

***Particle Detection
in
High-Energy Physics Experiments
Part -II***

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Silicon Detector:

1. Photon spectroscopy with high energy resolution
2. Vertex detection with high spatial resolution
3. Energy measurement of charged particles [few MeV]

Main advantages:

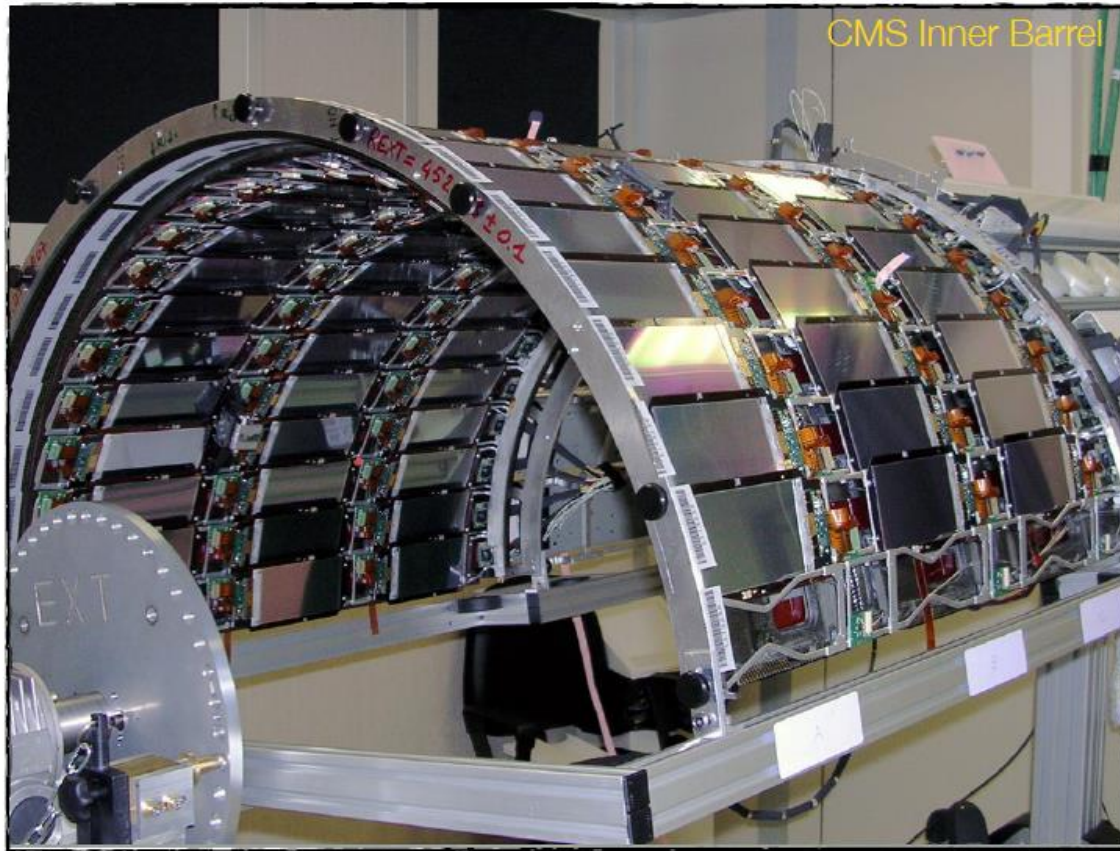
- (i) Possibility to produce small structures using micro-chip technology; 10 μm precision; relatively low costs ...
- (ii) Comparably low energy deposition per detectable electron-hole pair required ...

e.g. Silicon : 3.6 eV per electron-hole pair

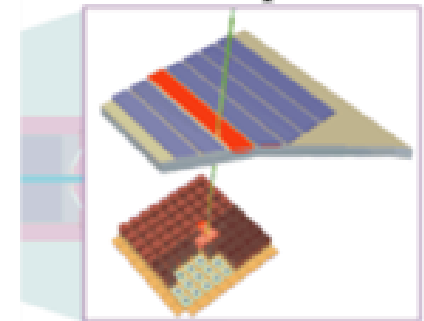
Ionization (LAr): O(30 eV) for a single ion; see later

Scintillators : O(100 eV) depending on light yield [typical 1-10%]

Silicon Tracker



Silicon strip detector



Pixel detector

Basic Semiconductor Detector:

Requirement:

Large sensitive region ...

We know:

$$d \approx x_p \approx \sqrt{\frac{2\epsilon U}{eN_A}}$$

Typical: $N_A = 10^{15}/\text{cm}^3$

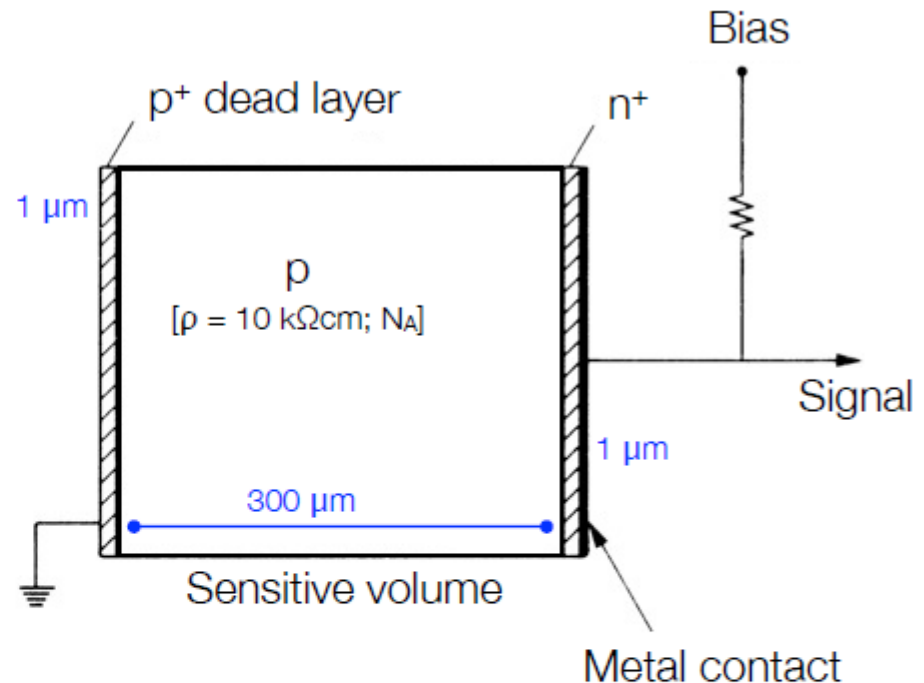
n^+ region highly doped: $N_D \gg N_A$

$$U = \frac{e}{2\epsilon} N_A d^2$$
$$= 100 \text{ V}$$

Electric field:

$$E = \frac{U}{d} = \frac{100 \text{ V}}{300 \cdot 10^{-6} \text{ m}} \approx 3 \cdot 10^5 \frac{\text{V}}{\text{m}}$$

[Safe. Breakdown limit at 10^7 V/m]



n^+ and p^+ needed to
allow metallic contacts ...
[High doping = small depletion zone]

Bias voltage supplied
through series resistor ...

Energy Resolution:

Comparison
of energy resolutions ...

Scintillator [NaI(Tl)]:

1 MeV photon; $\sigma/E \approx 2\%$; $\Delta E/E \approx 5\%$

[$N_i = 40000$ photons/MeV $\times \eta \times$ Q.E.; $\eta = 0.2$, Q.E. = 0.25; $\sigma/E = 1/\sqrt{N_i}$]

Semiconductor [Si]:

1 MeV photon; $\sigma/E \approx 0.06\%$; $\Delta E/E \approx 0.15\%$

[$N_i = 300000$ e/h-pairs/MeV; $\eta \approx 1$, Q.E. ≈ 1 ; $F = 0.1$ $\sigma/E = \sqrt{F}/\sqrt{N_i}$]

Energy resolution of a semiconductor detector
can be better by a factor 25 to 30.

This is indeed observed:
[for $E_\gamma = 1.33$ MeV]

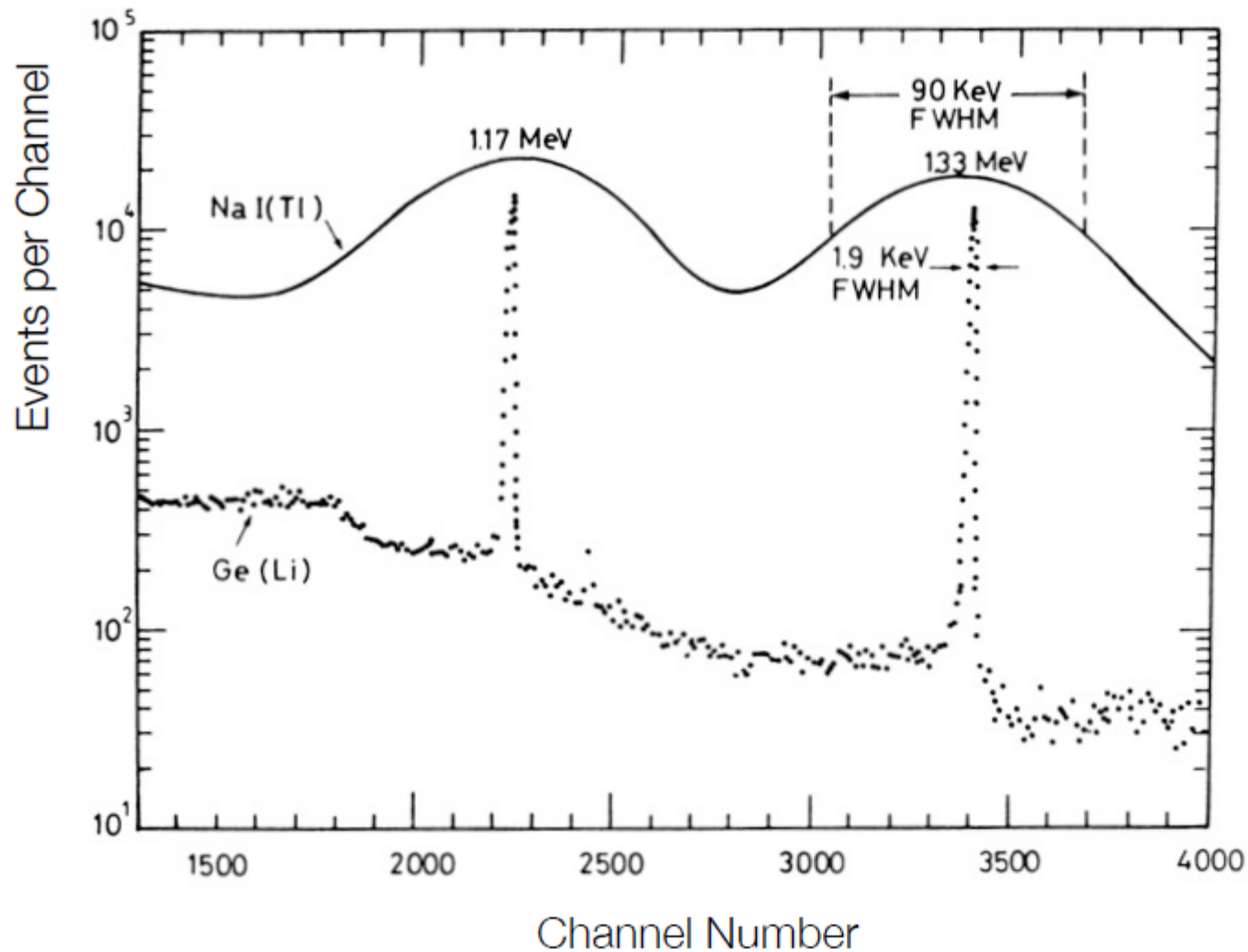
Ge(Li) Counter:

Resolution of 0.15% possible (at ~ 1 MeV)

NaI(Tl) Detector:

Resolution of about 6% (at ~ 1 MeV)

Energy Resolution:



Comparison
of two ^{60}Co γ -spectra
one measured with NaI(Tl)
and With Ge(Li) Detector

Position sensitive detector:

Motivation:

b-Quark tagging & life time measurements
via secondary vertex finding ...

e.g.: $p\bar{p} \rightarrow t\bar{t} + X$ [Tevatron]
 $\hookrightarrow b\bar{b}W^+W^-$

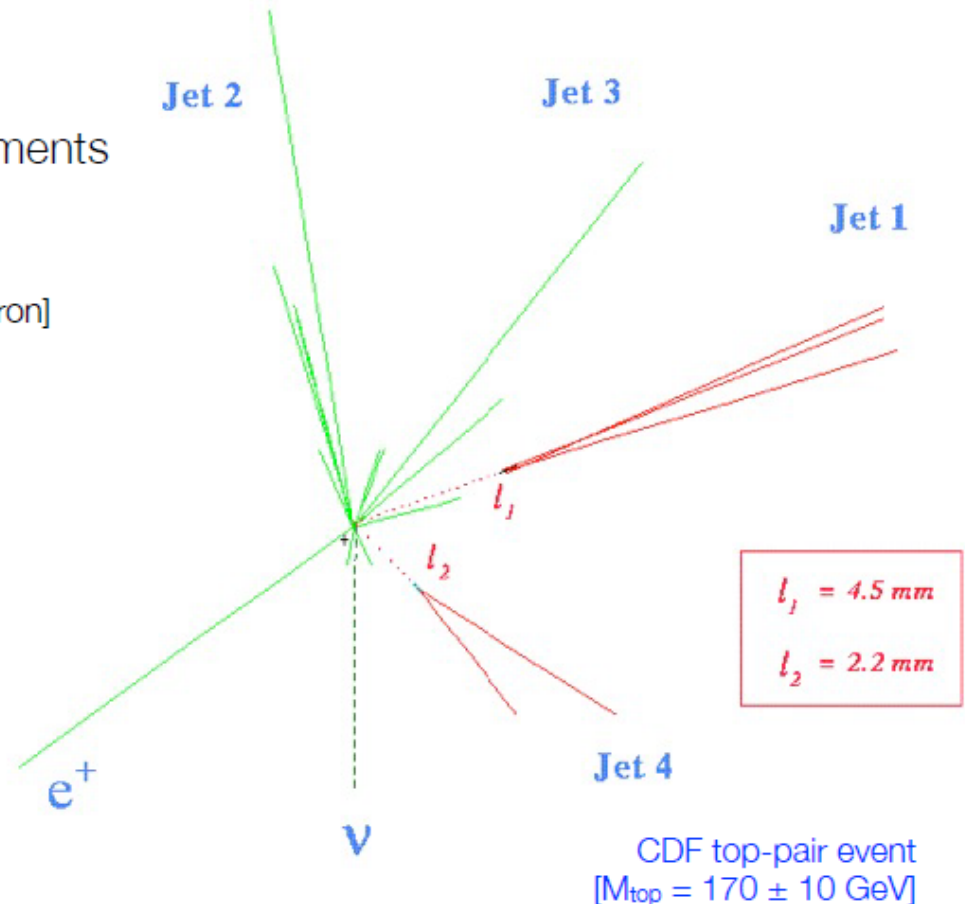
$pp \rightarrow H + X$ [LHC]
 $\hookrightarrow b\bar{b}$

Typical lifetime: $\tau = 10^{-12} \dots 10^{-13}$ s

$$\gamma c\tau = \gamma \cdot 3 \cdot 10^{10} \text{ cm/s} \cdot 10^{-13} \text{ s}$$
$$= \gamma \cdot 30 \text{ } \mu\text{m}$$

Thus:

To measure lifetime in picosecond regime
one needs spacial resolution of the order of 5 - 30 μm ...



Position Sensitive Detector:

Principle:

Segmentation
into strips, pads, pixels ...

Typical parameters:

Thickness: 150 - 500 μm

Strip separation (pitch): 20 - 150 μm

Resolution: 5 - 40 μm (pitch/ $\sqrt{12}$)

Charge collection: 20 ns

Charge integration: 120 ns

Operation voltage: 160 V

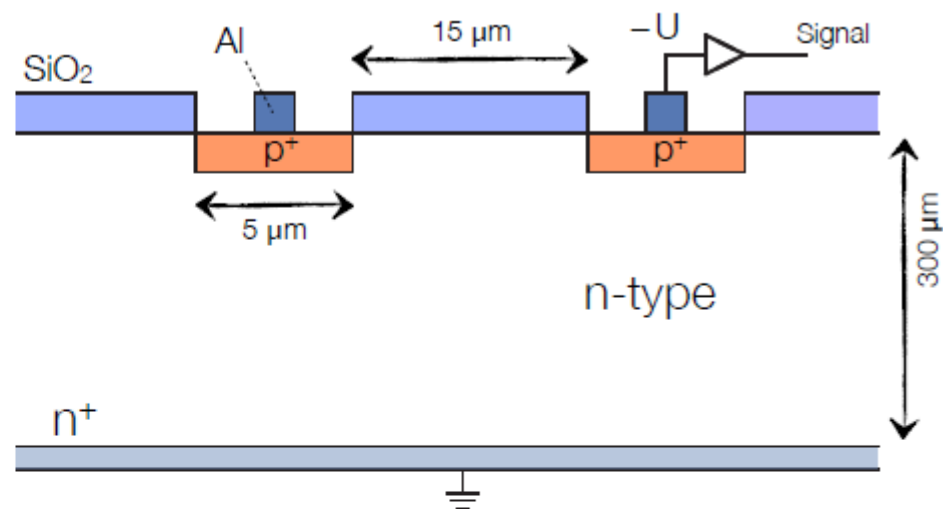
Output signal:

Total charge: $Q_{\text{out}} \sim 4 \text{ fC}$

Average energy loss of MIP: 300 eV/ μm ; Si: 3.6 eV/pair.

Thus 80 electron-hole pairs per μm ;

300 μm thickness $\rightarrow$ 25000 pairs/MIP



Schematics of
Silicon Strip Detector
[from 1983]

High resistive n-type silicon
onto which p+ diode strips with
aluminum contacts are implanted

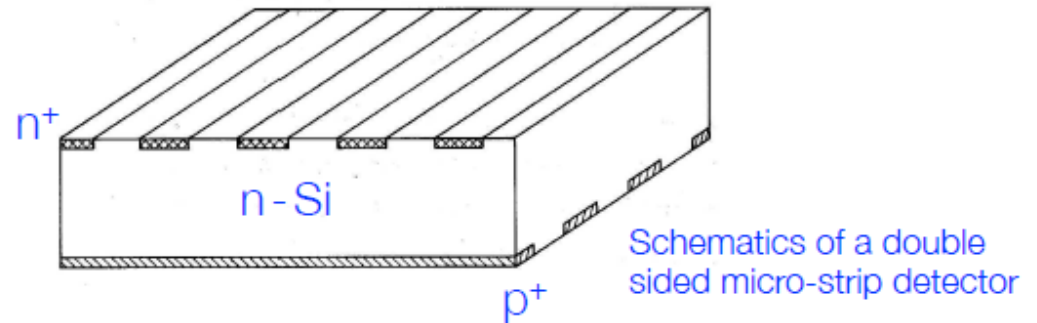
Double sided micro-strip detector:

Next step:

Double sided
micro-strip detectors ...

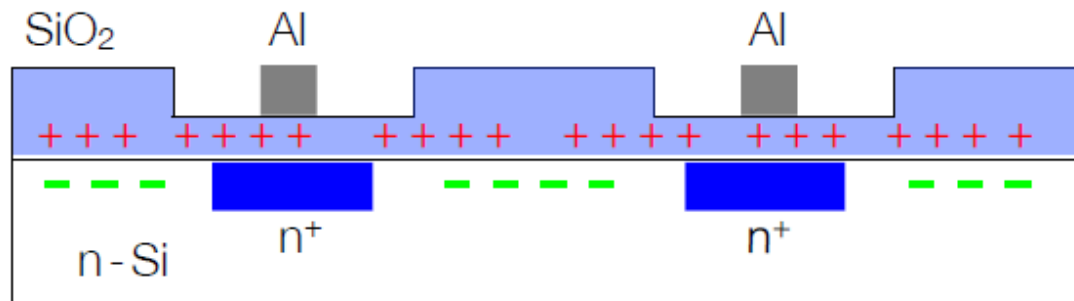
n^+ strips on one side
 p^+ strips on other side

Yields high spatial resolution
in both x and y direction ...



But:

Strips need insulation to avoid that positive
space charge attracts electrons from n -layer



Closeup of of n^+ side of
double sided micro-strip detector

Position Sensitive Detectors:

Pixel detectors:

Like micro-strips, but 2-dim. segmentation ...

Advantage:

As for micro-strips 2-dim. information,
but higher occupancy allowed;

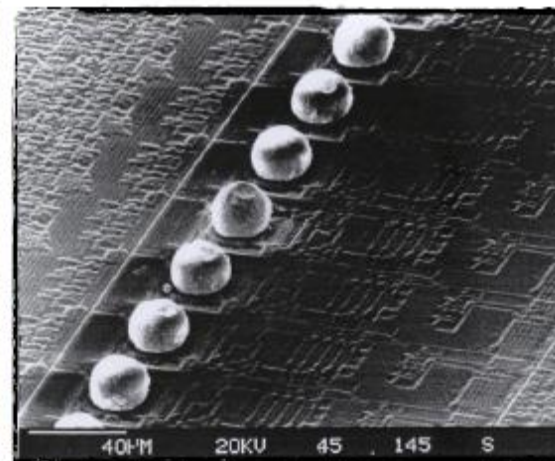
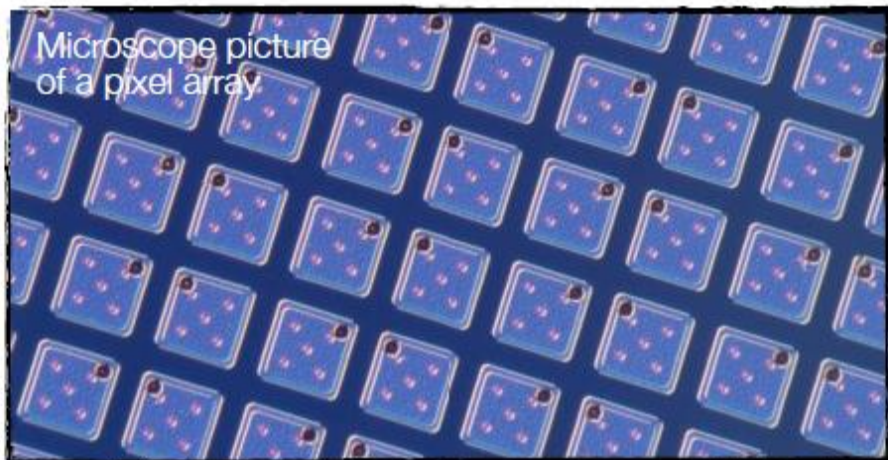
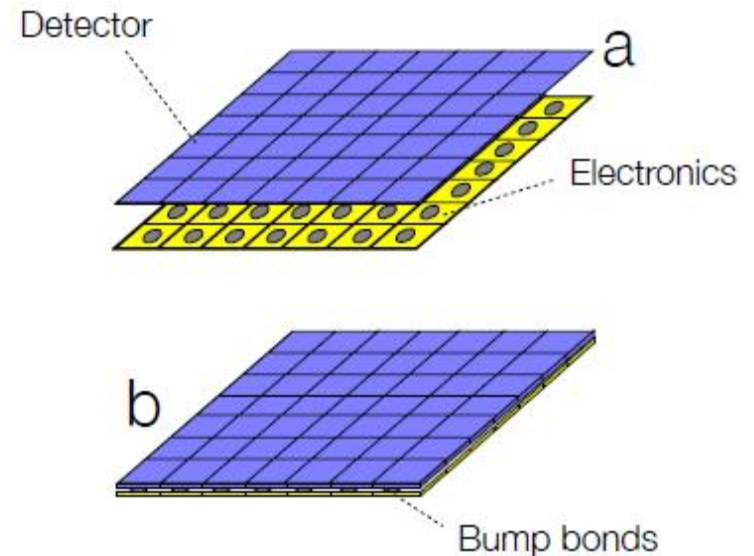
Lower noise due to lower capacitance ...

Disadvantage:

Huge number of readout channels;

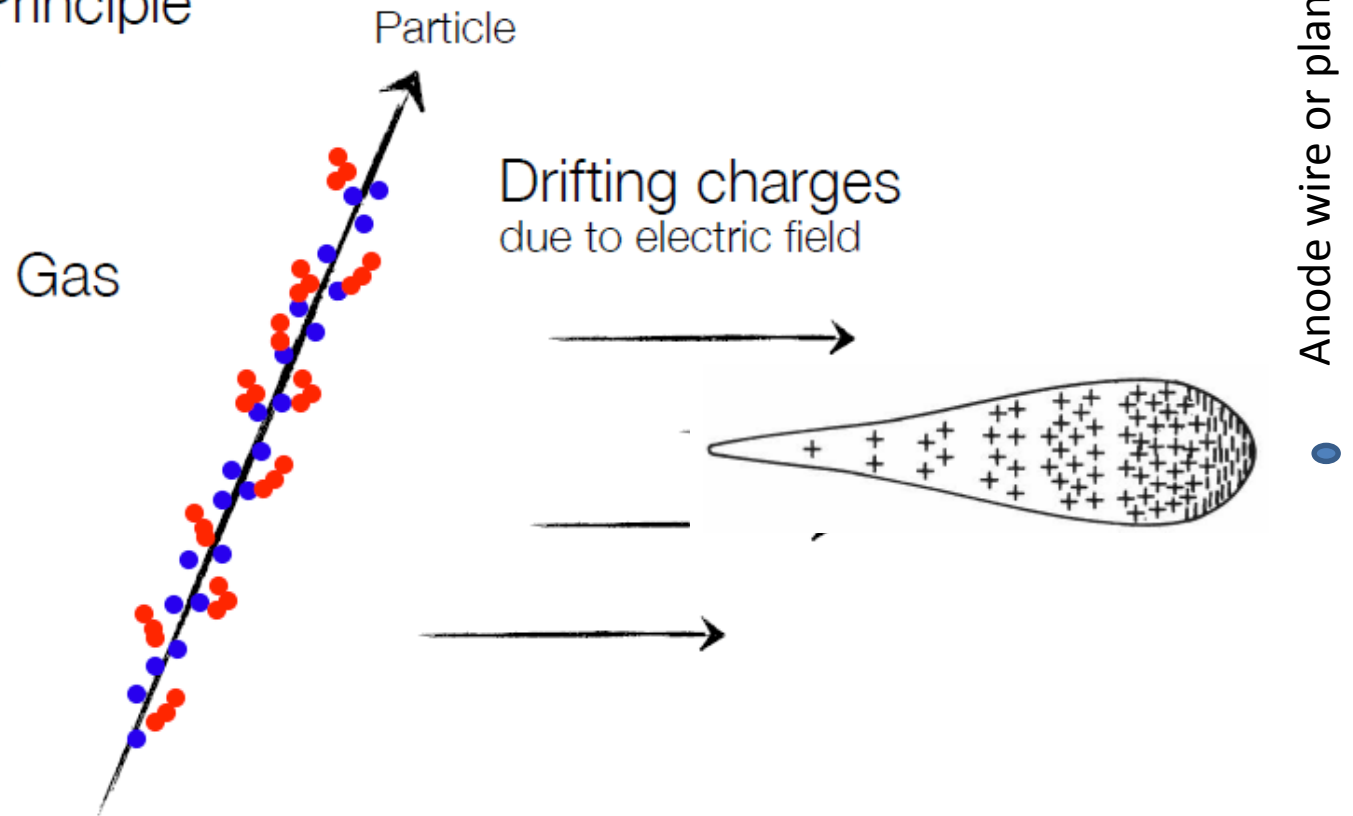
Complicated technology ("bump bonding")

Requires sophisticated readout architecture ...



Gas Detectors:-Basics

Schematic Principle
of gas detectors



- Primary Ionization
- Secondary Ionization (due to δ -electrons)

Gas Detectors:

Ionization mode:

full charge collection
no multiplication; gain ≈ 1

Proportional mode:

multiplication of ionization
signal proportional to ionization
measurement of dE/dx
secondary avalanches need quenching;
gain $\approx 10^4 - 10^5$

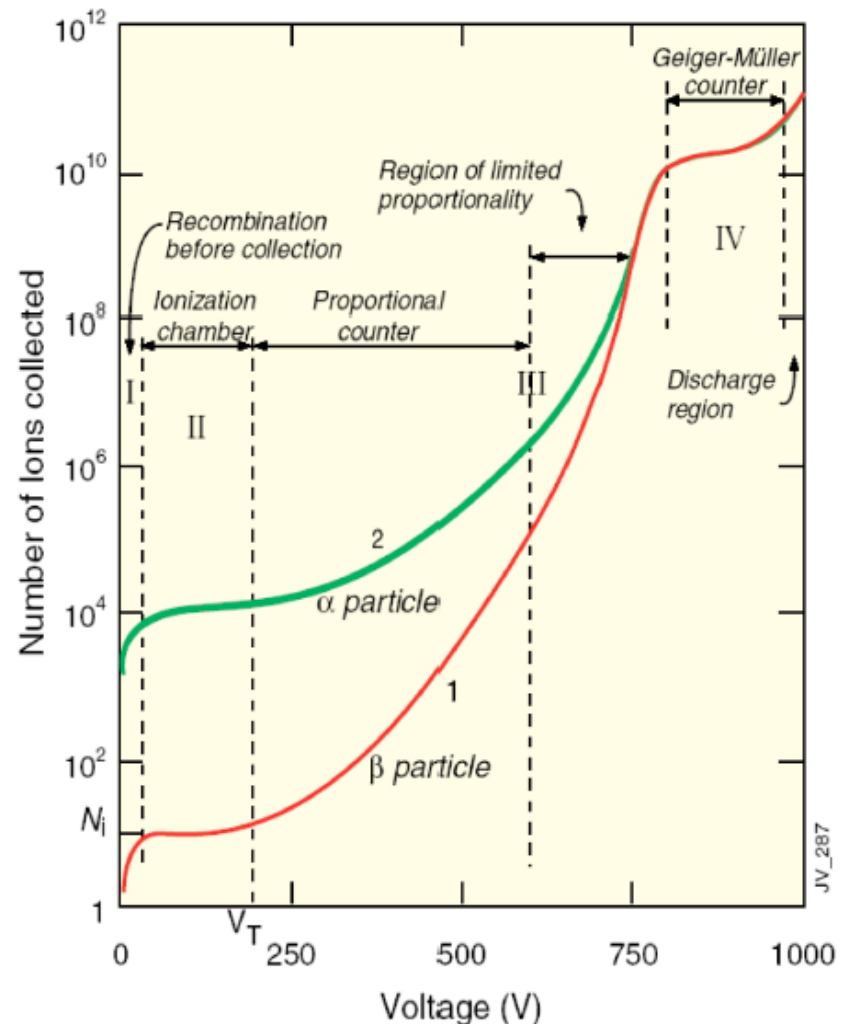
Limited proportional mode:

[saturated, streamer]

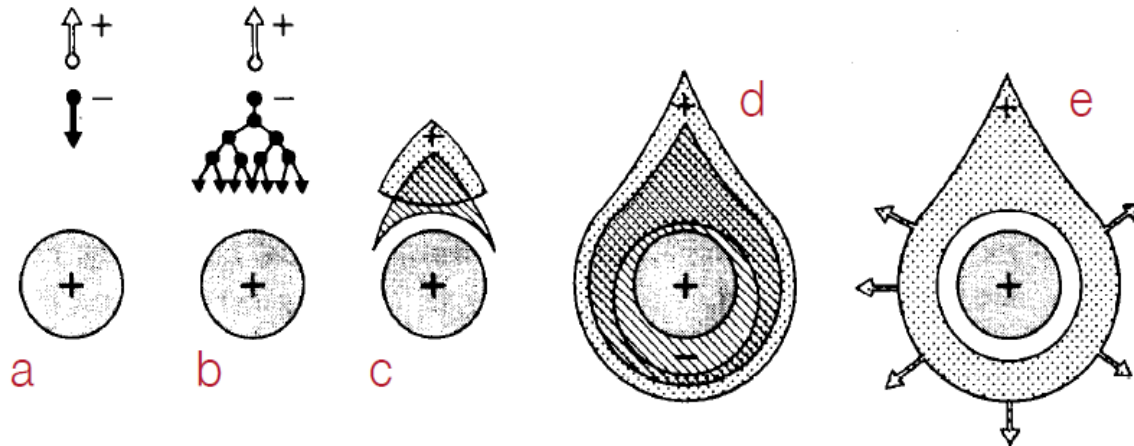
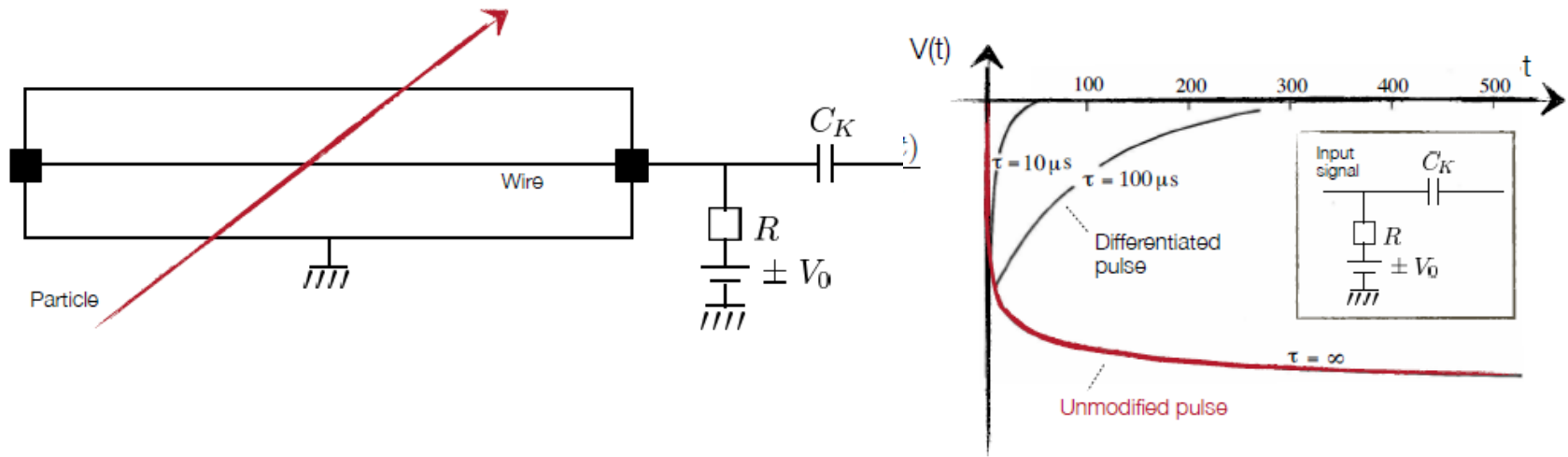
strong photoemission
requires strong quenchers or pulsed HV;
gain $\approx 10^{10}$

Geiger mode:

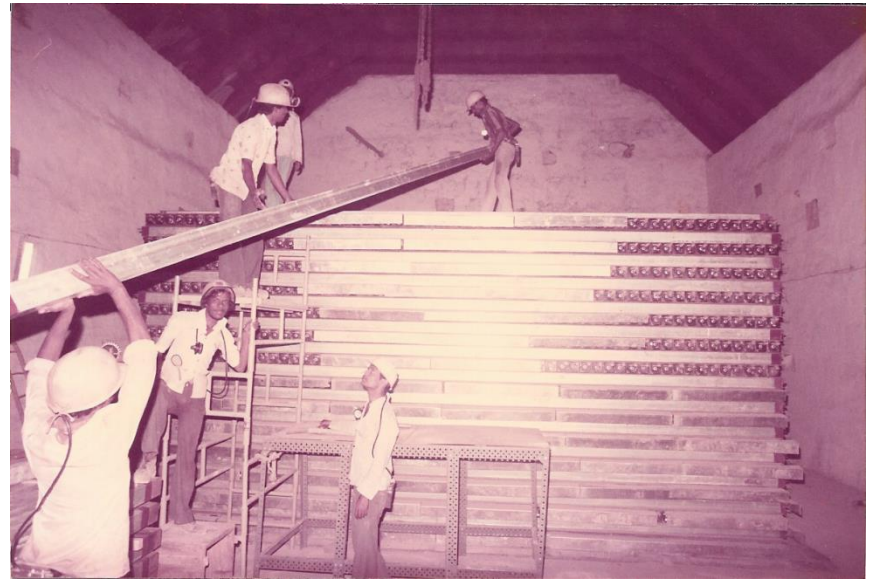
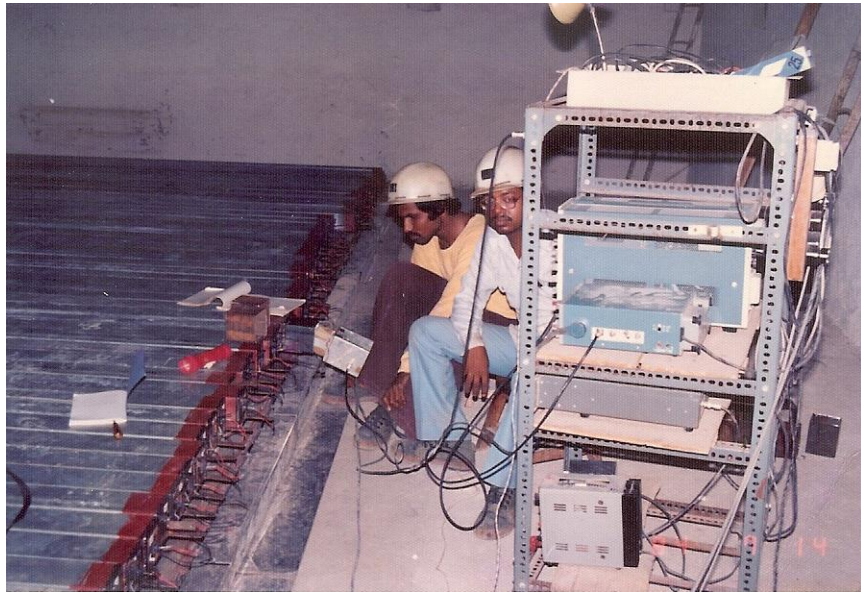
massive photoemission;
full length of the anode wire affected;
discharge stopped by HV cut

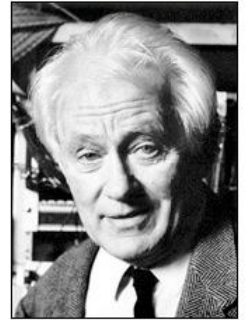


Single Wire Proportional counter



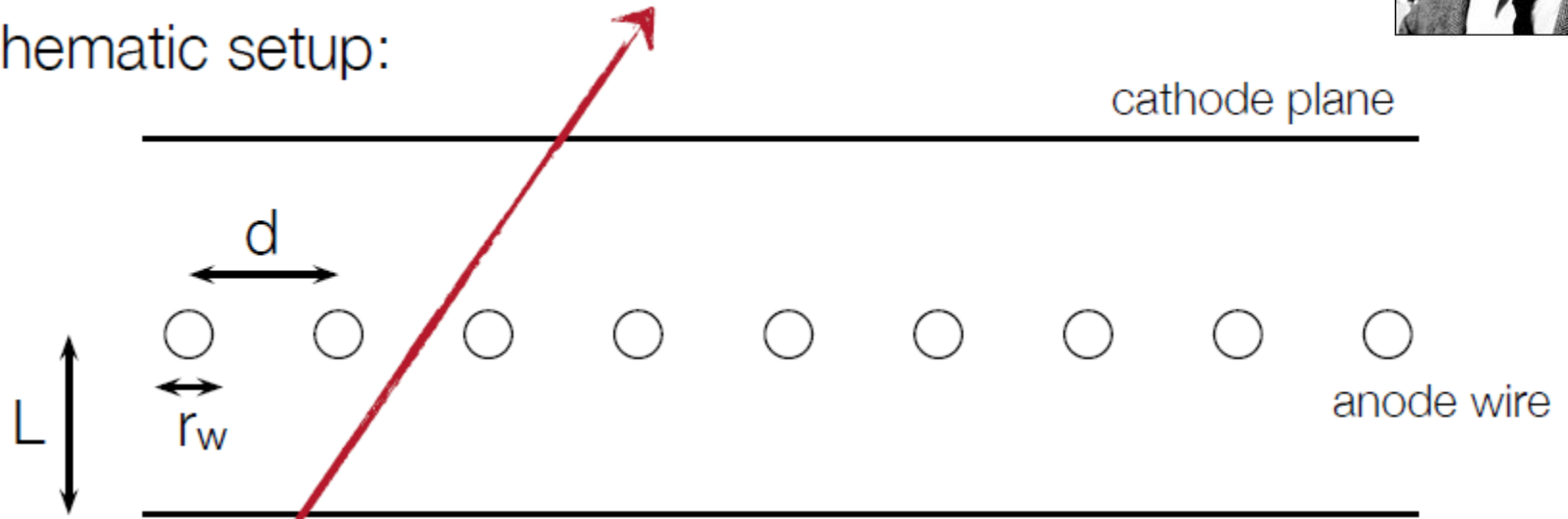
KGF Proportional counters





Multiwire Proportional Chambers

Schematic setup:



Parameters:

- $d = 2 - 4 \text{ mm}$
- $r_w = 20 - 25 \text{ } \mu\text{m}$
- $L = 3 - 6 \text{ mm}$
- $U_0 = \text{several kV}$
- Total area: $O(\text{m}^2)$

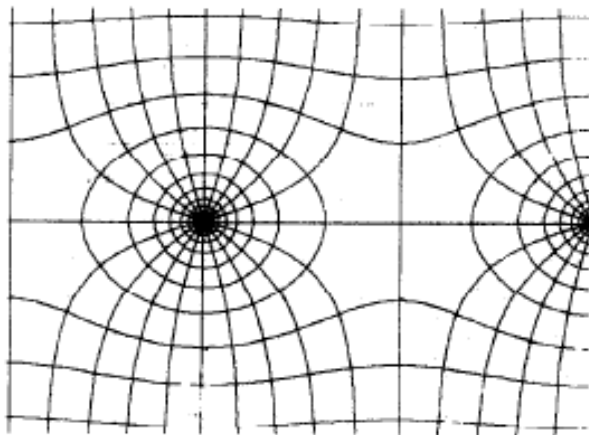
Features:

- Tracking of charged particles
- Some PID capabilities via dE/dx
- Large area coverage
- High rate capabilities

particle track

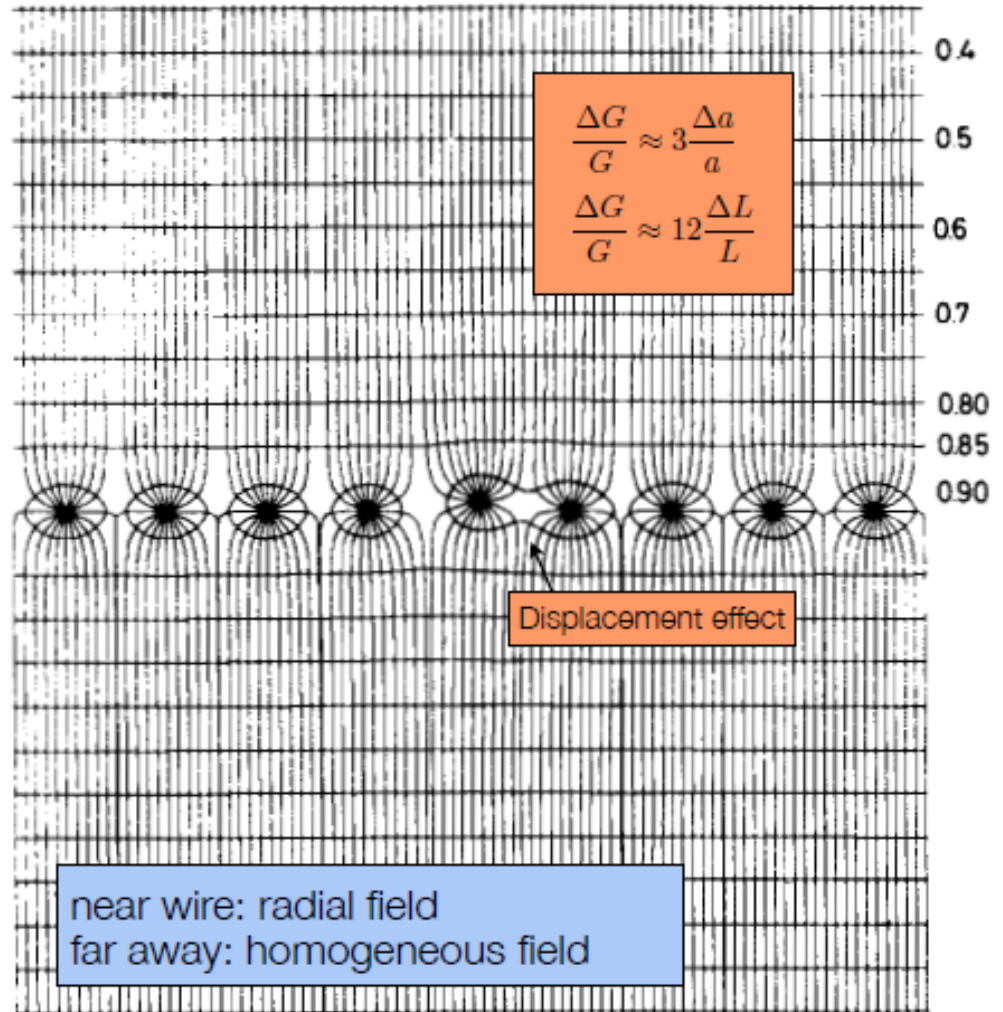
Multiwire Proportional Chambers

Electric field lines
and equipotentials



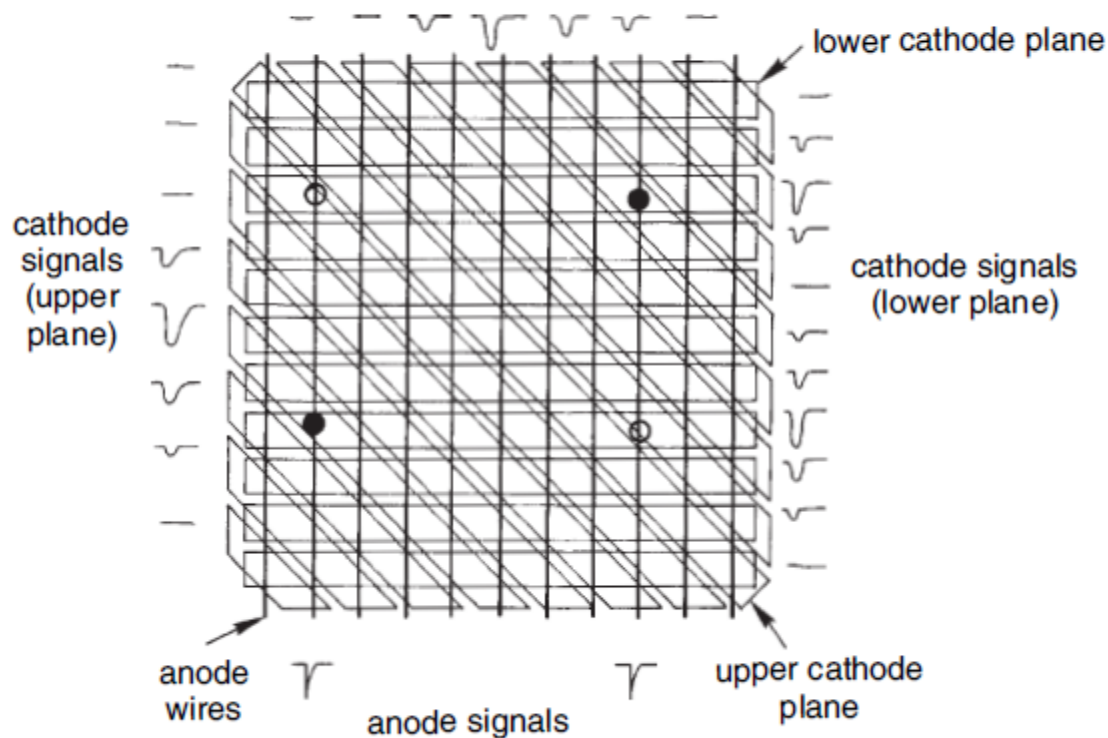
Small wire displacements
reduce field quality ...

Need high mechanical precision
both for geometry and wire tension ...
[electrostatics and gravitation; wire sag]



MWPC:

substantial functionality improvement
due to cathode strips/pads ...



Cathode readout
yields:

2-dim. information
true 2d: use pads ...

high spatial resolution
due to center of gravity reconstruction

resolving ambiguities
using second strip pattern or pads

Can wires be avoided?

Micro-strip Gas Chambers: (MSGC)

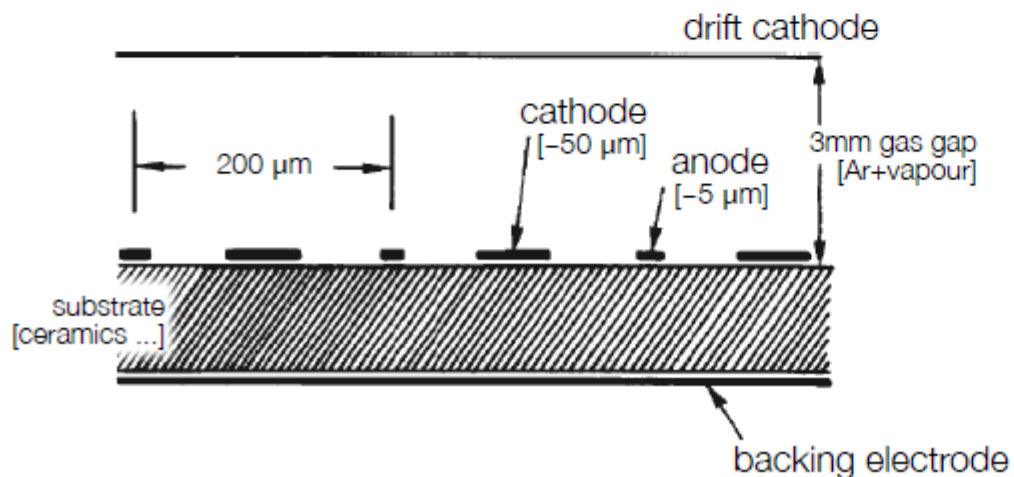
Can one avoid wires?

Anode realized via microstructures
on dielectrics ...

Simple construction (today)
Enhanced stability & flexibility
Improved rate capabilities

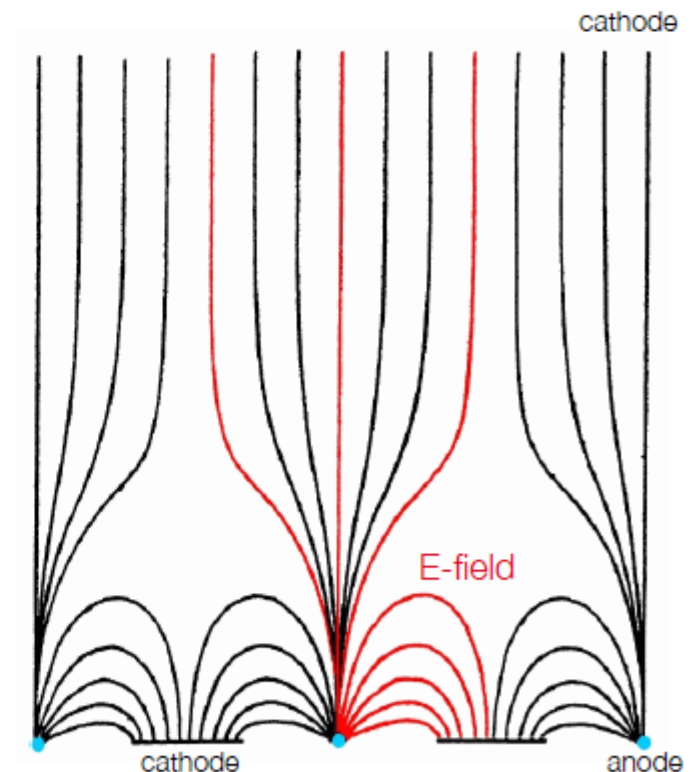
First MSGCs developed in 1990ies ...

Problems: charging of isolation structure
[$\rightarrow$ time-dependent gain; sparks, anode destruction]



Schematics of MSGC
field lines

high field directly above anode
ions drift only $100 \mu\text{m}$; yields low dead time ...



MSGC → GEM

MSGCs prone to aging problems ...
Solution: intermediate grid ...

e.g.: Micromegas
GEM detectors [Sauli, 1997]

Micromegas:

Fine cathode mesh collects ions
still fast; no wires ...

GEM (Gas Electron Multiplier):

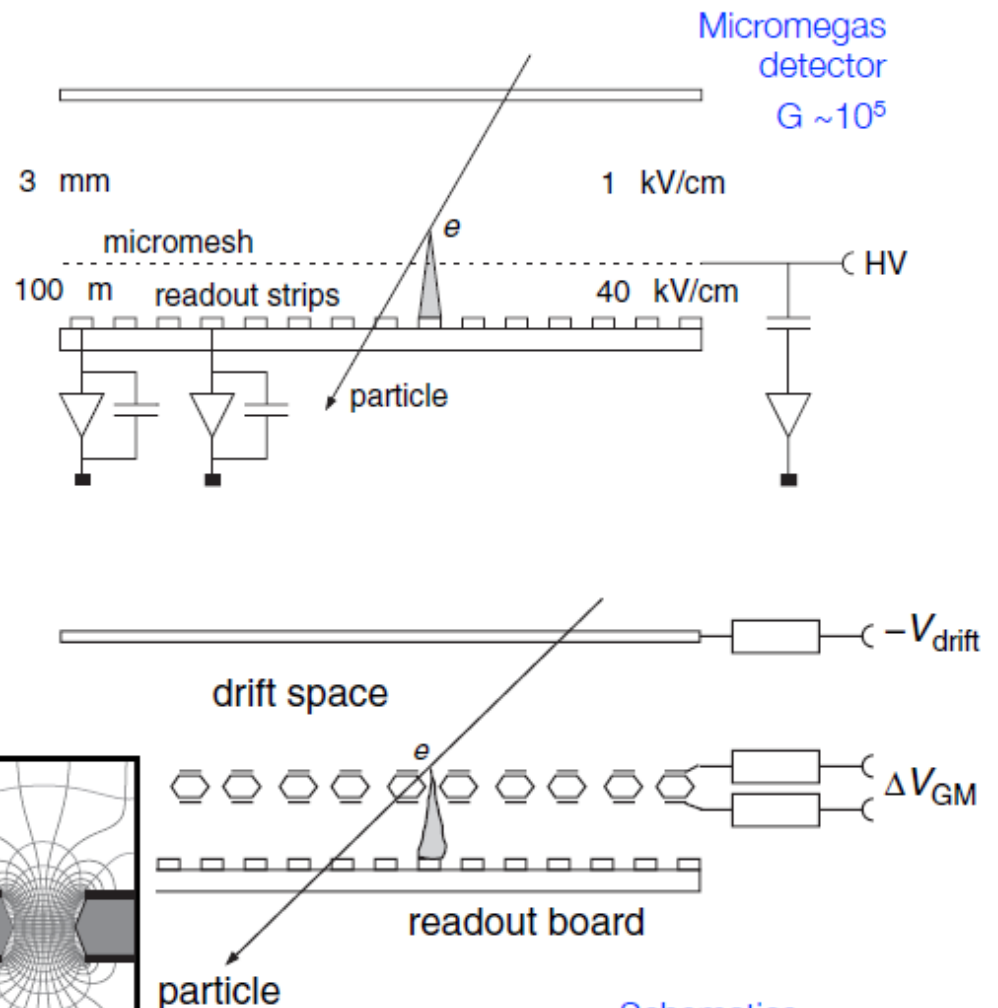
Thin insulating kapton foil
coated with metal film ...

Contains chemically produced
holes [100-200 μm]

Electrons are guided by high
electric drift field of GEMs ...

Avalanche production ...

Electrons drift to anode
GEM collects ions

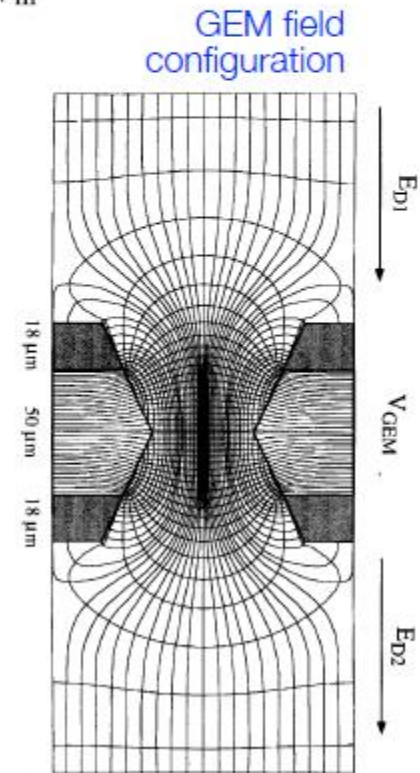
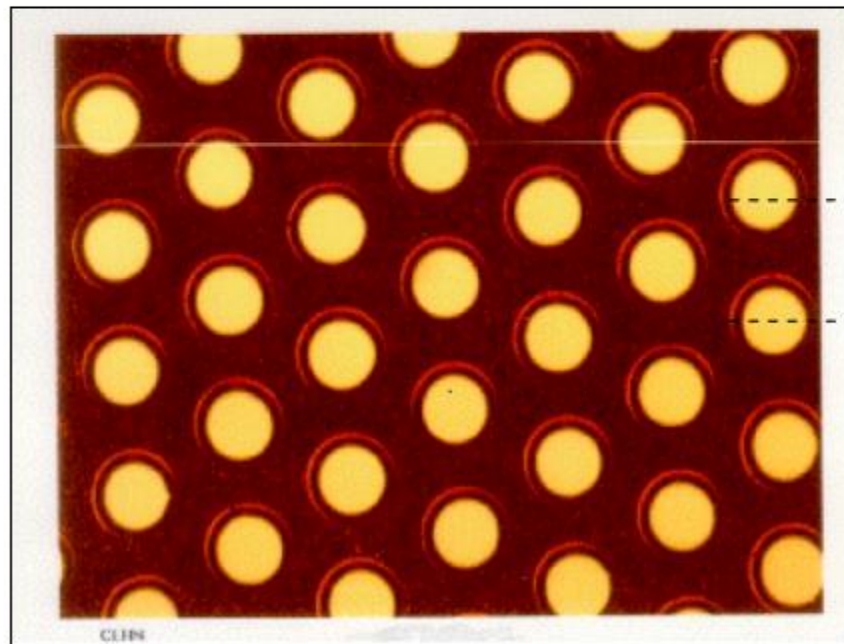
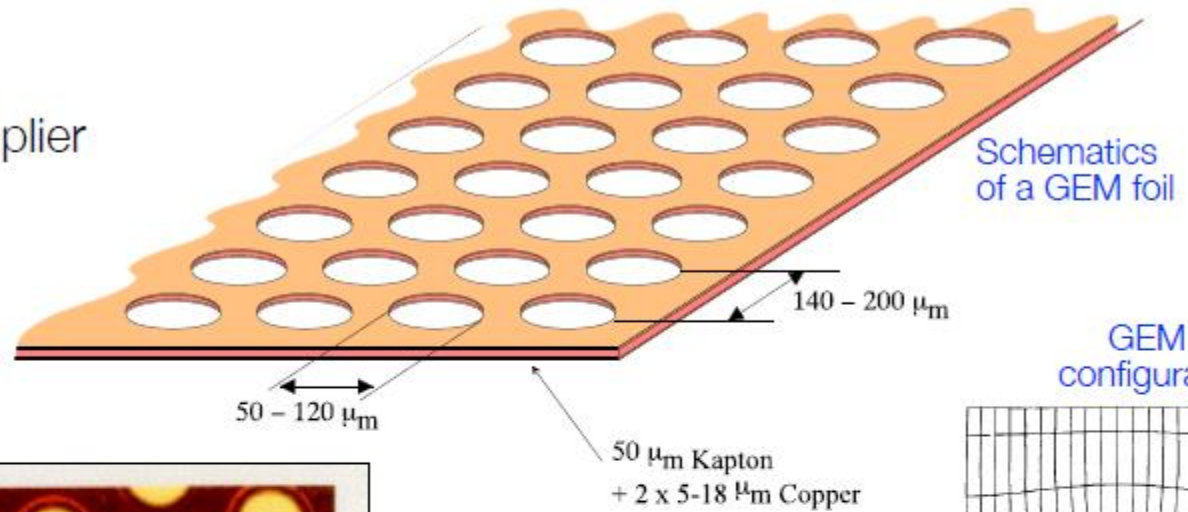


Schematics
of GEM detector

GEM foils:

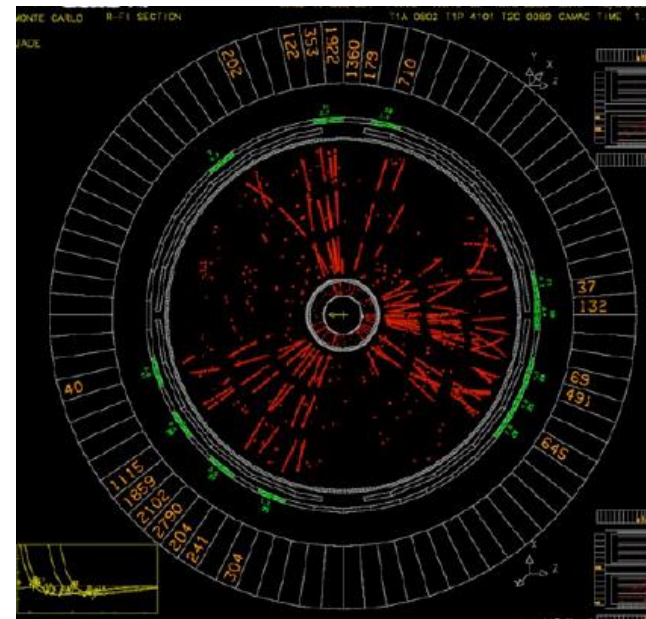
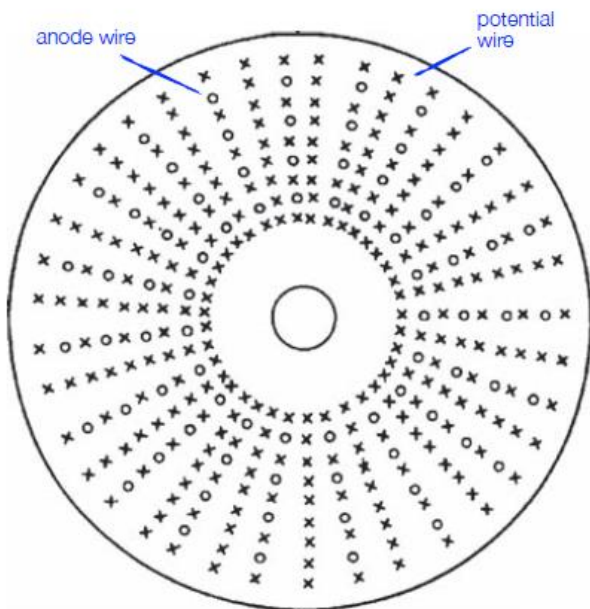
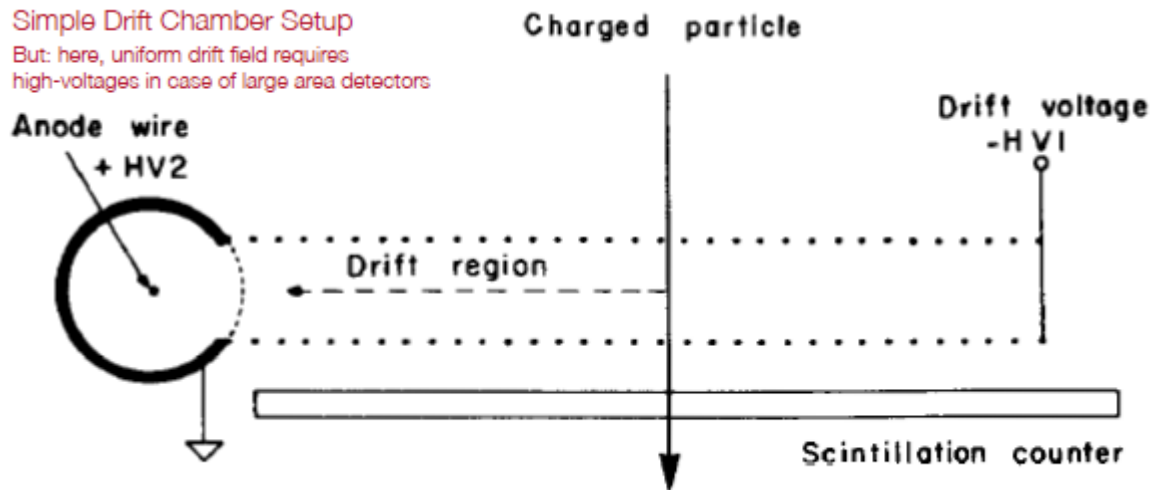
GEM

Gas Electron Multiplier



Drift Chambers:

Simple Drift Chamber Setup
But: here, uniform drift field requires
high-voltages in case of large area detectors



Time Projection Chambers: (TPC)

Electronic 'bubble chamber'
Full 3D reconstruction ...

xy : from wires and pads of MWPC ...
z : from drift time measurement

Momentum measurement ...
space point measurement
plus B field ...

Energy measurement ...
via dE/dx ...

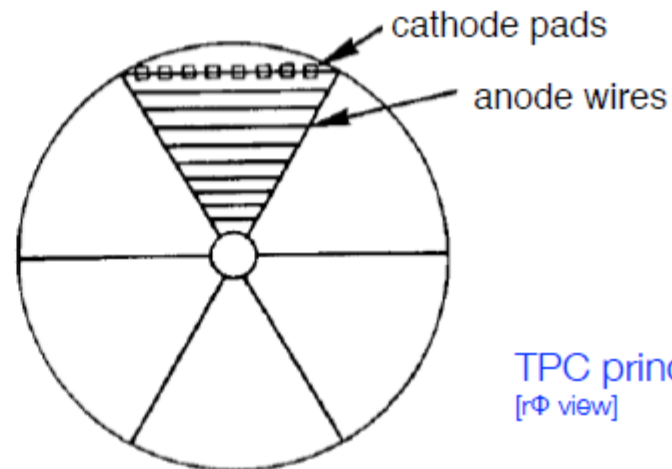
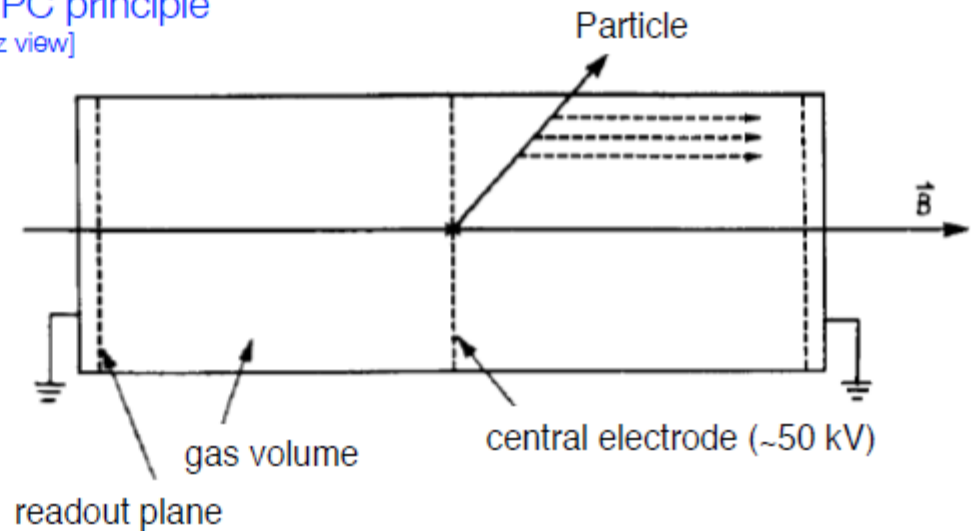
TPC setup:

(mostly) cylindrical detector
central HV cathode
MWPCs at end-caps of cylinder
 $B \parallel$ to $E \rightarrow$ Lorentz angle = 0

Charge transport :

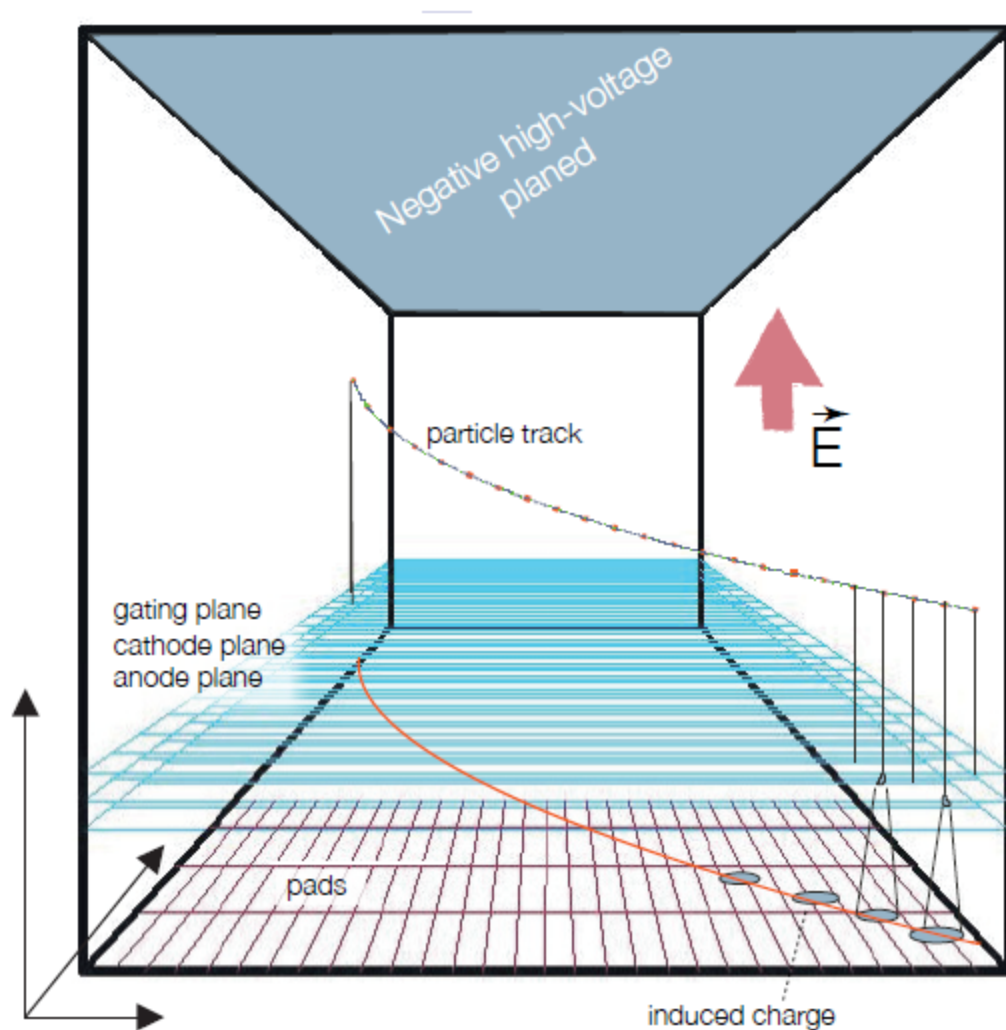
Electrons drift to end-caps
Drift distance several meters
Continuous sampling of induced
charges in MWPC

TPC principle
[rz view]



TPC principle
[rφ view]

TPC:



Advantages:

- Complete track within one detector yields good momentum resolution
- Relative few, short wires (MWPC only)
- Good particle ID via dE/dx
- Drift parallel to B suppresses transverse diffusion by factors 10 to 100

Challenges:

- Long drift time; limited rate capability [attachment, diffusion ...]
- Large volume [precision]
- Large voltages [discharges]
- Large data volume ...
- Extreme load at high luminosity; gating grid opened for triggered events only ...

Typical resolution:

- z : mm; x : 150 - 300 μm ; y : mm
- dE/dx : 5 - 10%

ALICE TPC:

ALICE TPC:

Length: 5 meter
Radius: 2.5 meter
Gas volume: 88 m³

Total drift time: 92 μ s
High voltage: 100 kV

End-cap detectors: 32 m²
Readout pads: 557568
159 samples radially
1000 samples in time

Gas: Ne/CO₂/N₂ (90-10-5)
Low diffusion (cold gas)

Gain: $> 10^4$

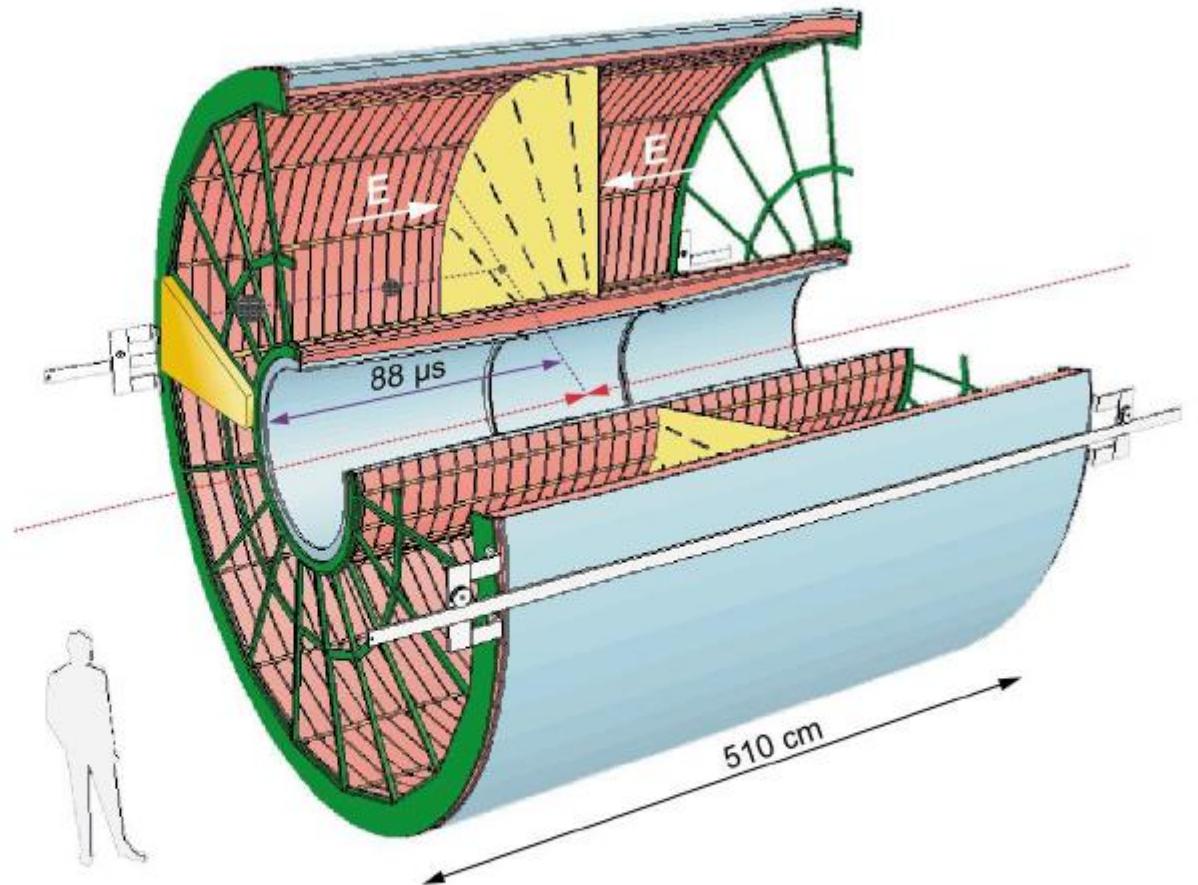
Diffusion: $\sigma_t = 250 \mu\text{m}$
Resolution: $\sigma \approx 0.2 \text{ mm}$

$\sigma_p/p \sim 1\%$ p; $\epsilon \sim 97\%$
 $\sigma_{dE/dx}/(dE/dx) \sim 6\%$

Magnetic field: 0.5 T

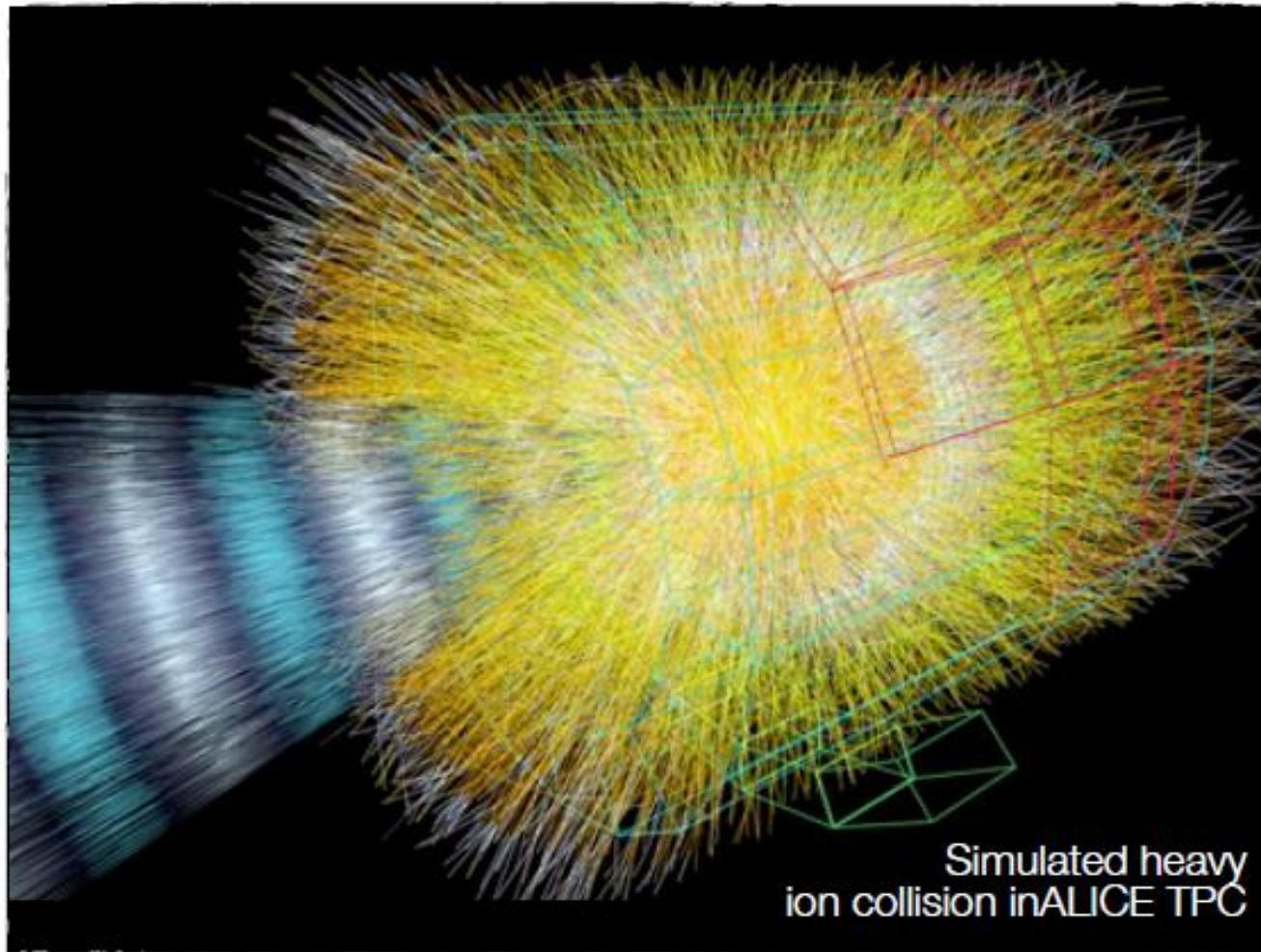
Pad size: 5x7.5 mm² (inner)
6x15 mm² (outer)

Temperature control: 0.1 K
[also resistors ...]



Material: Cylinder build from composite material of airline industry ($X_0 \approx 3\%$)

ALICE TPC:



Calorimeter:

Calorimeter:

Detector for **energy measurement** via total absorption of particles ...

Also: most calorimeters are **position sensitive** to measure energy depositions depending on their location ...

Principle of operation:

Incoming particle initiates particle shower ...

Shower Composition and shower dimensions depend on particle type and detector material ...

Energy deposited in form of: heat, ionization, excitation of atoms, Cherenkov light ...

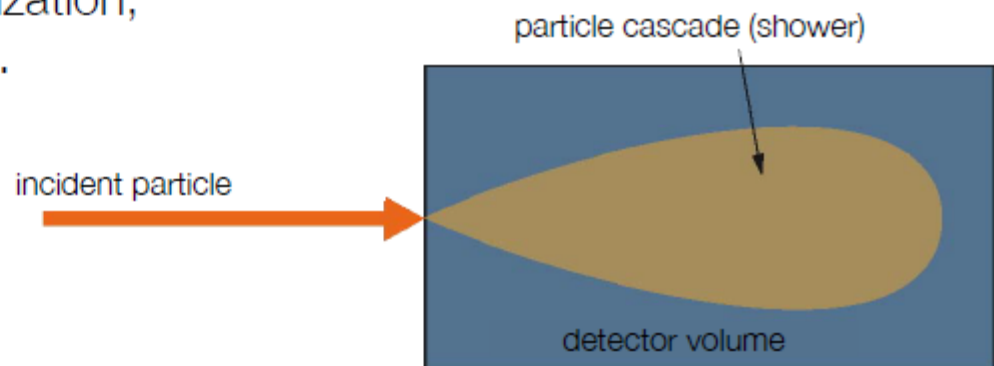
Different calorimeter types use different kinds of these signals to measure total energy ...

Important:

Signal ~ total deposited energy

[Proportionality factor determined by calibration]

Schematic of
calorimeter principle



Calorimeter:

Energy vs. momentum measurement:

Calorimeter:
[see below] $\frac{\sigma_E}{E} \sim \frac{1}{\sqrt{E}}$

Gas detector:
[see above] $\frac{\sigma_p}{p} \sim p$

e.g. ATLAS:

$$\frac{\sigma_E}{E} \approx \frac{0.1}{\sqrt{E}}$$

i.e. $\sigma_E/E = 1\% @ 100 \text{ GeV}$

e.g. ATLAS:

$$\frac{\sigma_p}{p} \approx 5 \cdot 10^{-4} \cdot p_t$$

i.e. $\sigma_p/p = 5\% @ 100 \text{ GeV}$

At very high energies one has to switch to calorimeters because their resolution improves while those of a magnetic spectrometer decreases with E ...

Shower depth:

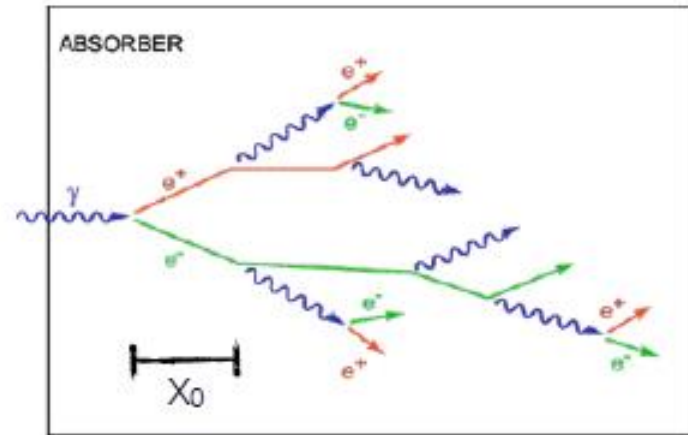
Calorimeter:
[see below] $L \sim \ln \frac{E}{E_c}$
[E_c : critical energy]

Shower depth nearly energy independent
i.e. calorimeters can be compact ...

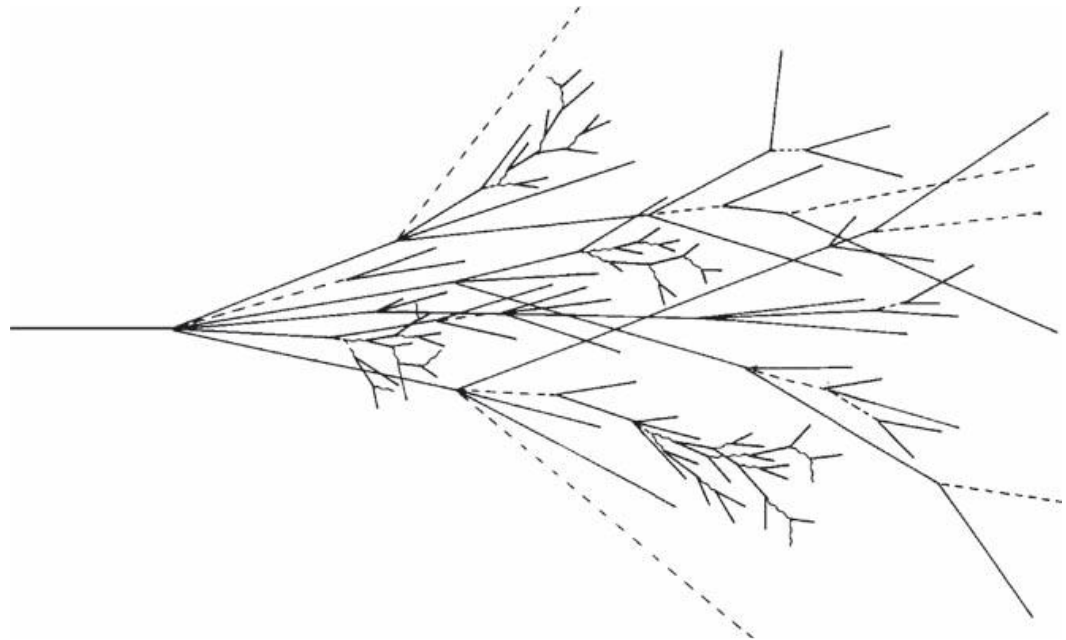
Compare with magnetic spectrometer: $\sigma_p/p \sim p/L^2$
Detector size has to grow quadratically to maintain resolution

Development of Shower

Electromagnetic Shower

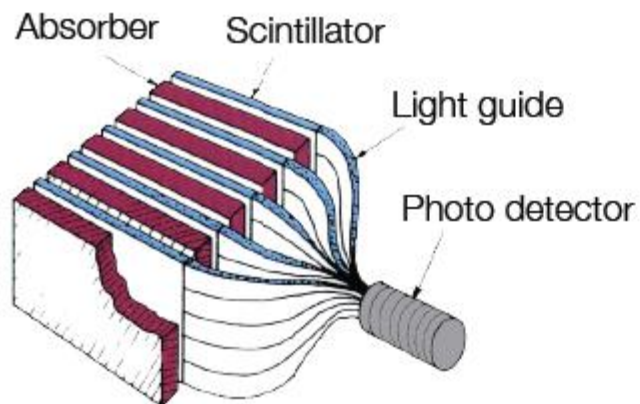


Hadronic Shower

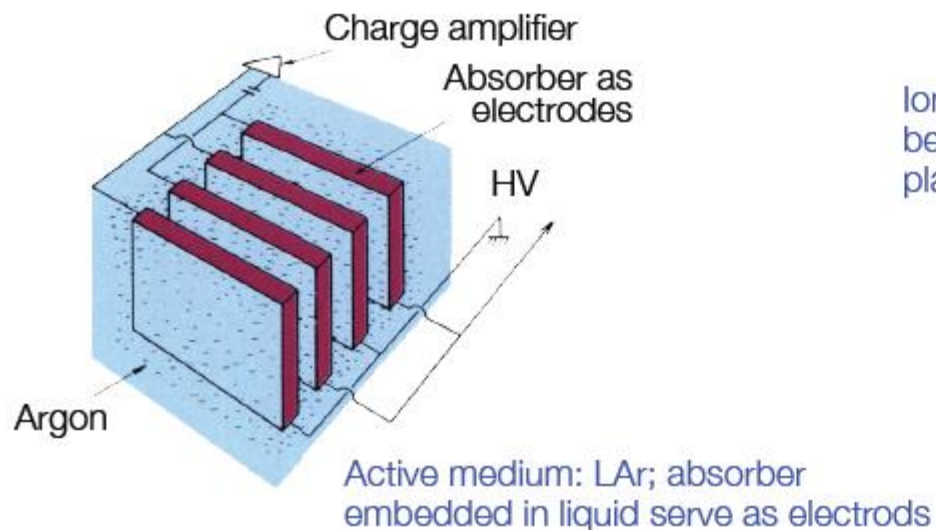
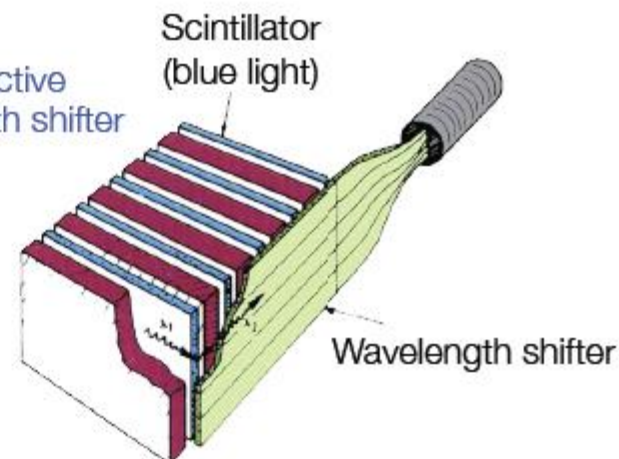


Sampling Calorimeter

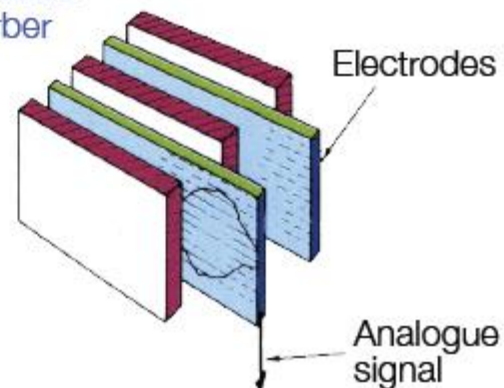
Scintillators as active layer;
signal readout via photo multipliers



Scintillators as active layer;
wave length shifter to convert light



Ionization chambers
between absorber
plates



Typical HEP Calorimeter:

Typical Calorimeter: two components ...

Electromagnetic (EM) +
Hadronic section (Had) ...

Different setups chosen for
optimal energy resolution ...

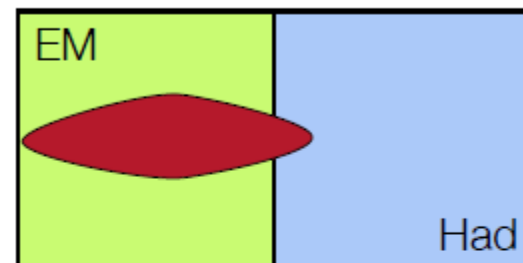
But:

Hadronic energy measured in
both parts of calorimeter ...

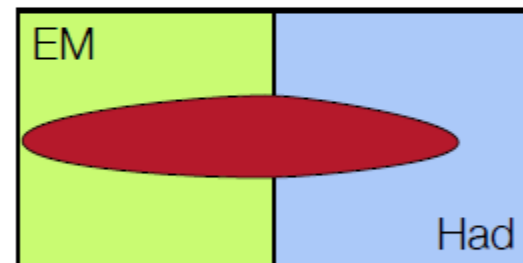
Needs careful consideration of
different response ...

Schematic of a
typical HEP calorimeter

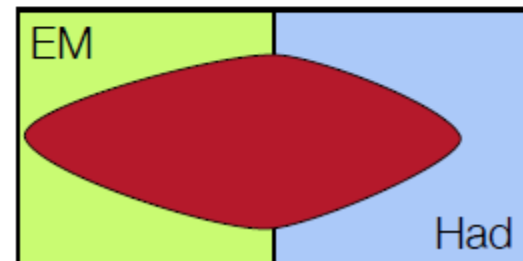
Electrons
Photons



Taus
Hadrons



Jets



Hadronic Energy Resolution:

Energy resolution:

e.g. inhomogeneities
shower leakage

e.g. electronic noise
sampling fraction variations

$$\frac{\sigma_E}{E} = \frac{A}{\sqrt{E}} \oplus B \oplus \frac{C}{E}$$

Fluctuations:

Sampling fluctuations

Leakage fluctuations

Fluctuations of electromagnetic
fraction

Nuclear excitations, fission,
binding energy fluctuations ...

Heavily ionizing particles

Typical:

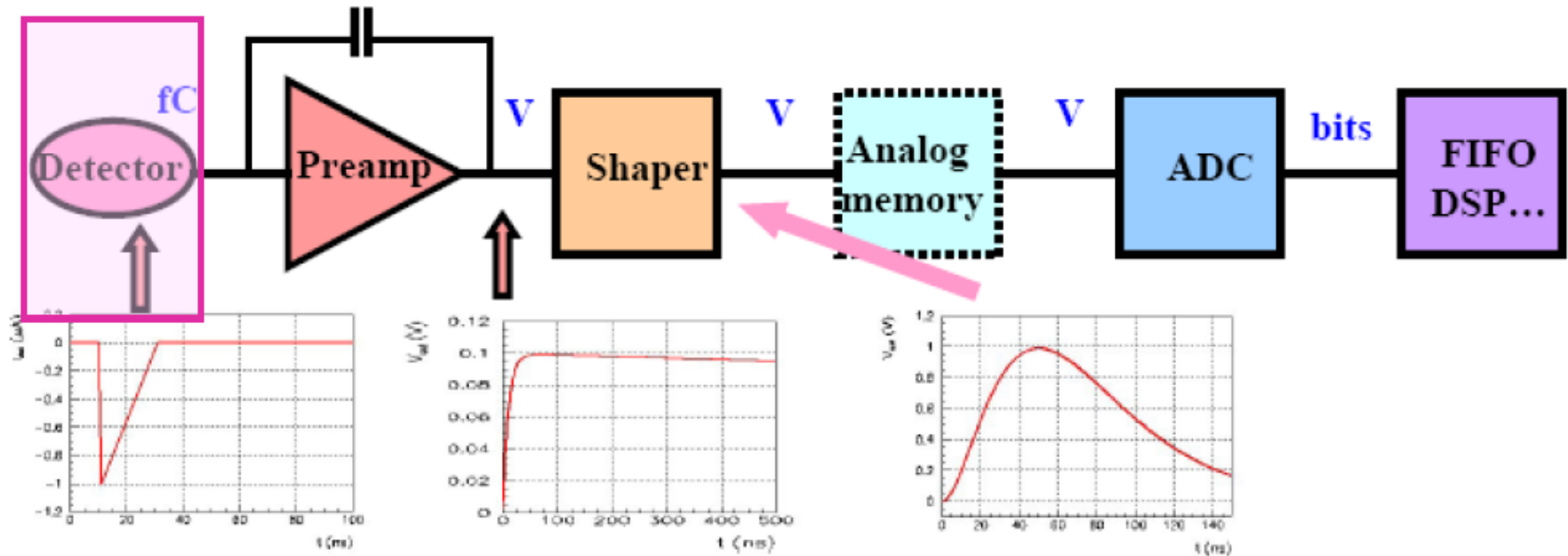
A: 0.5 – 1.0 [Record:0.35]

B: 0.03 – 0.05

C: few %

Path of the signal from Detection to Acquisition

Most front-ends follow a similar architecture



- Very small signals (fC) -> need amplification
- Measurement of amplitude (ADCs)
- Thousands to millions of channels

- **Acknowledgement**
- My thanks to Hans-Christian Schultz-Coulon, Kirchhoff-Institut für Physik for his excellent lecture notes which I have used extensively in my talk here.