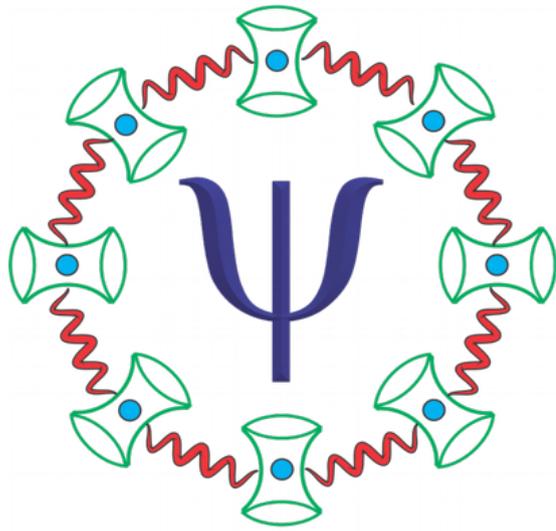


Mutual uncertainty, conditional uncertainty and strong subadditivity



QIPA



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(PRA 98, 032123)



<http://www.hri.res.in/~qic>

Uncertainty

- Which implies information. Statistical fluctuation
- **Intrinsic** and inevitability
- It carries through time
- Measures: variance, entropy

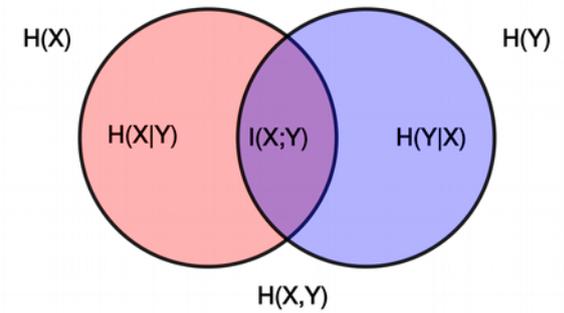


Motivations

- Sum uncertainty Relation: $\Delta(A + B) \leq \Delta A + \Delta B$.

—————▶ $M(A : B) := \Delta A + \Delta B - \Delta(A + B)$. PLA, 367, 177

$$\Delta A_i^2 = \langle A_i^2 \rangle - \langle A_i \rangle^2, \text{ where } \langle A_i \rangle = \text{Tr}[\rho A_i]$$



$$\hat{H} = \frac{1}{2}P^2 + \frac{1}{2}X^2$$

$$\Delta(A + B) = 0$$

$$1+1=0$$

Like entropy, the variance is

- Convex, $\Delta(\sum_i p_i A_i) \leq \sum_i p_i \Delta(A_i)$ $0 \leq p_i \leq 1$
 $\sum_i p_i = 1$

- Concave, for $\rho = \sum_\ell \lambda_\ell \rho_\ell$, $\Delta(A)_\rho \geq \sum_\ell \lambda_\ell \Delta(A)_{\rho_\ell}$

- +ve

- Variance is the 2nd moment whereas entropy includes all.



—————> Preparation uncertainty

Mutual and conditional uncertainty

- For $\{A_i; i = 1, 2, \dots, n\}$,

$$M(A_1 : A_2 : \dots : A_n) := \sum_{i=1}^n \Delta A_i - \Delta\left(\sum_{i=1}^n A_i\right).$$

- **CU:** $\Delta(A|B) := \Delta(A + B) - \Delta B.$

- **C-Variance:** $\Delta(A|B)^2 := \Delta(A + B)^2 - \Delta B^2.$

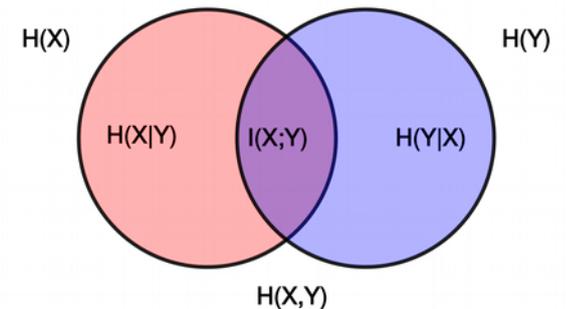
CU does't imply C-Var

Chain rule

- For sum uncertainty:

$$\Delta \left(\sum_{i=1}^n A_i \right) = \sum_{i=1}^n \Delta(A_i | A_{i-1} + \cdots + A_1).$$

$$\Delta(A + B) = \Delta(A) + \Delta(B|A)$$



Strong sub-additivity

- “Conditioning will not increase the entropy.”

$$S(\rho_{1|23}) \leq S(\rho_{1|2}).$$

- **Theorem.3** *If $M(B : C) = 0$, then $\Delta(A|B+C) \leq \Delta(A|B)$, i.e., conditioning on more observables reduces the uncertainty.*

Implies: This guarantees the +vity of MU.

- Equivalent one:

Inequality.1 Discarding the observable, one cannot increase the mutual uncertainty, i.e., $M(A : B) \leq M(A : B + C)$.

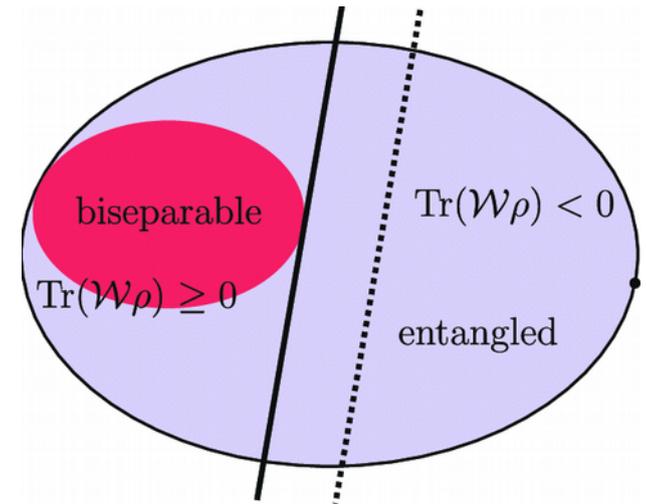
$$I(\rho_{12}) \leq I(\rho_{1(23)})$$

Physical implications

- Detection of entanglement
- and steering phenomenon

PPT criteria

Rev. Mod. Phys. **81**, 865



Entanglement detection



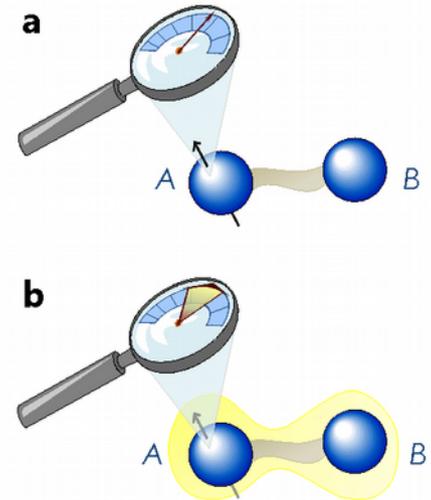
Tomography

Phys. Rev. A **68**, 032103

$$\|T\|_{KF} \leq \frac{d(d-1)}{2}$$

Quantum Inf. Comput. **7**, 624

$$\|X\|_{KF} := \sum_i \lambda_i(X) = \text{Tr}[\sqrt{X^\dagger X}]$$



N parties & d minension

$$\rho = \frac{1}{d^N} \mathbb{I}_{d^N} + \frac{1}{2d^{N-1}} [\vec{\sigma} \cdot \vec{r} \otimes \mathbb{I}_d^{\otimes N-1} + \dots + \mathbb{I}_d^{\otimes N-1} \otimes \vec{\sigma} \cdot \vec{r}_N] + \frac{1}{4d^{N-2}} \sum_{ij} [t_{ij0\dots 0} \sigma_i \otimes \sigma_j \otimes \mathbb{I}_d^{\otimes N-2} + \dots + t_{0\dots 0ij} \mathbb{I}_d^{\otimes N-2} \otimes \sigma_i \otimes \sigma_j] + \dots + \frac{1}{2^N} \sum_{i_1 \dots i_N} t_{i_1 \dots i_N} \sigma_{i_1} \otimes \dots \otimes \sigma_{i_N},$$

$$\|T^{(k)}\|_{KF} \leq \sqrt{(1/2^k) d^k (d-1)^k}$$

Quant. Inf. and Comp. **8**, 0773

$$\|T\|_{KF} \leq \frac{d(d-1)}{2}$$

Entanglement in Higher dimension

Theorem-4 For two qudit separable states and the set of observables $\{A_i\}$ and $\{B_i\}$ described above, $\sum_i \Delta(A_i|B_i)^2 \geq 2(d-1)$. This criteria is equivalent to

$\|T\|_{KF} \leq \frac{2(d-1)}{d} - \frac{1}{2}(|\vec{r}_1| - |\vec{r}_2|)^2$.

$$\|X\|_{KF} := \sum_i \lambda_i(X) = \text{Tr}[\sqrt{X^\dagger X}]$$

Compare with the extant criteria

$$A_i = \tilde{A}_i \otimes \mathbb{I}_d$$

$$\|T\|_{KF} \leq \frac{d(d-1)}{2}$$

$$\tilde{A}_i = \vec{a}_i \cdot \vec{\sigma}$$

$$\text{Tr}[\tilde{A}_i \tilde{A}_j] = 2\delta_{ij}$$

$$\tilde{A}_i = \sum_j \hat{\Theta}_{ij} \sigma_j, \text{ where } \Theta \in SO(d^2 - 1).$$

- It detects bound entangled states also.

- For $\rho = \frac{1}{4}[\mathbb{I}_4 + \frac{2}{5}(1-\alpha)\sigma_3 \otimes \mathbb{I}_2 - \frac{3}{5}(1-\alpha)\mathbb{I}_2 \otimes \sigma_3 - \alpha \sum_{i=1}^3 \sigma_i \otimes \sigma_i]$,

entangled for $\alpha > \frac{1}{19(5\sqrt{6}-6)} \blacktriangleright (0.3288)$ PPT criteria.

Theorem-4 $\longrightarrow \alpha > \frac{49}{74+5\sqrt{221}} \simeq 0.3303$

- Bound entangle state $\rho = \frac{1}{4}[\mathbb{I}_9 - \sum_i^4 |\psi_i\rangle\langle\psi_i|]$

$\|T\|_{KF} \simeq 3.1603$, Theorem-4 able to detect its entanglement

$$|\psi_0\rangle = |0\rangle(|0\rangle - |1\rangle)/\sqrt{2}, |\psi_1\rangle = (|0\rangle - |1\rangle)|2\rangle/\sqrt{2}, |\psi_2\rangle = |2\rangle(|1\rangle - |2\rangle)/\sqrt{2}$$

$$|\psi_3\rangle = (|1\rangle - |2\rangle)|0\rangle/\sqrt{2} \quad (|0\rangle + |1\rangle + |2\rangle)(|0\rangle + |1\rangle + |2\rangle)/3$$

N-qubit state

Proposition.1.– For pure N -qubit states with all pairwise correlation tensors of the form $T^{(2)} = \vec{r}_i \vec{r}_j^T$ ($i \neq j$) and the set of N observables $\{A_i\}$, the mutual uncertainty is $M(A_1 : \dots : A_N) = N - \sqrt{N}$, where r_i is the Bloch vector of i^{th} subsystem.

- For pure two qubits:

$$C = \frac{1}{2t} [2 + M(M - 4)].$$

where $C = |\langle \Psi | \sigma_2 \otimes \sigma_2 | \Psi^* \rangle|$,

Experimentally measurable with two observables.

Concurrence

$$T^{(2)} \longrightarrow [t_{0\dots 0ij}]$$

$$A_1 = \vec{a}_1 \cdot \vec{\sigma} \otimes \mathbb{I}_2 \otimes \mathbb{I}_2 \cdots, A_2 = \mathbb{I}_2 \otimes \vec{a}_2 \cdot \vec{\sigma} \otimes \mathbb{I}_2 \otimes \cdots$$

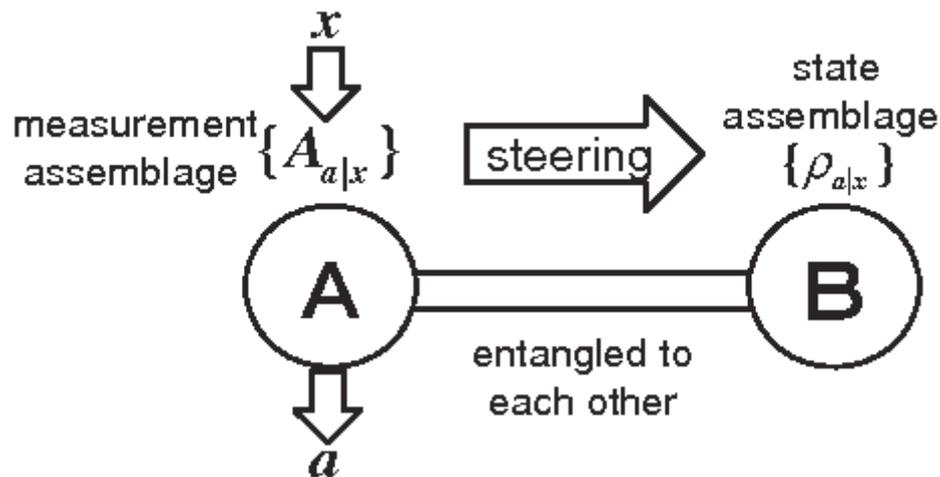
Quantum steering

- In nutshell

$$\tilde{\rho}_1^e = \sum_{\mu} p(\mu) \mathcal{P}(e|E, \mu) \rho_2^Q(\mu), \text{ where } F = \{p(\mu), \rho_2^Q(\mu)\}$$

and $\mathcal{P}(e|E, \mu)$ an ensemble prepared by Alice
stochastic map.

if Bob cannot find such F and $\mathcal{P}(e|E, \mu)$, \longrightarrow steering



Detecting steering

- SDP
- Local uncertainty at Bob's (Reid's criteria)

Phys. Rev. A **40**, 913

if $\text{Cov}(\check{A}, C) \neq 0$, $\Delta_{\text{inf}} A = \sqrt{\langle A - A_{\text{est}}(C) \rangle^2}$,

Much easier than SDP

Steering contd..

Proposition.2.– For any bipartite quantum state and any two observables, A and B , if $M_{\text{inf}}(A : B) < 0$, then the quantum state can demonstrate steering.

$$M_{\text{inf}}(A : B) = \Delta_{\text{inf}} A + \Delta_{\text{inf}} B - \Delta(A + B).$$

- **Werner state:**

$$\rho_W = p|\Psi^-\rangle\langle\Psi^-| + \frac{1-p}{4}\mathbb{I}_4,$$

$$p > \frac{1}{3} \quad \text{Entangled}$$

$$M_{\text{inf}}(A : B) = \sqrt{1-p^2} - 1/\sqrt{2}.$$

$$A = \sigma_x/2$$
$$B = \sigma_z/2$$

$$p > 1/\sqrt{2} \quad \text{Steerable}$$

Non-Gaussian state

- 2-mode squeeze vacuum + a single photon

$$\begin{aligned}
 W(X_1, P_{X_1}, X_2, P_{X_2}) &= \frac{1}{\pi^2} \exp[2 \sinh(2\alpha)(X_1 X_2 - P_{X_1} P_{X_2}) - \cosh(2\alpha) \sum_{i=1}^2 (X_i^2 + P_{X_i}^2)] [-\sinh(2\alpha)\{(P_{X_1} - P_{X_2})^2 \\
 &- (X_1 - X_2)^2\} + \cosh(2\alpha)\{(P_{X_1} - P_{X_2})^2 + (X_1 - X_2)^2\} - 1], \quad (18)
 \end{aligned}$$

$$\Delta_{\text{inf}} X_1^2 \Delta_{\text{inf}} P_{X_1}^2 = \frac{9}{2[3 \cosh(4\alpha) + 5]}$$

$$M_{\text{inf}}(X_1 : P_{X_1}) = \frac{\sqrt{3}}{2} \left(\frac{1}{\eta_-} + \frac{1}{\eta_+} \right) - (\eta_+ + \eta_-),$$

$$\eta_{\pm} = \sqrt{\cosh(2\alpha) \pm \cosh(\alpha) \sinh(\alpha)}.$$

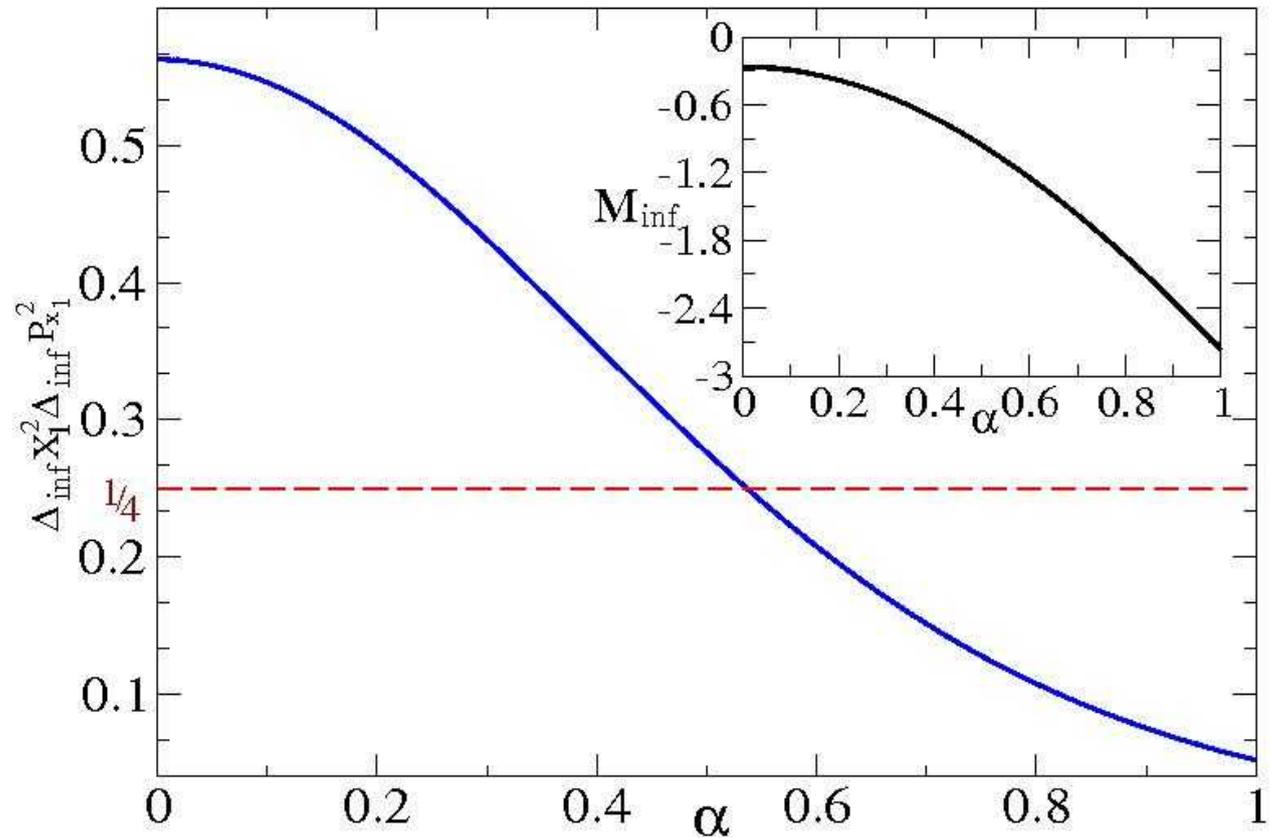
Contd.

- Where Reid's criteria fails, it doesn't.

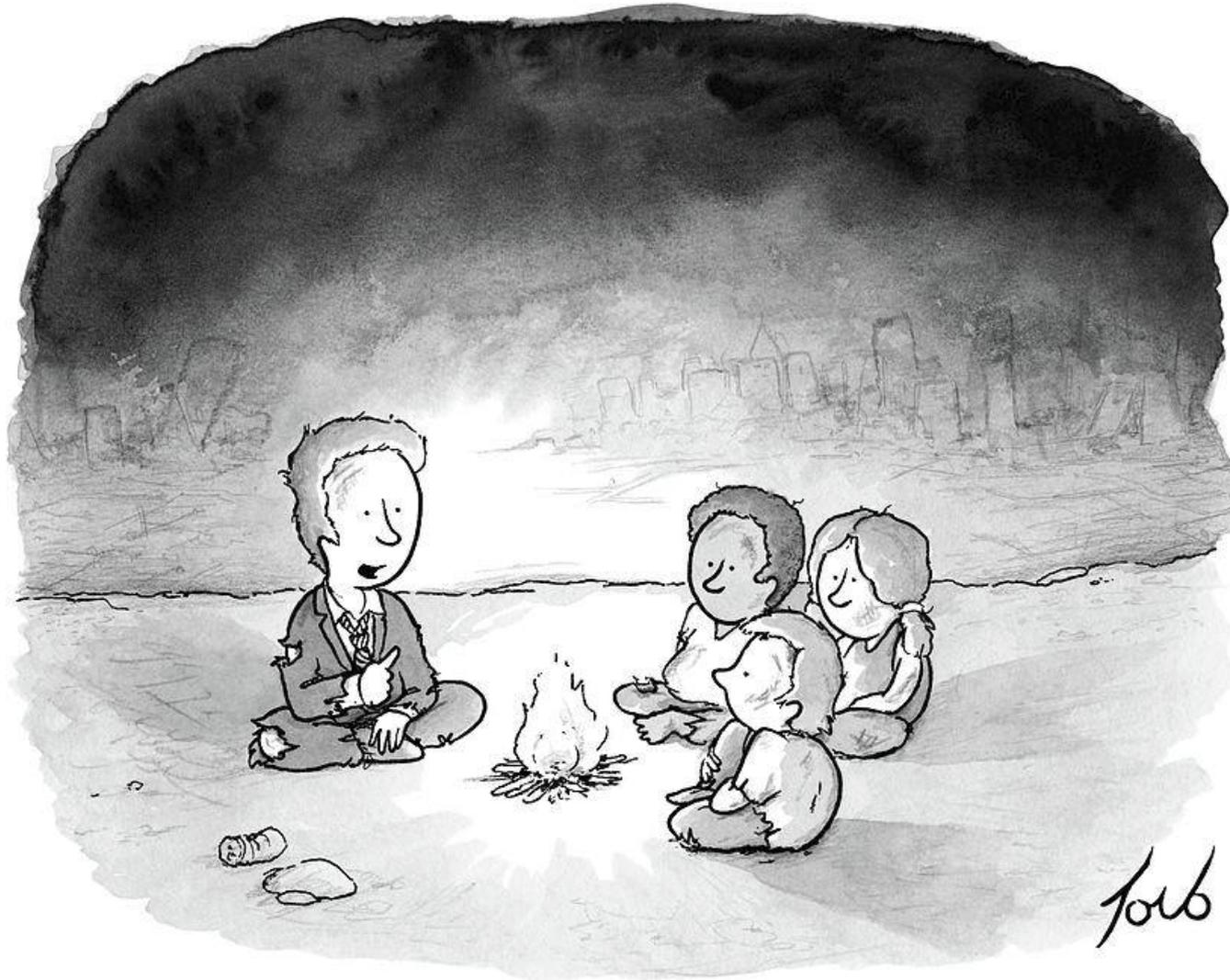
$$\Delta_{\text{inf}} X_1^2 \Delta_{\text{inf}} P_{X_1}^2 = \frac{9}{2[3 \cosh(4\alpha) + 5]}.$$

$$\Delta_{\text{inf}} X_1^2 \Delta_{\text{inf}} P_{X_1}^2 \geq 1/4$$

↑
Phys. Rev. A **40**, 913



Thanks.



"Yes, the planet got destroyed, but for a beautiful moment in time we created a lot of value for shareholders."