

Lower bounding convex roofs

Tobias Moroder¹

joint work with: Géza Tóth^{2,3}, Otfried Gühne¹

¹*Theoretical Quantum Optics, University of Siegen, Germany*

²*IKERBASQUE & Department of Physics, Bilbao, Spain*

³*Wigner Research Center, Budapest, Hungary*

QIPA Meeting, Allahabad, 12/08/2013

Motivation

- ▶ Convex roof, quantum information framework:
 - i) Function $f(|\psi\rangle)$ defined for pure states $|\psi\rangle$
 - ii) Extension to mixed state ρ via convex roof

$$f(\rho) = \inf_{\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|} \sum_i p_i f(|\psi_i\rangle) \quad (1)$$

- ▶ Examples:

- Bipartite concurrence: $C(|\psi\rangle_{AB}) = \sqrt{2(1 - \text{tr} \rho_A^2)}$
- Schmidt number: $|\psi\rangle_{AB} = \sum_{i=1}^r s_i |ii\rangle$, $s_i > 0$, $R_S(|\psi\rangle) = r$,
 $\Rightarrow SN(\rho_{AB}) = \inf_{p_i, |\psi_i\rangle} \max_i R_S(|\psi_i\rangle)$
- 3-tangle: $\tau(|\psi\rangle_{ABC}) = \dots$

\Rightarrow Entanglement quantification, Entanglement classification

- ▶ Lower bounds usually desired, but in general difficult to compute.

Motivation

- ▶ Convex roof, quantum information framework:
 - i) Function $f(|\psi\rangle)$ defined for pure states $|\psi\rangle$
 - ii) Extension to mixed state ρ via convex roof

$$f(\rho) = \inf_{\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|} \sum_i p_i f(|\psi_i\rangle) \quad (1)$$

- ▶ Examples:

- Bipartite concurrence: $C(|\psi\rangle_{AB}) = \sqrt{2(1 - \text{tr } \rho_A^2)}$
- Schmidt number: $|\psi\rangle_{AB} = \sum_{i=1}^r s_i |ii\rangle$, $s_i > 0$, $R_S(|\psi\rangle) = r$,
 $\Rightarrow SN(\rho_{AB}) = \inf_{p_i, |\psi_i\rangle} \max_i R_S(|\psi_i\rangle)$
- 3-tangle: $\tau(|\psi\rangle_{ABC}) = \dots$

\Rightarrow Entanglement quantification, Entanglement classification

- ▶ Lower bounds usually desired, but in general difficult to compute.

▶ This talk: Solution for polynomial functions f ◀

Main idea - Step I

⇒ explain idea via an example

- ▶ Linear entropy/Concurrence²

$$S_L(|\psi\rangle_{AB}) = [C(|\psi\rangle_{AB})]^2 = 2(1 - \text{tr} \rho_A^2) \quad (2)$$

can be rewritten as a linear operator on 2 copies

$$S_L(|\psi\rangle_{AB}) = 2[1 - \sum_i \text{tr}(\sigma_i^A \otimes \mathbb{1} |\psi\rangle \langle \psi|)^2] \quad (3)$$

$$= 2[1 - \sum_i \text{tr}(\sigma_i^A \mathbb{1} \otimes \sigma_i^{A'} \mathbb{1} |\psi\rangle \langle \psi| \otimes |\psi\rangle \langle \psi|)] \quad (4)$$

$$= \text{tr}[(2(\mathbb{1}_{ABA'B'} - \sum_i \sigma_i^A \mathbb{1} \sigma_i^{A'} \mathbb{1})) |\psi\rangle \langle \psi| \otimes |\psi\rangle \langle \psi|] \quad (5)$$

$$= \text{tr}(L_{ABA'B'} |\psi\rangle \langle \psi|^{\otimes 2}) \quad (6)$$

- ▶ More general: Any polynomial of degree n in the density operator can be written as a linear operator on n copies.

Main idea - Step II

- ▶ Evaluate $[\rho_{AB} = \sum q_i |\phi_i\rangle \langle \phi_i|]$:

$$S_L(\rho_{AB}) = \sum_i q_i S_L(|\phi_i\rangle) = \sum_i q_i \text{tr}[L |\phi_i\rangle \langle \phi_i| \otimes |\phi_i\rangle \langle \phi_i|] \quad (7)$$

$$= \text{tr}[L(\sum_i q_i |\phi_i\rangle \langle \phi_i| \otimes |\phi_i\rangle \langle \phi_i|)] = \text{tr}(L \omega_{ABA'B'}) \quad (8)$$

- ▶ The state $\omega = \omega_{ABA'B'}$ satisfies:

- i) separable state for $AB|A'B'$,
- ii) only on symmetric subspace $Sym(\mathcal{H}_{AB}^{\otimes 2})$,
- iii) reduced state $\rho_{AB} = \text{tr}_{A'B'}(\omega)$.

Main idea - Step II

- ▶ Evaluate $[\rho_{AB} = \sum q_i |\phi_i\rangle \langle \phi_i|]$:

$$S_L(\rho_{AB}) = \sum_i q_i S_L(|\phi_i\rangle) = \sum_i q_i \text{tr}[L |\phi_i\rangle \langle \phi_i| \otimes |\phi_i\rangle \langle \phi_i|] \quad (7)$$

$$= \text{tr}[L(\sum_i q_i |\phi_i\rangle \langle \phi_i| \otimes |\phi_i\rangle \langle \phi_i|)] = \text{tr}(L \omega_{ABA'B'}) \quad (8)$$

- ▶ The state $\omega = \omega_{ABA'B'}$ satisfies:

- i) separable state for $AB|A'B'$,
- ii) only on symmetric subspace $Sym(\mathcal{H}_{AB}^{\otimes 2})$,
- iii) reduced state $\rho_{AB} = \text{tr}_{A'B'}(\omega)$.

- ▶ Thus, if we optimize we obtain

$$S_L(\rho_{AB}) \geq \min \text{tr}(L \omega) \quad (9)$$

s.t. ω satisfies *i) – iii)*

- ▶ “ \leq ” Any ω with *i) – iii)* gives a decomposition as in (8).

Main idea

- ▶ Equivalent formulation:

$$S_L(\rho_{AB}) = \min \operatorname{tr}(L \omega) \quad (10)$$

s.t. $\omega = \omega_{ABA'B'}$ separable $AB|A'B'$,

$$\omega \in \mathcal{L}(\operatorname{Sym}(\mathcal{H}_{AB}^{\otimes 2})), \operatorname{tr}_{A'B'}(\omega) = \rho_{AB}$$

- ▶ Relax separability condition via

- PPT

$$S_L(\rho_{AB}) \geq \min \operatorname{tr}(L \omega) \quad (11)$$

s.t. $\omega = \omega_{ABA'B'}$ PPT $AB|A'B'$,

$$\omega \in \mathcal{L}(\operatorname{Sym}(\mathcal{H}_{AB}^{\otimes 2})), \operatorname{tr}_{A'B'}(\omega) = \rho_{AB}$$

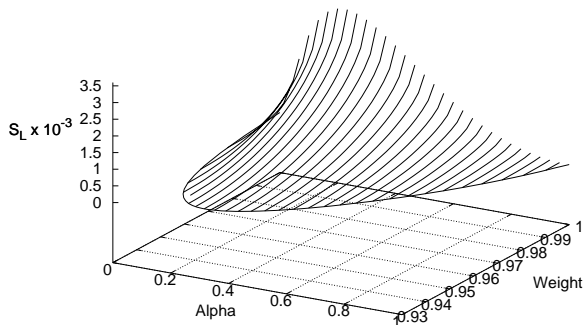
- Hierarchy of symmetric extensions [Doherty *et al.* PRL (2003)] \Rightarrow Complete solution

- ▶ Advantages:

- Optimization is a semidefinite program (certified solution)
- Dual provides “quantitative entanglement witnesses”
- Clear picture about strength

Example I - $3 \otimes 3$ Horodecki state

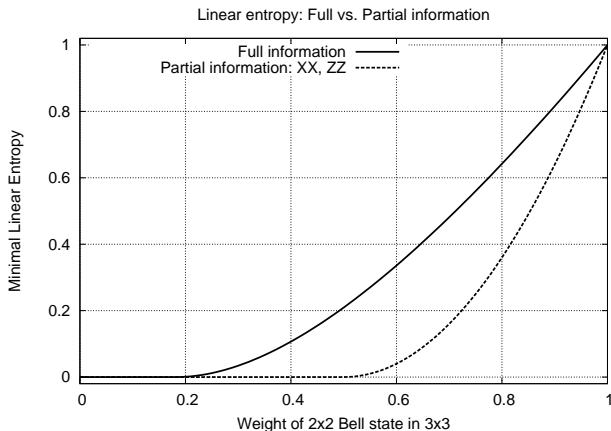
- ▶ Family of PPT entangled states: $p\rho_H(\alpha) + (1 - p)\mathbb{1}/9$



- ▶ Lower bound is positive at least for all
 - NPT entangled state
 - states with no symmetric extensions to 2 copies of A, B
- ▶ Zero for states with PPT sym. ext. to 3 copies of A, B

Example II - Partial information

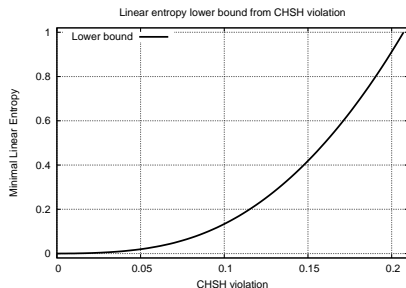
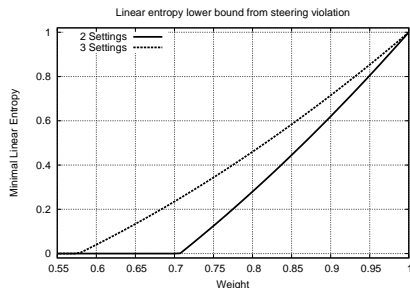
- ▶ Estimate entanglement from a few expectation value
- ▶ Example: 2-qubit Bell state embedded in $3 \otimes 3$
- ▶ Partial info.: only $\langle \sigma_x \otimes \sigma_x \rangle, \langle \sigma_z \otimes \sigma_z \rangle$ on the “qubit”



⇒ Just add constraint: $\rho_{AB}(x) = \rho_{AB}^{\text{data}} + \rho_{AB}^{\text{open}}(x) \geq 0$

Example III - Steering & device-independent scenario

- ▶ Estimate entanglement from a few measurement results and no knowledge about performed measurement
- ▶ Steering: only Alice apparatus uncharacterized
- ▶ Full dvi: both uncharacterized \Rightarrow Bell inequality



\Rightarrow Map ρ_{AB} in $?\otimes?$ to fixed finite dim. output state $\chi_{\bar{A}\bar{B}}$ by local maps \Rightarrow Estimate S_L via partial information on $\chi_{\bar{A}\bar{B}}$

Schmidt number

- ▶ Measure for the Schmidt number is given by $[\lambda_i = \lambda_i(\rho_A)]$

$$N_2(|\psi\rangle) = \sum_{i < j} \lambda_i \lambda_j = 4S_L(|\psi\rangle), \quad N_2(|\psi_{r=1}\rangle) = 0 \quad (12)$$

$$N_3(|\psi\rangle) = \sum_{i < j < k} \lambda_i \lambda_j \lambda_k, \quad N_3(|\psi_{r=2}\rangle) = 0 \quad (13)$$

$$N_4(|\psi\rangle) = \dots \quad (14)$$

- ▶ N_3 can be expressed as linear operator on 3 copies

$$N_3(|\psi\rangle) \propto 1 - 3 \operatorname{tr}(\rho_A^2) + 2 \operatorname{tr}(\rho_A^3) = \operatorname{tr}(D_3 |\psi\rangle \langle \psi|^{\otimes 3}) \quad (15)$$

- ▶ Equivalent formulation/Lower bound:

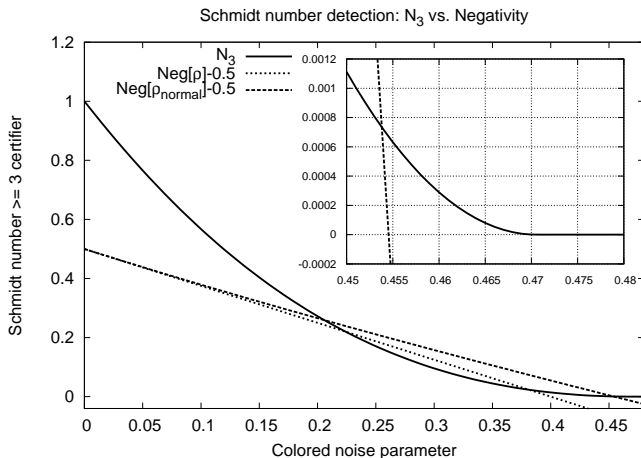
$$N_3(\rho_{AB}) = \min \operatorname{tr}(D_3 \omega_{123}) \quad (16)$$

s.t. $\omega = \omega_{123}$ fully separable

$$\omega \in \mathcal{L}(\operatorname{Sym}(\mathcal{H}_{AB}^{\otimes 3})), \operatorname{tr}_{23}(\omega) = \rho$$

Example IV - Schmidt number

- ▶ $|\psi\rangle \propto |00\rangle + |11\rangle + |22\rangle$ mixed with colored “qubit” noise



⇒ N_3 beats Schmidt number estimate via negativity

Tangle - GHZ vs. W class

- ▶ Distinguishes between the two SLOCC classes of 3 qubit multiparticle entanglement:

$$|W\rangle \propto |001\rangle + |010\rangle + |100\rangle, \quad \tau^2(|W\rangle) = 0 \quad (17)$$

$$|GHZ\rangle \propto |000\rangle + |111\rangle, \quad \tau^2(|GHZ\rangle) = 1 \quad (18)$$

- ▶ Tangle² can be expressed as linear operator on 4 copies

$$\tau^2(|\psi\rangle) = C_{A|BC}^2 - C_{AB}^2 - C_{AC}^2 = \text{tr}(T |\psi\rangle \langle \psi|^{\otimes 4}) \quad (19)$$

- ▶ Equivalent formulation/Lower bound:

$$\tau^2(\rho_{ABC}) = \min \text{tr}(T\omega_{1234}) \quad (20)$$

s.t. $\omega = \omega_{1234}$ fully separable

$$\omega \in \mathcal{L}(\text{Sym}(\mathcal{H}_{ABC}^{\otimes 4})), \text{tr}_{234}(\omega) = \rho_{ABC}$$

Can we still compute it?

- ▶ Resources:

- a) $Sym((\mathbb{C}^8)^{\otimes 4}) \cong \mathbb{C}^{330} \Rightarrow$ 110.000 parameters
- b) 2 PPT 1|234, 12|34: 960×960 , 1296×1296 matrix

☹ Just slightly too big!

- ▶ Ways to reduce resources:

- ▶ Reduced rank

- e.g. $r(\rho_{ABC}) = 2$ then ω_{1234} effectively $Sym((\mathbb{C}^2)^{\otimes 4}) \cong \mathbb{C}^5$
- Works up to $r(\rho_{ABC}) = 6$: ≈ 15.000 parameters

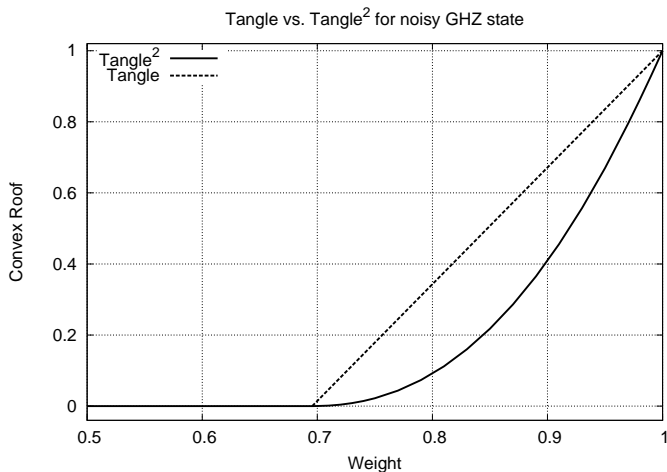
- ▶ Symmetries

$$g_i \rho g_i = \rho, \forall i, g_i^{\otimes 4} T g_i^{\otimes 4} = T \Rightarrow g_i^{\otimes 4} \omega g_i^{\otimes 4} = \omega \quad (21)$$

- GHZ diagonal states: ≈ 10.000 parameters
- Permutationally invariant states: ≈ 20.000 parameters
- Real input $\rho = \rho^T$: parameters/2
- ▶ Are there other solutions [Carathéodory $r(\omega) \leq 64$, normal form, methods to reduce rank,...]?

Example V - Tangle vs. Tangle²

- ▶ GHZ state mixed with white noise
- ▶ Required weight $p \geq 0.6955$ [Eltschka, Siewert PRL (2012)]



Summary

- ▶ Convex roof for polynomial functions f :

- i) $f(|\psi\rangle) = \text{tr}(L |\psi\rangle \langle\psi|^{\otimes n})$

- ii) Equivalent formulation/Lower bound:

$$f(\rho) = \min_{p_i, |\psi_i\rangle} \sum_i p_i f(|\psi_i\rangle) \quad (22)$$

$$= \min \text{tr}(L \omega_{1\dots n}) \quad (23)$$

s.t. $\omega = \omega_{1\dots n}$ fully separable

$$\omega \in \mathcal{L}(\text{Sym}(\mathcal{H}^{\otimes n})), \text{tr}_{2\dots n}(\omega) = \rho$$

- ▶ Examples:

- ▶ Linear entropy [with partial information]
 - ▶ Schmidt number
 - ▶ GHZ vs. W class for 3 qubits

- ▶ Properties:

- ☺ Semidefinite program (certified optimization)
 - ☹ Numerical expensive, in particular for higher polynomials
 - ☺ Employ lots of known techniques from separability questions