

# Fault tolerant quantum metrology

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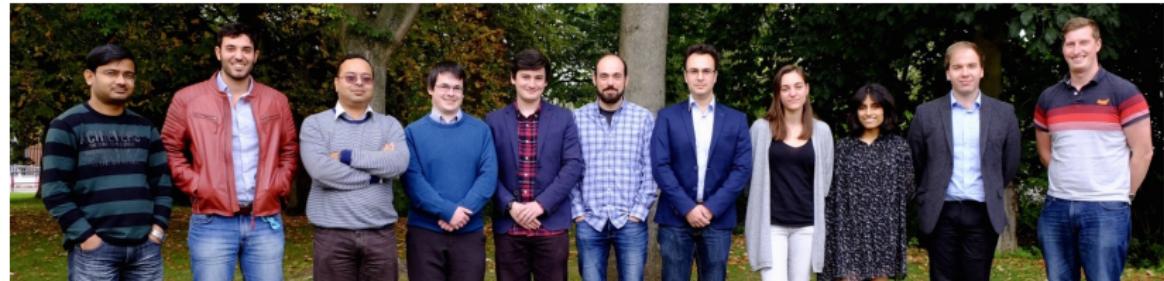


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- Samuele Ferracin
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- Andrew Jackson
- Theodoros Kapourniotis
- Max Marcus
- Francesco Albarelli



# Enhancement using quantum probes

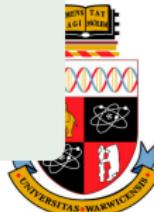
- Probe made of constituents qubits - quantum 2-level systems
- Using a probe  $|\Psi\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}}$ , and

$$R_z(\phi) = e^{-i\phi Z}, \quad Z = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$|\psi\rangle \rightarrow [R_z(\phi)] \rightarrow \text{circuit symbol}$$

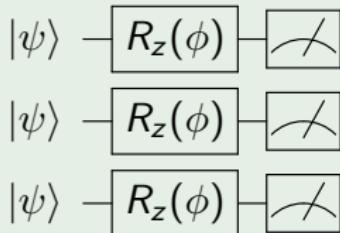
- $|\Psi\rangle = \frac{|0\rangle + e^{-i\phi}|1\rangle}{\sqrt{2}}$
- $M = X$
- $\langle M \rangle \sim \sin^2(\phi/2)$

- $|\Psi\rangle = \frac{|00\rangle + e^{-2i\phi}|11\rangle}{\sqrt{2}}$
- $M = X \otimes X$
- $\langle M \rangle \sim \sin^2 \phi$

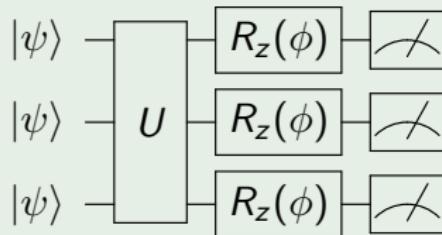


# Enhancement using quantum probes

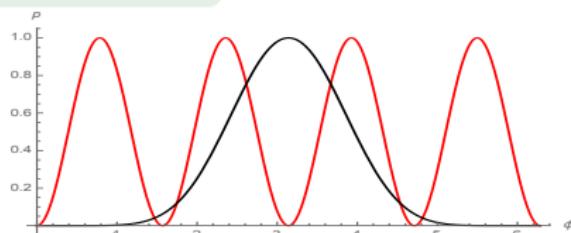
- Probe made of constituents, say  $N$  qubits, in a state  $|\Psi\rangle$
- Field  $R_z(\phi) = e^{-i\phi Z}$
- Measurement  $M$



- $|\Psi\rangle = \left( \frac{|0\rangle + e^{-i\phi}|1\rangle}{\sqrt{2}} \right)^{\otimes N}$
- $\langle M \rangle \sim (\sin^2(\phi/2))^N$



- $|\Psi\rangle = \frac{|00\dots000\rangle + e^{-iN\phi}|11\dots111\rangle}{\sqrt{2}}$
- $\langle M \rangle \sim \sin^2(N\phi/2)$



# Estimation theory

- Precision: variance of a ‘units-corrected’ estimator  $\tilde{\phi}$

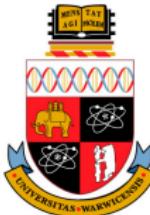
$$\Delta\phi = \left\langle \left( \frac{\tilde{\phi}}{|\partial\tilde{\phi}/\partial\phi|} - \phi \right)^2 \right\rangle^{1/2}$$

- Cramér-Rao: Variance of estimator lower bounded

$$\Delta\phi \geq \frac{1}{\sqrt{\nu F_\phi}} \geq \frac{1}{\sqrt{\nu Q_\phi}},$$

- $F_\phi$  is the classical Fisher information
- $Q_\phi$  is the quantum Fisher information

Helstrom, Holevo, Braunstein, Caves, ...



# Estimation theory

- Precision: variance of a ‘units-corrected’ estimator  $\tilde{\phi}$

$$\Delta\phi = \left\langle \left( \frac{\tilde{\phi}}{|\partial\tilde{\phi}/\partial\phi|} - \phi \right)^2 \right\rangle^{1/2}$$

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$$\Delta\phi \geq \frac{1}{\sqrt{\nu F_\phi}} \geq \frac{1}{\sqrt{\nu Q_\phi}},$$

- $F_\phi$  is the classical Fisher information
- $Q_\phi$  is the quantum Fisher information

Helstrom, Holevo, Braunstein, Caves, ...

- Quantum metrology is mostly about bounds on the variance

We will provide exact variances



# Paedagogical examples

## GHZ (Greenberger-Horne-Zeilinger) states

$$\frac{|00\cdots 00\rangle + |11\cdots 11\rangle}{\sqrt{2}}$$

$$Q_\phi = N^2$$

$$\Delta\phi \sim \frac{1}{N}$$

Bollinger *et al.*, PRA, **54**, R4649, (1996)

## N00N states

$$\frac{|N, 0\rangle + |0, N\rangle}{\sqrt{2}}$$

$$Q_\phi = N^2$$

$$\Delta\phi \sim \frac{1}{N}$$

Kok *et al.*, Phys. Rev. A **65**, 052104 (2002)

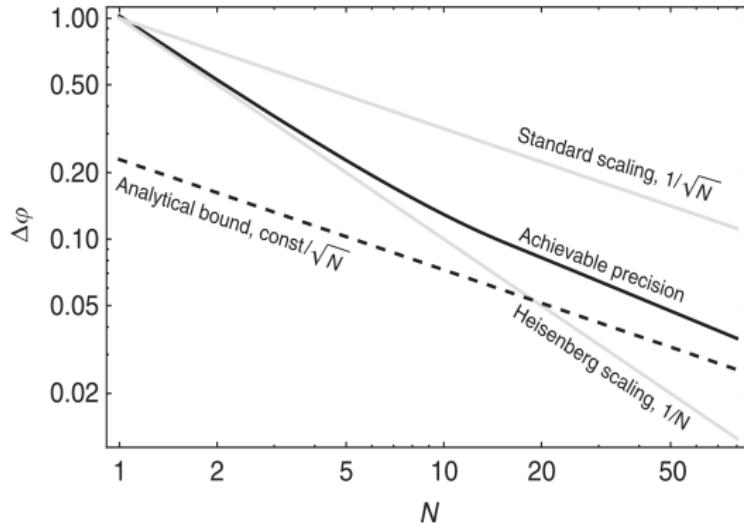
## Classical scaling

$$Q_\phi \sim N$$

$$\Delta\phi \sim \frac{1}{\sqrt{N}}$$

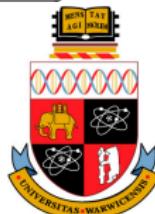


# Noisy quantum metrology



Demkowicz-Dobrzanski/Kolodynski/Guta, Nat. Comms. 3, 1063, (2012)

- ✗ No quantum scaling with noise
- ✗ Quantum scaling lost no matter how small the noise
- ✗ Quantum-enhanced scaling is impossible!



# Error corrected quantum metrology

- Attempts to recover quantum-enhanced scaling

Preskill, quant-ph:0010098

- Assume specific forms of noise (and all others to be absent)

W. Dur *et al.*, PRL, 112, 080801 (2014)

D. A. Herrera-Martí *et al.*, PRL, 115, 200501 (2015)

- Assume short sensing times to commute noise to the end

E. M. Kessler *et al.*, PRL, 112, 150802 (2014)

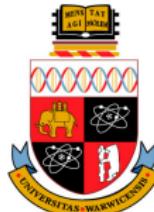
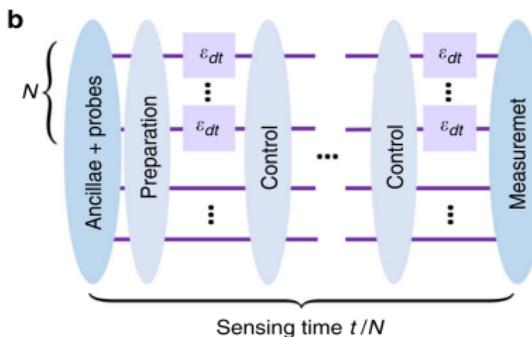
X.-M. Lu *et al.*, Nat. Comms. 6, 7282 (2015)

- Assume instantaneous, perfect correction & control operations

G. Arrad *et al.*, PRL, 112, 150801 (2014)

Demkowicz-Dobrzanski *et al.* PRX, 7, 1, (2017)

Zhou *et al.*, Nat. Comms. 9, 78, (2018)



- Error correction helps iff generator  $\notin$  span of noise

# Our results

- Separate noise into two types
  - beyond our control: associated with the field ( $R_z(\phi)$ )
  - under our control: devices (prepare/measure probes & ancillae)
- Introduce noise thresholds
- ✓✓ Show better devices counter more noise beyond our control
- Give actual variances (not bounds and scalings)
- Retrieve more information with **local, full-rank Pauli noise**
- No assumption on time-scales

arXiv.org > quant-ph > arXiv:1807.04267

Quantum Physics

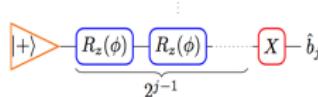
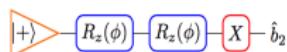
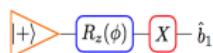
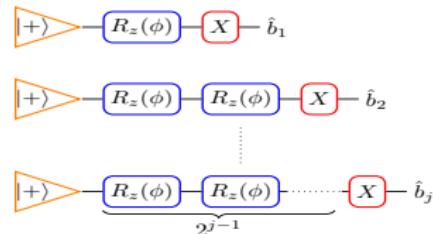
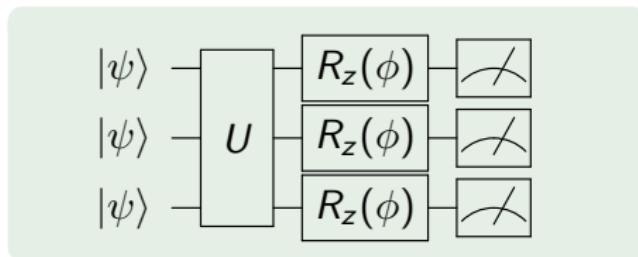
## Fault-tolerant quantum metrology

Theodoros Kapourniotis, Animesh Datta

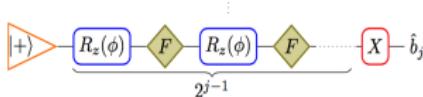
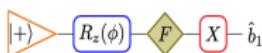
(Submitted on 11 Jul 2018)



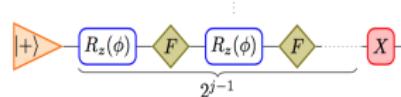
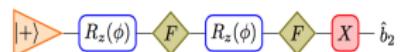
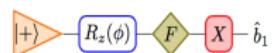
# Fault tolerance (FT) quantum metrology



(a) Protocol Ia



(b) Protocol Ib



(c) Protocol Ic

**Ia:** Noise everywhere, FT nowhere

**Ib:** Noiseless/noisy devices + field noise, FT for **field only**

**Ic:** Noise everywhere, FT everywhere



# Fault tolerant quantum metrology

Since  $\phi \in \mathbb{R}$ , we cannot have

- ✗ a stabiliser code that is transversal for  $R_z(\phi) = e^{-i\phi Z}$

Jochym-O'Connor/Kubica/Yoder, PRX 8, 021047 (2018)

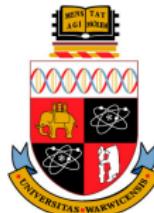
## Digital phase estimation

$$\phi = 2\pi \times 0.b_0 b_1 b_2 \dots = b_0 \pi + b_1 \frac{\pi}{2} + b_2 \frac{\pi}{4} + \dots$$

$$\text{If } T_n = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\frac{2\pi}{2^n}} \end{pmatrix}, R_z(\phi) = T_1^{b_0} T_2^{b_1} T_3^{b_2} \dots$$

Since  $\phi$  is unknown, we cannot use

- ✗ gate synthesis to acquire a FT gate set
- ✗ magic state distillation
- ✗ state twirling to diagonalise noise in the magic state basis



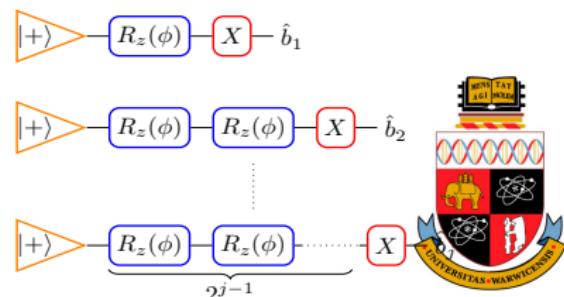
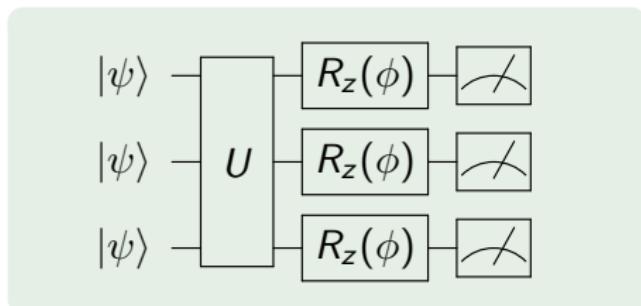
# Performance of FT quantum metrology depends on

- Noise model

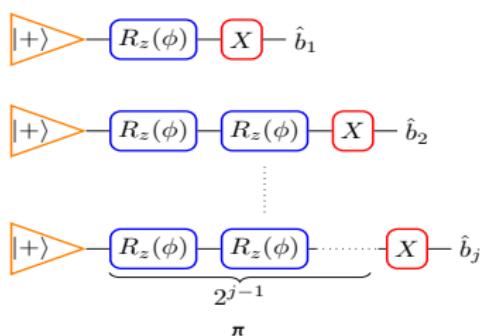
$$\mathcal{E}(\rho) = (1 - p)\rho + p(p_x X \rho X + p_y X Z \rho Z X + p_z Z \rho Z),$$

$$0 \leq p, p_x, p_y, p_z \leq 1, \quad p_x + p_y + p_z = 1$$

- Local, full-rank noise that applies everywhere
- Error correcting code
- Estimator (encapsulates the protocol)



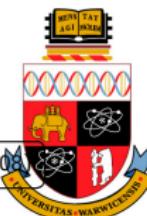
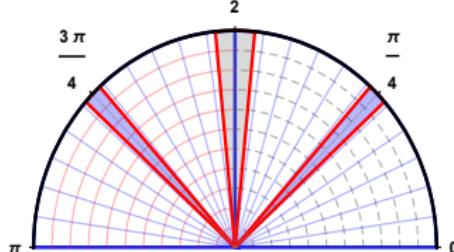
# Bit-wise estimator



- Set  $j = 1$
- $p_1$ : Prob. of obtaining +1
- $\hat{p}_1$ : Estimate of obtaining +1
- $\text{prob}(|\hat{p}_1 - p_1| \leq \delta) \geq 1 - e^{-2M\delta^2}$
- $\delta = |\cos^2(\frac{\pi}{4}) - \cos^2(\frac{\pi}{4} - \frac{\gamma}{2})| = |\sin \gamma|/2$
- If  $0 \leq \phi < \frac{\pi}{2} - \gamma$ ,

$$\text{prob}(0 \leq \hat{\phi} < \frac{\pi}{2}) \geq 1 - e^{-2M\delta^2}$$

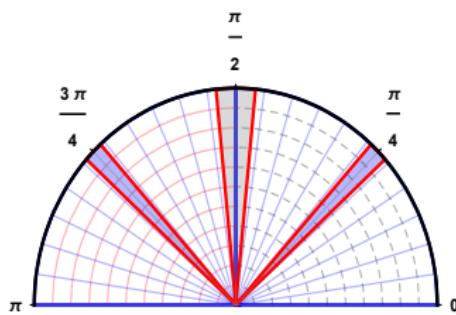
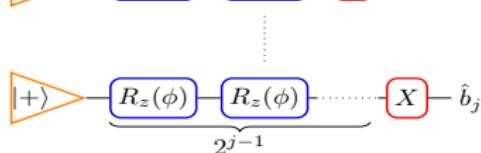
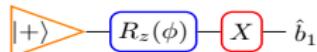
$$\Rightarrow \text{prob}(\hat{b}_1 = b_1 = 0) \geq 1 - e^{-2M\delta^2}$$



Rudolph/Grover, PRL. 91, 217905, (2003)

Ji/Wang/Duan/Feng/Ying, IEEE Trans. Inf. Th. 54, 5172 (2008)

# Bit-wise estimator



## Protocol Ia

For  $j = 1, \dots, t$

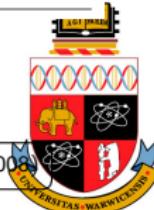
1. Repeat  $M$  times:
  - (i) Prepare  $|+\rangle$ .
  - (ii) Interrogate field  $2^{j-1}$  times.
  - (iii) Measure  $X$ .
2. Calculate  $\hat{p}_j$  as the fraction of the  $+1$  measurement outcomes out of  $M$ . If  $\hat{b}_{j-1} = 0$  set  $\hat{\phi}_j = \cos^{-1}(2\hat{p}_j - 1)$  in  $[0, \pi]$ , or else in  $[\pi, 2\pi]$ . If
  - (i)  $\hat{b}_{j-1}\pi \leq \hat{\phi}_j < \hat{b}_{j-1}\pi + (\pi/2 - \gamma)$ , set  $\hat{b}_j = 0$ .
  - (ii)  $\hat{b}_{j-1}\pi + (\pi/2 + \gamma) \leq \hat{\phi}_j \leq \hat{b}_{j-1}\pi + \pi$ , set  $\hat{b}_j = 1$ .Otherwise output estimate up to bit  $j - 1$  and exit.
3. If  $j \neq t$  increase  $j$  by one and go to step 1, otherwise exit and output

$$\hat{\phi} = \hat{b}_1 \frac{\pi}{2} + \hat{b}_2 \frac{\pi}{4} + \dots + \hat{b}_t \frac{\pi}{2^t}$$

- Convergences if outputs the first  $t$  bits with confidence  $\epsilon$

Rudolph/Grover, PRL. 91, 217905, (2003)

Ji/Wang/Duan/Feng/Ying, IEEE Trans. Inf. Th. 54, 5172 (2008)



# Ia: No fault tolerance

- Local, full-rank noise

$$\mathcal{E}(\rho) = (1-p)\rho + p(p_x X \rho X + p_y X Z \rho Z X + p_z Z \rho Z),$$

- For  $\hat{b}_j$ , failure probability  $1 - (1-p)^{2^{j-1}}$
- Protocol Ia converges if  $p < p_{\text{th}}$  which is the solution of

$$1 - (1-p)^{2^{j-1}} = |\sin \gamma|/2.$$

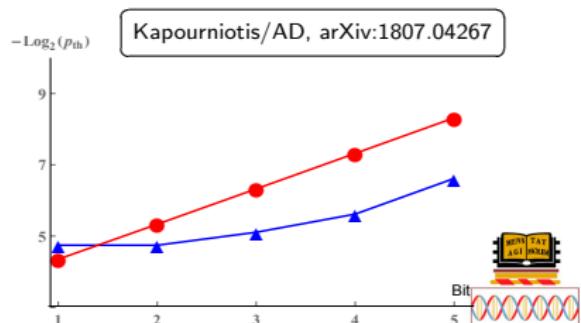
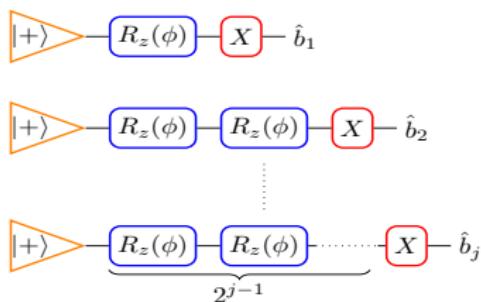
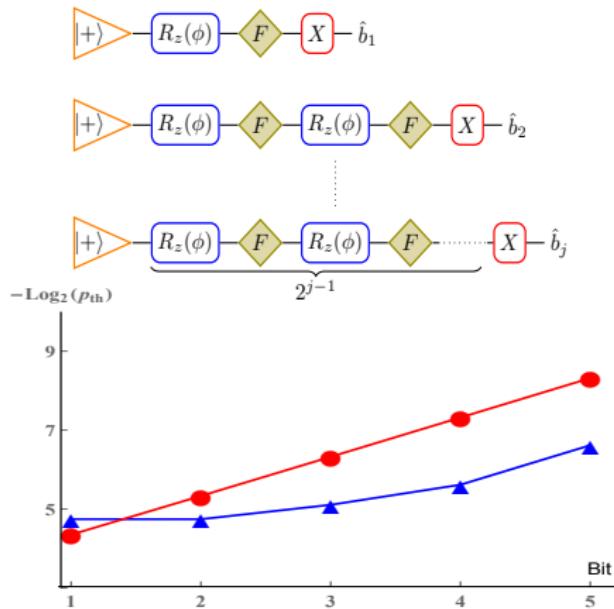


Figure: Protocol Ia,  $\gamma = \pi/32$ . (Later!)

# Ib: Fault tolerance for noise beyond our control



(assume noiseless devices)

## Protocol Ib

For  $j = 1, \dots, t$

1. Repeat  $M$  times
  - (i) Prepare probe  $|+\rangle$ . Set  $k = 1$ .
  - (ii) Prepare ancilla  $|0\rangle$ . Apply CNOT between probe and ancilla. Encode probe by QRM( $1, j+2$ ).
  - (iii) Interrogate field transversally with probe. Apply error detection on probe. Restart (i) if syndrome measurements reject.
  - (iv) Teleport by measuring probe in logical  $X$  and adapting Pauli frame accordingly (See Fig. (3)).
  - (v) If  $k < 2^{j-1}$ , increase  $k$  by one, use ancilla as new probe and return to (ii).
  - (vi) Measure  $X$ .
2. Step 2 of Protocol Ia with  $\gamma$  replaced by  $\gamma'$ .
3. If  $j \neq t$  increase  $j$  by one and go to step 1, otherwise exit and output

$$\hat{\phi} = \hat{b}_1 \frac{\pi}{2} + \hat{b}_2 \frac{\pi}{4} + \dots + \hat{b}_t \frac{\pi}{2^t}$$

Figure: Protocol Ia, Protocol Ib.  $\gamma = \frac{\pi}{32}$ .

- Threshold is solution

$$1 - (1 - p_{\text{err}}^X)^{2^{j-1}} (1 - p_{\text{err}}^Z)^{2^{j-1}} = |\sin(\gamma')|/2$$



# Ib: Fault tolerance for noise beyond our control

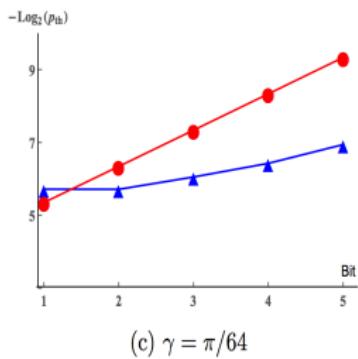
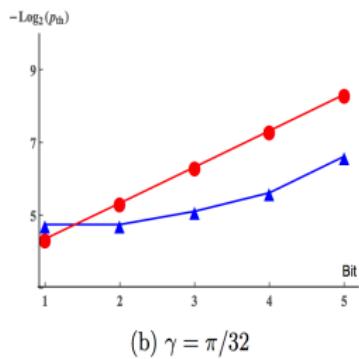
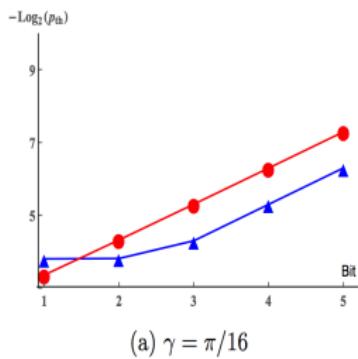
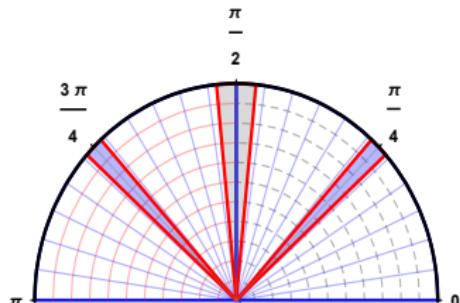
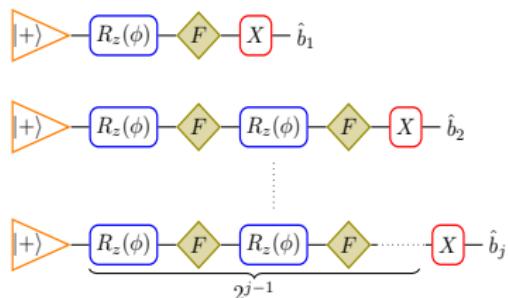
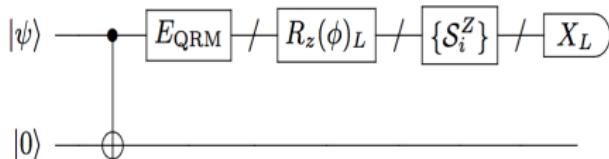
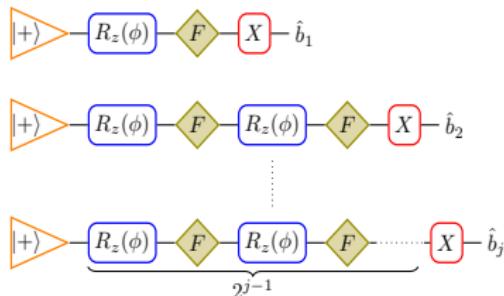


Figure: Protocol Ia, Protocol Ib (assume noiseless devices).

# Ib: Fault tolerance for noise beyond our control



**Figure:** For  $k = 1$ ,  $|\psi\rangle = |+\rangle$ , else output of  $k - 1$ .  $\{\mathcal{S}_i^Z\}$ : Z stabiliser measurements,  $X_L$ : logical  $X$  syndromes. Only  $R_z(\phi)_L$  is noisy.

- Error detection, NOT correction

(assume noiseless devices)

## Protocol Ib

For  $j = 1, \dots, t$

1. Repeat  $M$  times

- (i) Prepare probe  $|+\rangle$ . Set  $k = 1$ .
- (ii) Prepare ancilla  $|0\rangle$ . Apply CNOT between probe and ancilla. Encode probe by QRM( $1, j+2$ ).
- (iii) Interrogate field transversally with probe. Apply error detection on probe. Restart (i) if syndrome measurements reject.
- (iv) Teleport by measuring probe in logical  $X$  and adapting Pauli frame accordingly (See Fig. (3)).
- (v) If  $k < 2^{j-1}$ , increase  $k$  by one, use ancilla as new probe and return to (ii).
- (vi) Measure  $X$ .

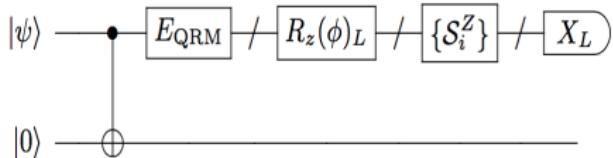
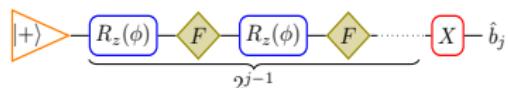
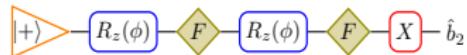
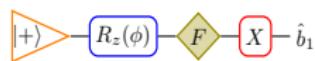
2. Step 2 of Protocol Ia with  $\gamma$  replaced by  $\gamma'$ .

3. If  $j \neq t$  increase  $j$  by one and go to step 1, otherwise exit and output

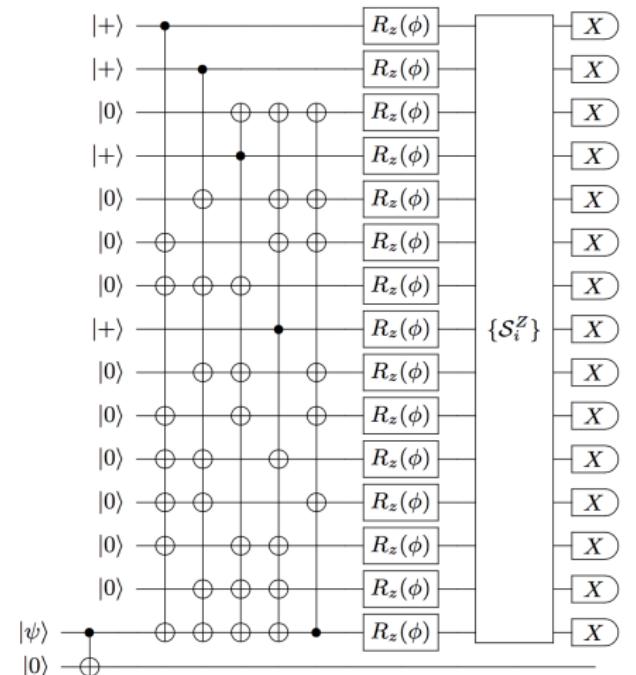
$$\hat{\phi} = \hat{b}_1 \frac{\pi}{2} + \hat{b}_2 \frac{\pi}{4} + \dots + \hat{b}_t \frac{\pi}{2^t}$$



# Ib: Fault tolerance application of traversal $R_z(\phi)$



(assume noiseless devices)

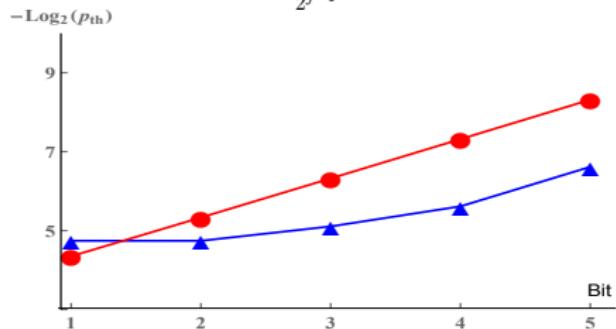
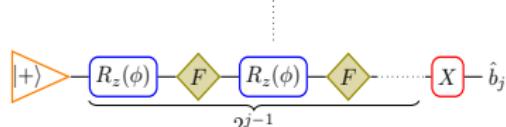
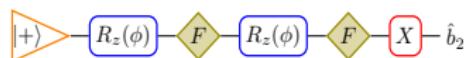


**Figure:** FT application of transversal  $R_z(\phi)$  using QRM(1,4).  
Only  $R_z(\phi)$  is noisy.

- Error detection, NOT correction



## Ib: Fault tolerance for noise beyond our control



(assume noiseless devices)

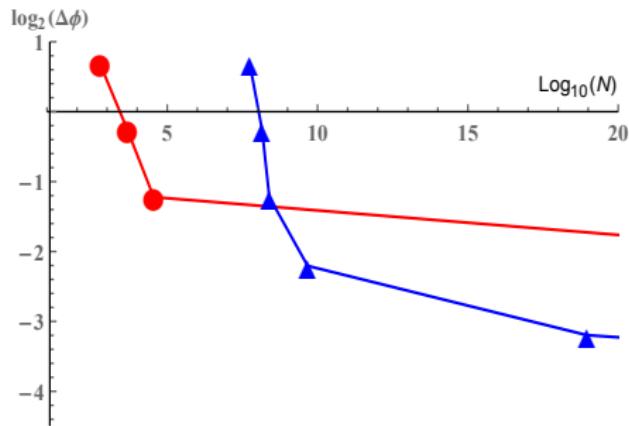


Figure: **Ia, Ib.**  $\gamma = \pi/32$ ,  $p = 0.63\%$

Figure: **Protocol Ia, Protocol Ib.**  $\gamma = \frac{\pi}{32}$ .

- Threshold is solution to

$$1 - (1 - p_{\text{err}}^X)^{2^{j-1}} (1 - p_{\text{err}}^Z)^{2^{j-1}} = |\sin(\gamma')|/2$$



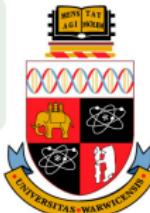
# Quantum Reed-Muller codes (very briefly)

- Quantum stabiliser codes

Steane, IEEE Trans. Inf. Th. 45, 1701 (1999)

- QRM( $1, n+1$ ) transversal for  $T_j = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\frac{2\pi}{2^j}} \end{pmatrix}$ , for  $j \leq n$
- Transversality allows logical operation on  $2^{n+1} - 1$  qubits  $\times$
- Operate as error detecting code for full-rank noise in  $T_n$  gates
- QRM( $1, n+1$ ) **not** transversal for  $T_j$  for  $j > n$

$$\gamma' = \gamma - 2 \arctan \left( \frac{\sin(2^{j+1}\gamma)}{(2^{j+2}-1) + \cos(2^{j+1}\gamma)} \right) \sim \gamma - O(2^{-j})$$



# Noisy devices

- Noisy non-transversal encoding
- Noisy syndrome measurements
- Assume device noise independent of  $\phi$

$$\mathcal{E}(\rho) = (1 - p')\rho + p'(p'_x X \rho X + p'_y X Z \rho Z X + p'_z Z \rho Z),$$

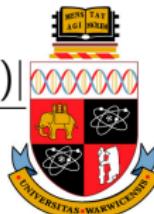
in addition of noise beyond our control

- Threshold equation changes  
from

$$1 - (1 - p_{\text{err}}^X)^{2^{j-1}} (1 - p_{\text{err}}^Z)^{2^{j-1}} = \frac{|\sin(\gamma')|}{2}$$

to

$$1 - (1 - p'_{\text{err}}^X)^{2^{j-1}} (1 - p'_{\text{err}}^Z)^{2^{j-1}} (1 - p')^{3 \times 2^{j-1} + 2} = \frac{|\sin(\gamma')|}{2}$$



# Noisy devices

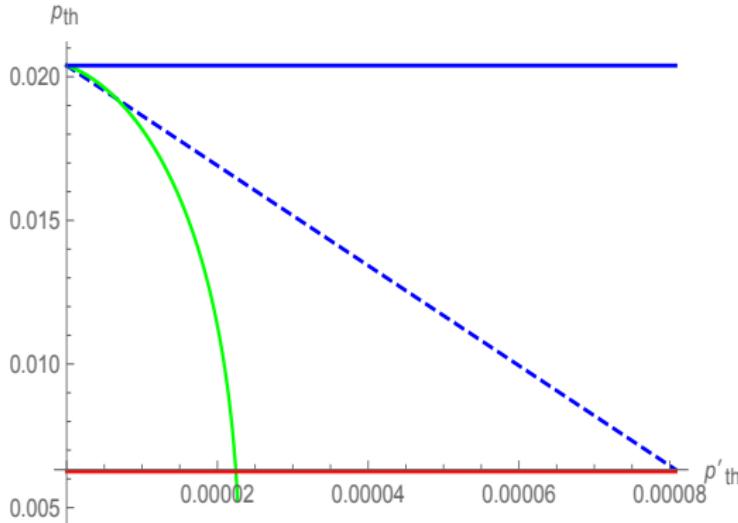


Figure:  $j = 4, \gamma = \pi/32$ . Protocol Ia.  
Ib with device noise (dashed).  
Ib without device noise (solid).  
(Later!)

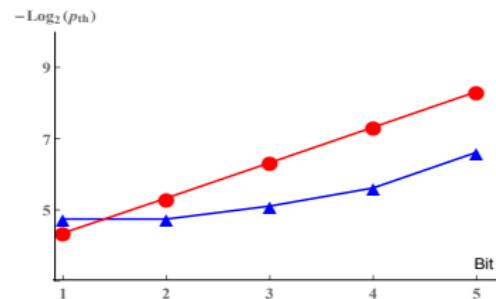
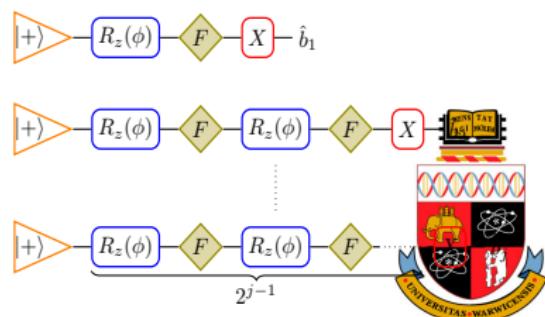
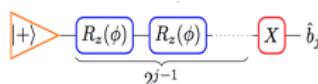
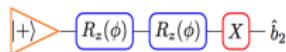
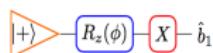
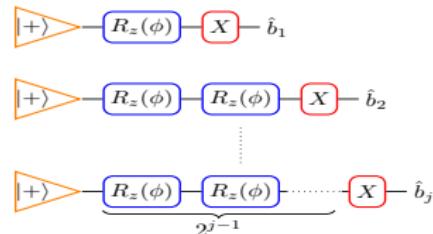
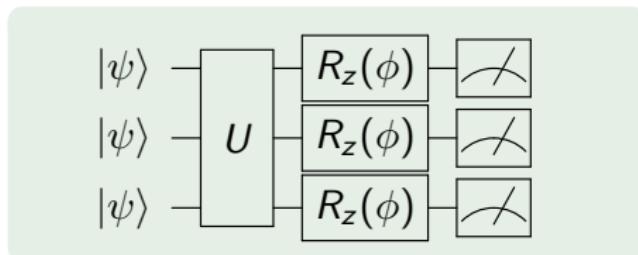


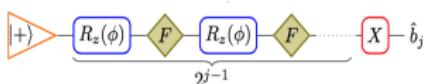
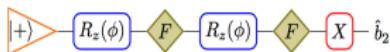
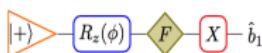
Figure:  $\gamma = \pi/32, p = 0.63\%$



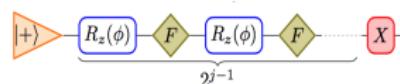
# Fault tolerance (FT) quantum metrology



(a) Protocol Ia



(b) Protocol Ib



(c) Protocol Ic

Ia: Noise everywhere, FT nowhere

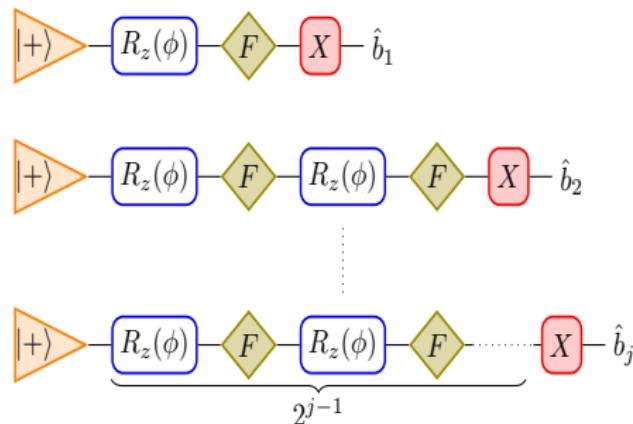
Ib: Noiseless/noisy devices + field noise, FT for **field only**

Ic: Noise everywhere, FT everywhere



# Ic: Fault tolerance everywhere

- Local full rank noise everywhere



---

## Protocol Ic

For  $j = 1, \dots, t$

1. Repeat  $M$  times
  - (i) Prepare  $|+_L\rangle$  using FT procedure employing the Steane code and switch to QRM(1,  $j+2$ ). Set  $k = 1$ .
  - (ii) Prepare ancilla  $|0_L\rangle$  using FT procedure employing QRM(1,  $j+2$ ). Apply transversal FT CNOT between probe and ancilla.
  - (iii) Interrogate field transversally with probe. Apply error detection on probe. Restart (i) if syndrome measurements reject.
  - (iv) Teleport by measuring probe in logical  $X$  and adapting Pauli frame accordingly (See Fig. (9) in Appendix F2).
  - (v) If  $k < 2^{j-1}$ , increase  $k$  by one, use ancilla as new probe and return to (ii).
  - (vi) FT measurement of logical  $X$ .
2. Step 2 and 3 of Protocol Ib.



# Fault tolerance everywhere

- Local full rank noise everywhere

---

## Protocol Ib

For  $j = 1, \dots, t$

- Repeat  $M$  times

- Prepare probe  $|+\rangle$ . Set  $k = 1$ .
- Prepare ancilla  $|0\rangle$ . Apply CNOT between probe and ancilla. Encode probe by QRM(1,  $j+2$ ).
- Interrogate field transversally with probe. Apply error detection on probe. Restart (i) if syndrome measurements reject.
- Teleport by measuring probe in logical  $X$  and adapting Pauli frame accordingly (See Fig. (3)).
- If  $k < 2^{j-1}$ , increase  $k$  by one, use ancilla as new probe and return to (ii).
- Measure  $X$ .

- Step 2 of Protocol Ia with  $\gamma$  replaced by  $\gamma'$ .

- If  $j \neq t$  increase  $j$  by one and go to step 1, otherwise exit and output

$$\hat{\phi} = \hat{b}_1 \frac{\pi}{2} + \hat{b}_2 \frac{\pi}{4} + \dots + \hat{b}_t \frac{\pi}{2^t}$$

---

## Protocol Ic

For  $j = 1, \dots, t$

- Repeat  $M$  times

- Prepare  $|+_L\rangle$  using FT procedure employing the Steane code and switch to QRM(1,  $j+2$ ). Set  $k = 1$ .
- Prepare ancilla  $|0_L\rangle$  using FT procedure employing QRM(1,  $j+2$ ). Apply transversal FT CNOT between probe and ancilla.
- Interrogate field transversally with probe. Apply error detection on probe. Restart (i) if syndrome measurements reject.
- Teleport by measuring probe in logical  $X$  and adapting Pauli frame accordingly (See Fig. (9) in Appendix F2).
- If  $k < 2^{j-1}$ , increase  $k$  by one, use ancilla as new probe and return to (ii).
- FT measurement of logical  $X$ .

- Step 2 and 3 of Protocol Ib.

- Switching between QRM(1,  $n$ ) and Steane code (QRM(1,3))



# Why code switching?

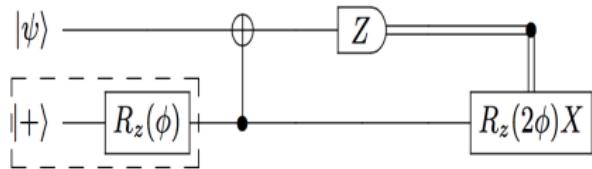


Figure: Gate distillation

$$R_z(\phi)X R_z^\dagger(\phi) \propto R_z(2\phi)X$$

- For  $b_k, \sim R_z(2^k\phi)$
- Effective noise  $\sim 2^k p$
- Loose FT advantage

## Protocol Ic

For  $j = 1, \dots, t$

1. Repeat  $M$  times
  - (i) Prepare  $|+_L\rangle$  using FT procedure employing the Steane code and switch to QRM(1,  $j+2$ ). Set  $k = 1$ .
  - (ii) Prepare ancilla  $|0_L\rangle$  using FT procedure employing QRM(1,  $j+2$ ). Apply transversal FT CNOT between probe and ancilla.
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  - (v) If  $k < 2^{j-1}$ , increase  $k$  by one, use ancilla as new probe and return to (ii).
  - (vi) FT measurement of logical  $X$ .
2. Step 2 and 3 of Protocol Ib.

- Switching between QRM(1,  $n$ ) and Steane code (QRM(1,3))
- Error correction in the device (still error detection in field)

$$1 - (1 - p_{\text{err}}''X)^{2^{j-1}} (1 - p_{\text{err}}''Z)^{2^{j-1}} (1 - p_{\text{EC}})^{3 \times 2^{j-1} + j + 1} = \frac{|\sin \gamma'|}{2}$$



# Fault tolerance everywhere

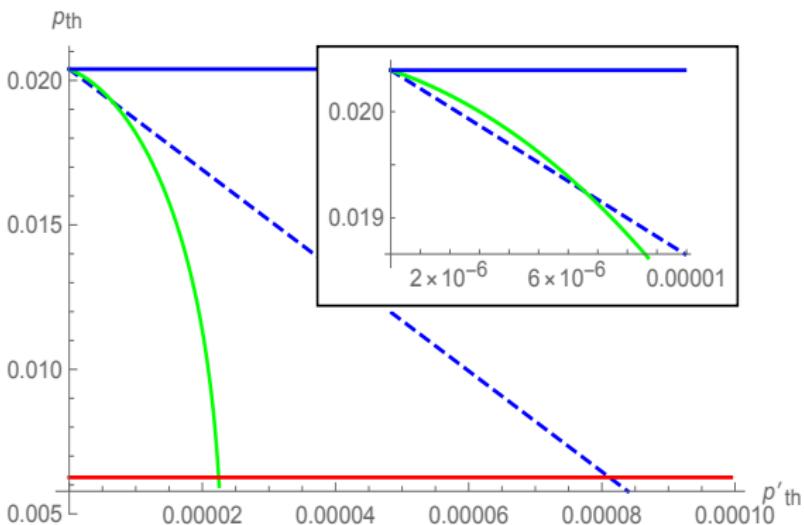
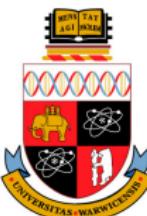


Figure:  $j = 4, \gamma = \pi/32$ . Protocol Ib with device noise (dashed).  
Ib without device noise (solid). Protocol Ic



- ✓ Better devices counter noise beyond our control
- Small improvements!!!

So,

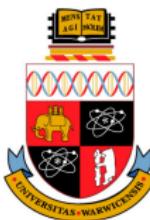
Fault tolerant quantum metrology can ...

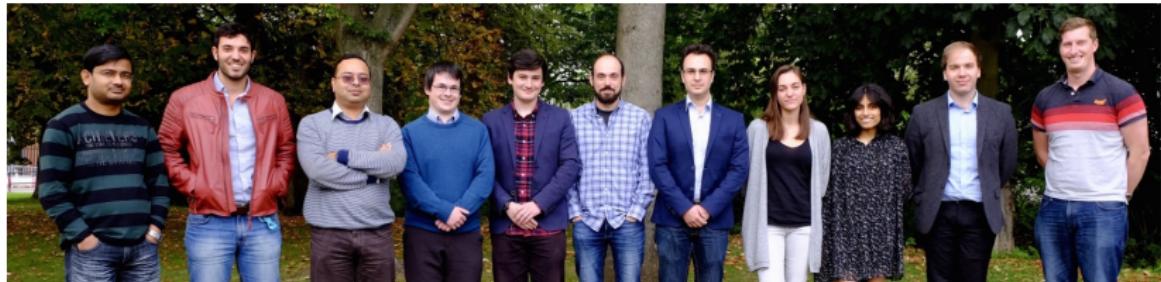
- protect from noise everywhere (fairly loose assumptions)
- estimate higher bits, giving more precision
- counter noise beyond our control with better devices ✓

Much room for improvements. We need ...

- codes with better rates ( $\text{QRM}(1, n)$  has rate  $\sim 2^n$  ✗)
- estimator and code optimisation in tandem
- better understanding of fault tolerance on unknown gates

Kapourniotis/AD, arXiv:1807.04267





- Dominic Branford
- Samuele Ferracin
- Jamie Friel
- Evangelia Bisketzi
- Aiman Khan
- Andrew Jackson
- Theodoros Kapourniotis
- Max Marcus
- Francesco Albarelli

Thank you!

[www.warwick.ac.uk/qinfo](http://www.warwick.ac.uk/qinfo)



## Quantum sensors for fundamental physics

We are looking for a post-graduate student to join the quantum information science group of **Animesh Datta** at the University of Warwick. The goals of this theoretical project are to produce the design principles for quantum sensors that can tackle some of the most fundamental open problems in physics. Instances include the direct detection of dark matter, testing the validity of quantum mechanics in macroscopic systems, searching for time variation of fundamental constants, and the direct detection of gravitational waves from exotic sources. The principle underlying all of these quests is the precise sensing of physical observables such as exquisitely small forces, phases, displacements and temperature.

The student must be interested in a close interplay of quantum metrology, quantum information science, quantum optics, and quantum mechanics.

[www.warwick.ac.uk/qinfo](http://www.warwick.ac.uk/qinfo)

