Optimal Binning of the Primordial Power Spectrum

Paniez Paykari CEA - Saclay

Paykari & Jaffe - APJ 2009

PFNG - India

* * * * *

 \bigstar

Primordial Power Spectrum

- e Primordial PS → physics of the structure formation in the early Universe
- Our path to its measurements is through the different surveys such as CMB and LSS surveys
- Surveys outcome → a PS that is a <u>convolution</u> of the primordial PS and a transfer function (which holds the cosmological parameters) → induces degeneracy between them
- We want to <u>quantify</u> this degeneracy

CMB & Matter Power Spectra

Primordial PS

$$\Delta_{\zeta}^2(k) \simeq A k^{n_s - 1}$$

Matter PS

$$P_{\delta}(k) = T^2(k)\Delta_{\zeta}^2(k)$$

Measurable quantity \rightarrow galaxy PS $P_g(k) = b^2(k)P_{\delta}(k)$

CMB PS

BPS

$$Radiation TF$$

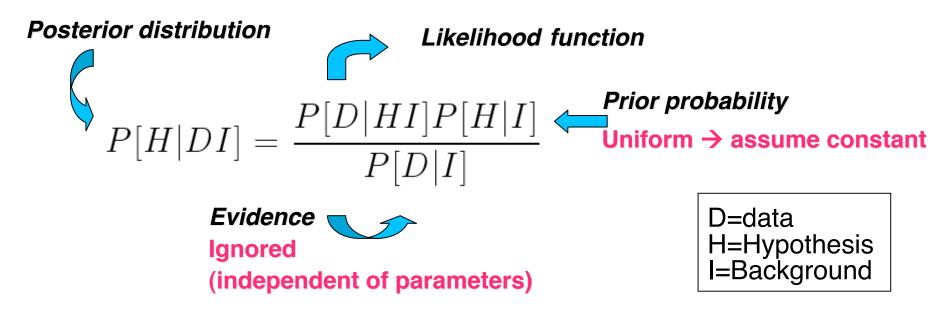
 $C_{\ell} = 2\pi \int_{0}^{\infty} d\ln k \Delta_{\zeta}^{2}(k) \Delta_{\ell}^{2}(k)$

Combine all Available Information

- Above PS <u>cannot</u> provide enough constraints on parameters on their own inherent degeneracies
- Use them all → they depend on parameters in different ways → can constrain parameters in different directions
- One interesting quantity to measure is the primordial PS measured in bins

Data Analysis

• Bayesian Statistics; takes into account prior knowledge



• A proportionality relation

 $P[H|DI] \propto P[D|HI] = \mathcal{L}$

Optimal Quadratic Estimator

Likelihood function is a

- Multivariate Gaussian (around the peak)
- Maximum in the parameter space where the parameters have the best estimated values
- $\,\,\,\sim\,\,\ln{\cal L}\,$ is maximised at the same place in the parameter space as ${\cal L}\,$

$$\frac{\partial \ln \mathcal{L}}{\partial \lambda} \big|_{\lambda = \bar{\lambda}} = 0$$

Taylor Expansion around the peak

$$F = \langle \mathcal{F} \rangle = \left\langle -\frac{\partial^2 \ln \mathcal{L}}{\partial \lambda^2} \right\rangle$$

Fisher Matrix Covariance Matrix

- Cramer-Rao equality → smallest error for a parameter is
 - Marginalised $\sqrt{(F^{-1})}$
 - Non-marginalised

 $1/\sqrt{F_{\alpha\alpha}}$

- **#** Fisher matrix
 - → generally used to determine the sensitivity of a particular survey to a set of parameters
 - → suitable for forecasting and optimizing future experiments
- Obtain a Fisher matrix for CMB and galaxy surveys \rightarrow add them to take both into account

CMB Fisher Matrix

Window En

Fisher matrix for

$$F_{\ell\ell'} = f_{sky} \frac{2\ell + 1}{2} \delta_{\ell\ell'} [C_{\ell} + w^{-1} e^{\ell^2 \sigma^2}]^{-2}$$

Weight function

$$w = (\sigma_{pix}^2 \Omega_{pix})^{-1}$$

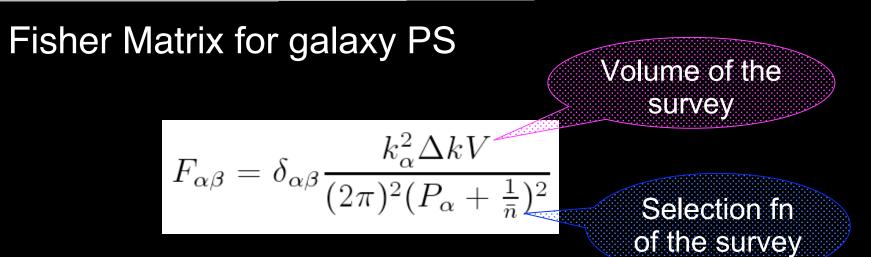
Fisher matrix for other parameters

$$F_{\alpha\beta} = \sum_{\ell} \frac{1}{(\delta C_{\ell})^2} \frac{\partial C_{\ell}}{\partial \lambda_{\alpha}} \frac{\partial C_{\ell}}{\partial \lambda_{\beta}}$$

For primordial power spectrum we have

$$\frac{\partial \mathcal{C}_{\ell}}{\partial \Delta_{\zeta}^{2}(k)} = 2\ell(\ell+1) \int_{k_{min}^{B}}^{k_{max}^{B}} dk \left| \Delta_{\ell}(k) \right|^{2}$$

Galaxy Surveys Fisher Matrix



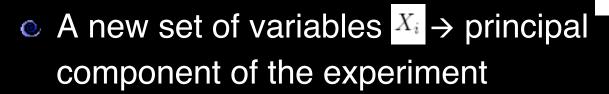
For primordial PS

$$\frac{\partial P_g(k)}{\partial \Delta_{\zeta}^2(k')} = 4 \times 2\pi^2 k T^2(k) \delta_{kk'}$$

Principal Component Analysis

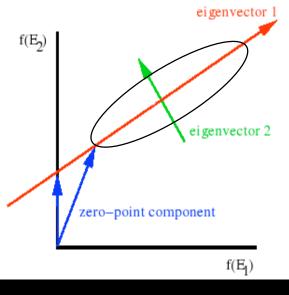
• Inverse of FM \rightarrow Covariance matrix

- $\bullet \rightarrow$ hard to interpret errors



$$\underline{C} = \underline{E}^T \underline{\Lambda} \underline{E}$$

$$\underline{X} = \underline{EO}$$



Hermitian Square Root

- Another approach to remove correlation between parameters → <u>Hermitian sqrt</u> of the FM → linear transformation on the parameter space
- Oces not give orthogonal basis
 - \rightarrow think of it as a <u>window function</u>

$$\underline{F}^{1/2} = \underline{E}^T \underline{\Lambda}^{1/2} \underline{E} \qquad (\underline{F}^{-1/2})^T \underline{F}(\underline{F}^{-1/2}) = diag$$

$$H_{nm} = \frac{(F^{1/2})_{nm}}{\sum_n (F^{1/2})_{nm}} \qquad \sum_n H_{nm} = 1$$
Define a windowed PS
$$\tilde{P}_m = \sum_n H_{nm} P(k_n)$$

Model of the Universe (Parameter Space)

- A geometrically flat, adiabatic, LCDM no dark energy
- parameter space
 - Primordial PS bins
 - Cosmological parameters

$$\Delta_{\zeta}^2(k) = \sum_B w_B(k)Q_B, \ w_B = \begin{cases} 1 & k \in B \\ 0 & k \notin B \end{cases}$$

SDSS(BRG)
$$V = (1h^{-1}\text{Gpc})^3 \& \overline{n} = \frac{10^5}{V}$$

Planck(HFI)

- 1

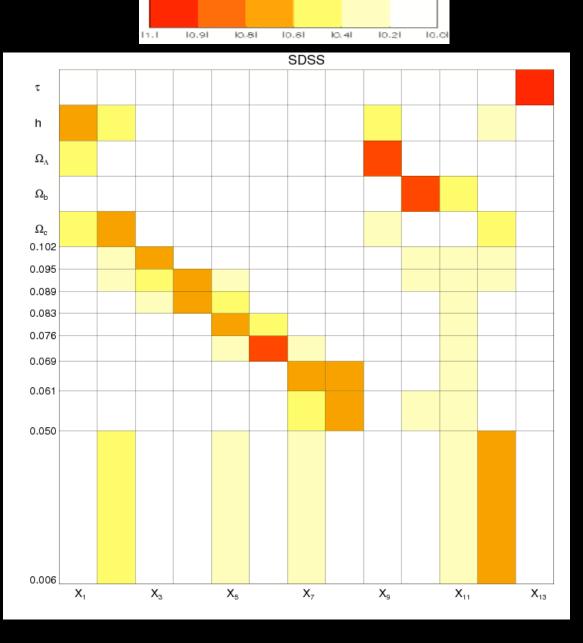
$$\nu = 100 \text{GHz}, \ \theta_{fwhm} = 10.7', \ \sigma_{pix} = 1.7 \times 10^{-15}, \ w^{-1} = 0.028 \times 10^{-15}, \ f_{sky} = 0.5$$

Optimal Binning

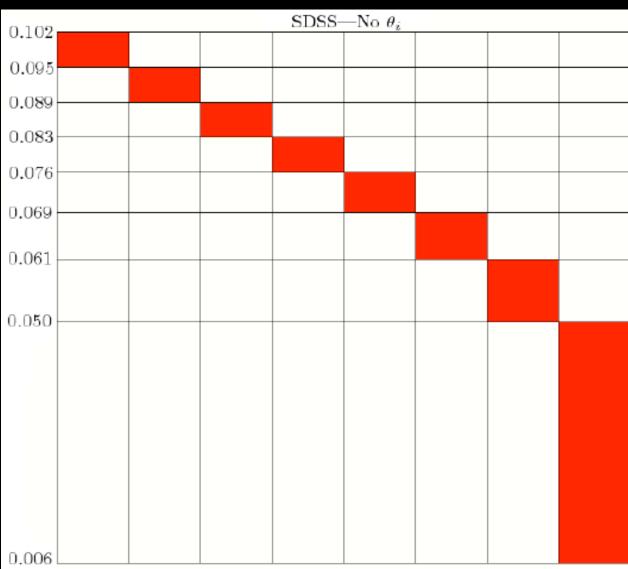
e Bins with same contribution to the Fisher matrix → <u>same S/N</u>

- Signal \rightarrow amplitude of the primordial PS in the bin
- Noise → sqrt of diagonal elements of Fisher matrix for that bin
- SDSS \rightarrow 8 bins
- Planck \rightarrow 23 bins

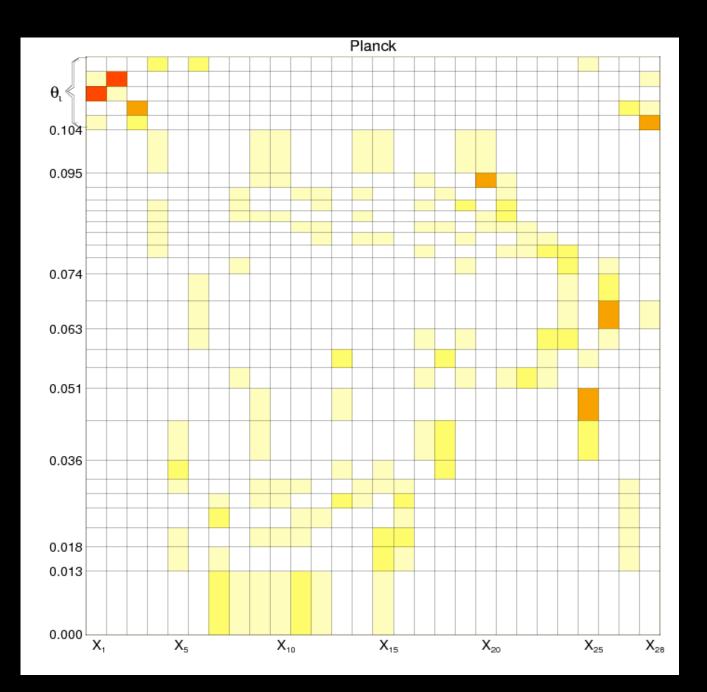
parameters

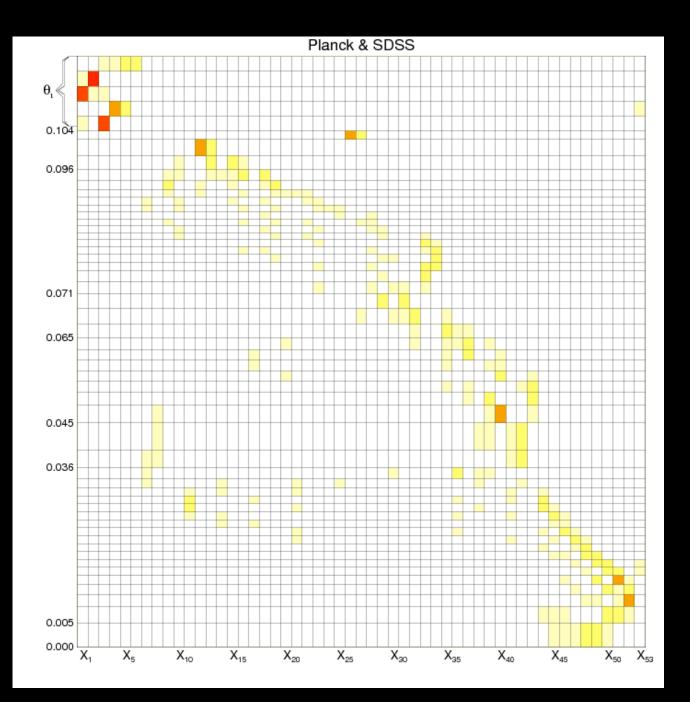


Principal components

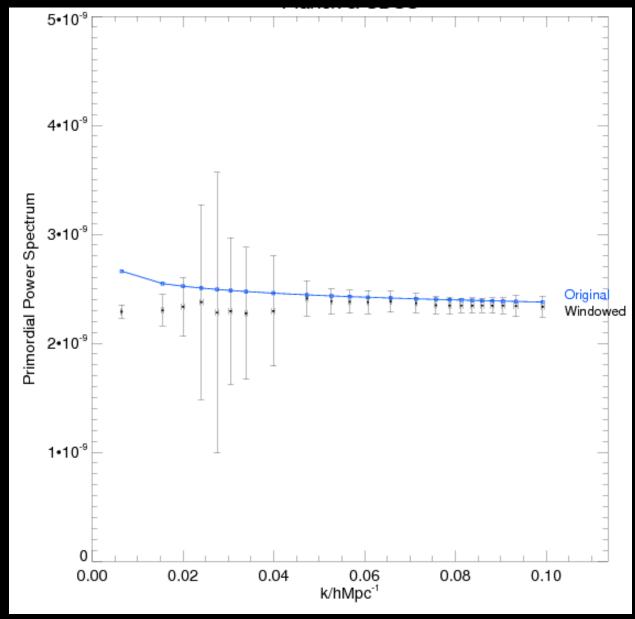


 $X_1 \quad X_2 \quad X_3 \quad X_4 \quad X_5 \quad X_6 \quad X_7 \quad X_8$

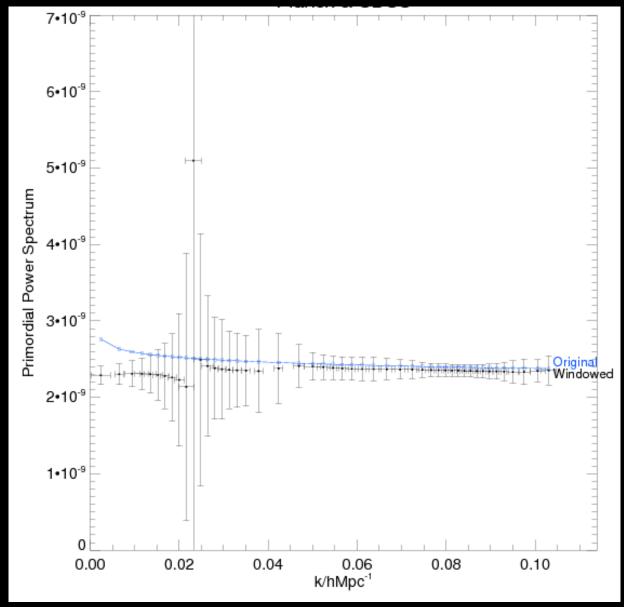




WINDOWED POWER SPECTRUM PLANCK



WINDOWED POWER SPECTRUM PLANCK & SDSS



Conclusions

- Induced degeneracy between cosmological parameters and PPS limit our ability to recover this PS completely
 - In case of Planck, even a perfect survey could recover the primordial PS completely, due to the projection effects.
- Quantify the degeneracy → parameter space with carefully chosen bins of the PPS and a set of cosmological parameters
- **<u>PCA:</u>**
 - measure cosmological parameters better
 - constrain small scales better than large scales
 - SDSS & Planck
 - break degeneracy between them
 - increase the resolution of the PPS significantly
- E <u>Hsqrt:</u>
 - divert correlation between errors of the bins to correlation between bins themselves
 - quantify the difference between the actual and measured PPS 20

Thank you