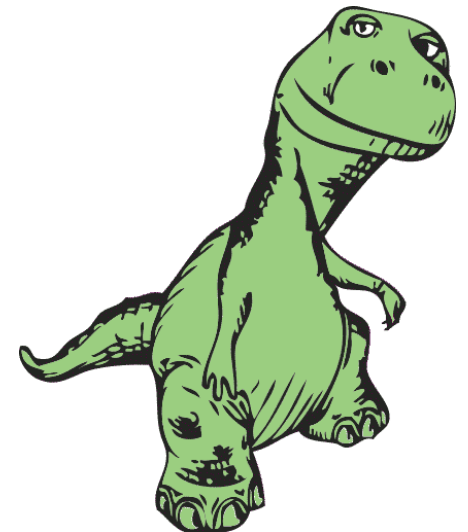


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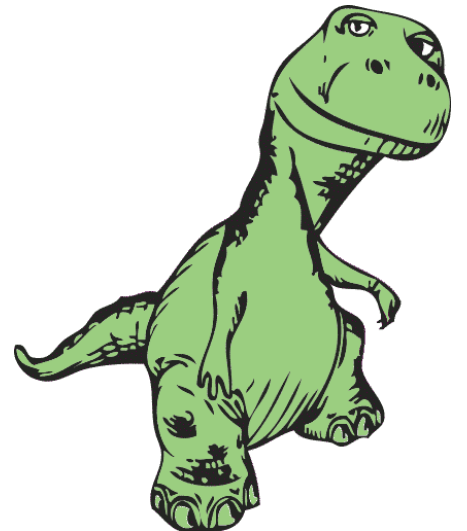


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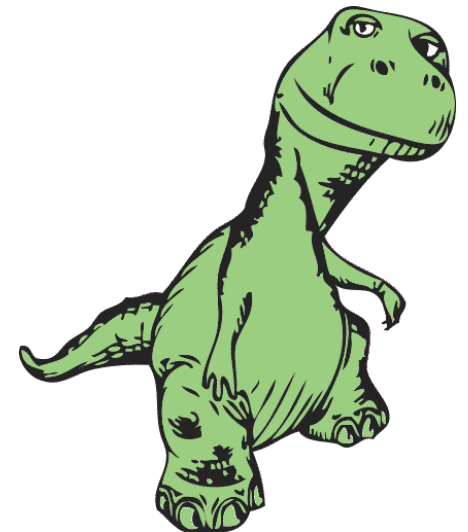
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Has negligible mutual info for $\epsilon \rightarrow 0$

but Pearson correlation
always >1 !



Simple criterion for entanglement detection!!

Just measure two complementary properties. Are the correlations greater than perfect correlation on one?

⇒ The state is entangled!




Simple to measure and simple to optimize.

Unfortunately: not very effective in finding entanglement in random states

- Entanglement as correlation among complementary observables
- Using different measures of correlation:
 - Mutual info
 - Pearson correlation
- Some theorems and some conjectures



The most correlated states are entangled
but ent states are not the most correlated



Correlations on
complementary prop.
help understanding
entanglement