

Electroweak Symmetry Breaking & Vacuum Stability

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Recap: Lecture 1

- ▷ The SM Higgs sector at tree level ensures *unitarity* and renormalisability of the theory of massive weak gauge boson
- ▷ Unitarity conditions necessitates the introduction of a physical Higgs field into the spectrum leading to the *Gauge Hierarchy Problem*
- ▷ This is a consequence of absence of symmetry to protect the Higgs mass

Reference: Contino, TASI lecture; 1005.4269

Lecture 1: SM Higgs sector

$$\mathcal{L}_{Higgs} \supset (D_\mu H(x))^\dagger (D^\mu H(x)) + V_0(h(x)) - \frac{y_u}{\sqrt{2}} \bar{Q}_l \tilde{H} u_R - \frac{y_d}{\sqrt{2}} \bar{Q}_l H(x) d_R$$

$$V_0(h) = -\frac{1}{2} m_h^2 h^2 + \frac{\lambda}{4} h^4$$

- ▷ The tree level Lagrangian is Unitary & Renormalizable
- ▷ The tree level relation between the W mass and Z mass reproduced
- ▷ Gives a gauge invariant mass to the massive gauge boson and fermions

Lecture 2: Radiative corrections at 1-loop

$$V_{1-loop}(h) = \sum_{n=0}^{\infty} \frac{h^n}{n!} \Gamma^{(n)}(P_i = 0)$$

C-W effective potential: DR+MS

$$V_{1-loop}(h) = \frac{1}{64\pi^2} \sum_i n_i m_i(h)^4 \left(\log \left[\frac{m_i(h)^2}{\mu^2} \right] - c_i \right)$$

$$n_{h,\chi,W,Z,t} = 1, 3, 6, 3, -12$$

$$c_{h,\chi,t} = 2/3$$

$$c_{W,Z} = 5/6$$

$m_h \rightarrow$ Higgs dependent mass
 $\mu \rightarrow$ regularisation scale

Reference: Quiros; hep-ph/990131

Lecture 2: RG improved C-W potential

$$V(h, \lambda_i, \mu) = V_0(h, \lambda_i, \mu) + V_{1-loop}(h, \lambda_i, \mu)$$

Regularization scale invariance of Green's Function imply:

$$\mu \frac{dV}{d\mu} = \left(\mu \frac{\partial V}{\partial \mu} + \beta_{\lambda_i} \frac{\partial V}{\partial \lambda_i} - \gamma h \frac{\delta}{\delta h} \right) V = 0$$

Solution to RGE:

$$\mu \rightarrow \mu(t) \equiv \mu_0 e^t$$

$$\lambda(t) \rightarrow \lambda_i(t)$$

$$\text{where } \frac{d\lambda_i}{dt} = \beta_{\lambda_i}$$

- ▷ Identify $\mu(t) \rightarrow h(x)$
- ▷ Perturbative expansion of CW-pot
- ▷ Can compute the beta functions directly

$$h(x) \rightarrow h(t, x) \equiv \exp\left(-\int_0^t \gamma(t') dt'\right) h(x)$$

RG improved 1-loop effective Higgs potential

$$V_{SM}(h) \sim \frac{\lambda_{eff}(h)}{4} h^4$$

C-W Potential:

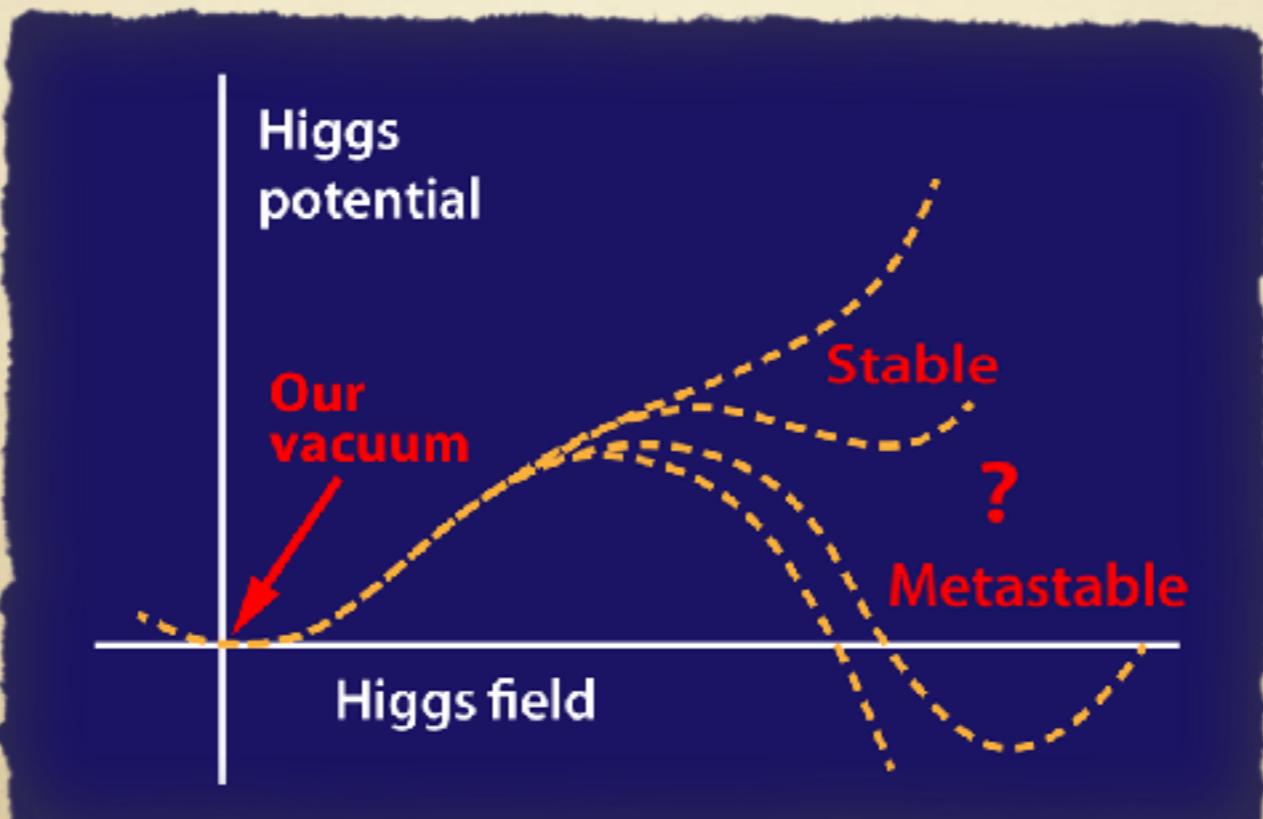
$$\begin{aligned} \lambda_{eff}(h) = & e^{4\Gamma} \left[\frac{\lambda(h)}{4} \text{ -RG improved tree} \right. \\ & + 3 \left(\frac{g_1^2 + g_2^2}{4} \right)^2 \left(\log \left[\frac{g_1^2 + g_2^2}{4} \right] - 5/6 \right) \text{ -Z boson} \\ & + 6 \left(\frac{g_2^2}{2} \right)^2 \left(\log \left[\frac{g_2^2}{2} \right] - 5/6 \right) \text{ -W boson} \\ & - 12 \left(\frac{y_t^2}{2} \right)^2 \left(\log \left[\frac{y_t^2}{2} \right] - 3/2 \right) \text{ -t quark} \\ & + \left(\frac{3\lambda}{2} \right)^2 \left(\log \left[\frac{3\lambda}{2} \right] - 3/2 \right) \text{ -Higgs} \\ & \left. + 3 \left(\frac{\lambda}{2} \right)^2 \left(\log \left[\frac{\lambda}{2} \right] - 3/2 \right) \right] \text{ -Goldstone} \end{aligned}$$

Beta Function:

$$\begin{aligned} \frac{d\lambda(h)}{d \log h} = \beta_\lambda = & \frac{1}{16\pi^2} (12\lambda^2 - 3y_t^4 \\ & + 6y_t^2\lambda + \frac{9}{16} \left[g_2^2 + \frac{2}{3}g_2^2g_1^2 + \frac{3}{25}g_1^4 \right] \\ & + \frac{9}{2}\lambda \left[g_2^2 + \frac{g_1^2}{5} \right]) \end{aligned}$$

SM Vacuum: Stability

- The behaviour of the potential at large field related to Higgs quartic coupling behaviour at high scale
- Depending on the running of the quartic the potential may take various forms:



Metastable Vacuum: Lifetime calculations

WKB Approximation:

$$P \sim A e^{-B}$$

Bounce:

$$B = S_E = \int_{\infty}^{\infty} \mathcal{L}_E$$

Bounce calculation for SM:

$$\mathcal{L}_E = -\frac{1}{2} \left(\frac{\partial h}{\partial \tau} \right)^2 - \frac{1}{2} \left(\frac{\partial h}{\partial x_i} \right)^2 - V(h)$$

E-L Eqn.

$$\frac{\partial^2 h}{\partial \tau^2} + \frac{\partial^2 h}{\partial x_i^2} = \frac{\partial V(h)}{\partial h}$$



$\text{SO}(4)$

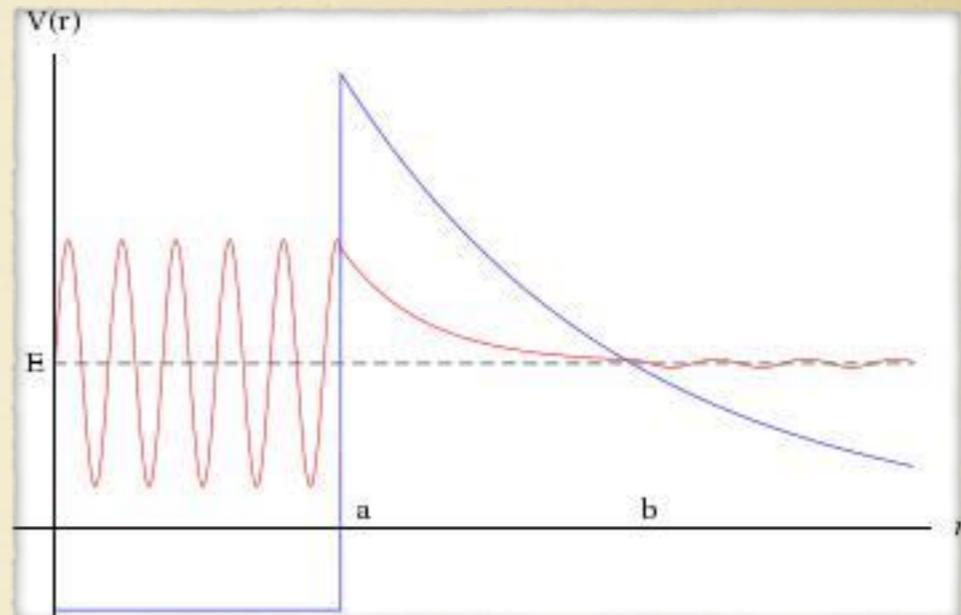
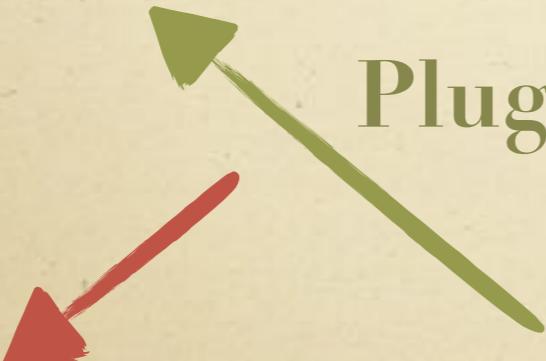
$\rho \equiv \sqrt{\tau^2 + x_i x_i}$ symmetry

Plug back

Solution

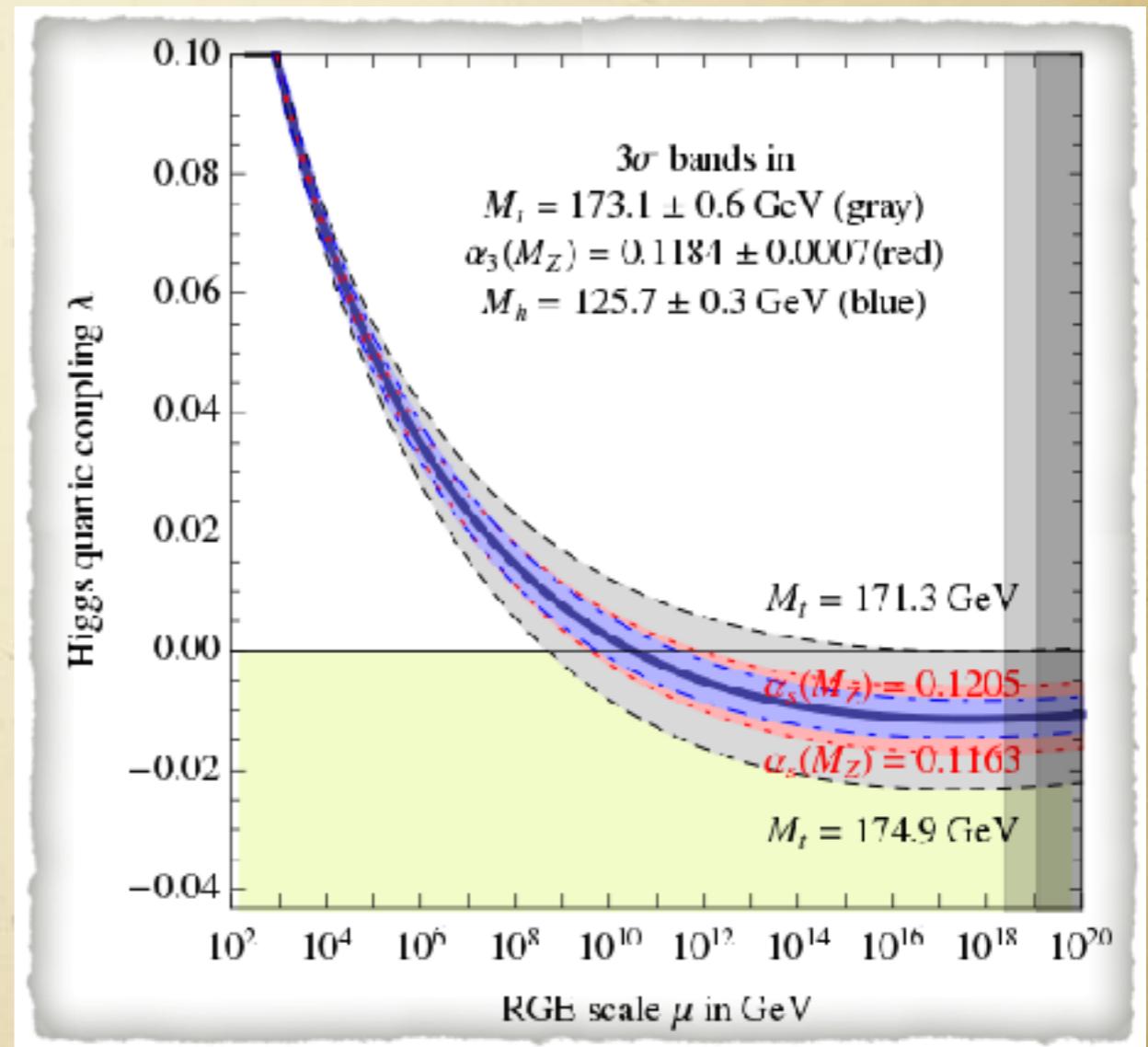
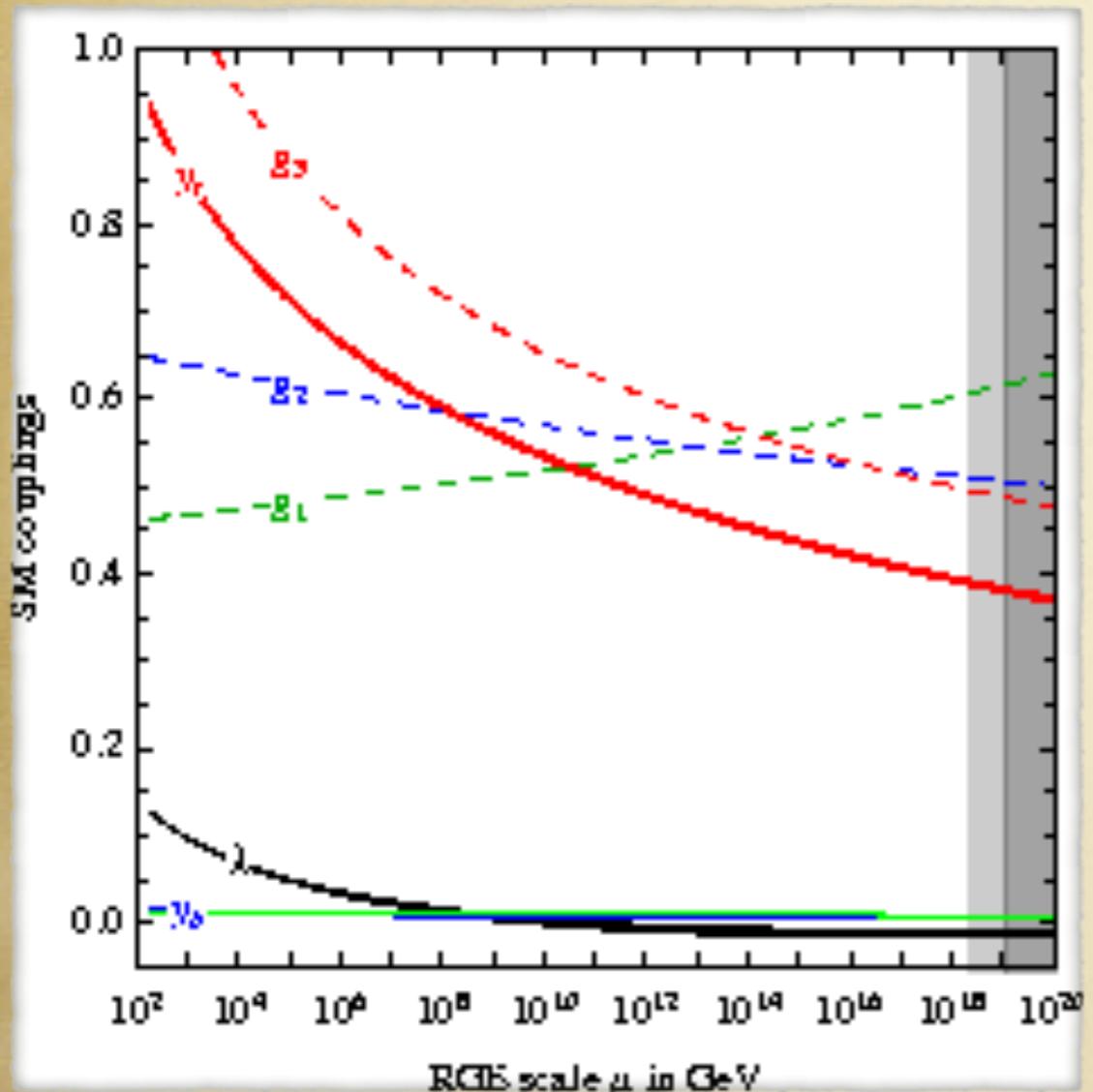
$$\frac{\partial^2 h(\rho)}{\partial \rho^2} + \frac{3}{\rho} \frac{\partial h(\rho)}{\partial \rho} = \lambda h^3(\rho)$$

$$B = S_E = \frac{8\pi^2}{3\lambda}$$



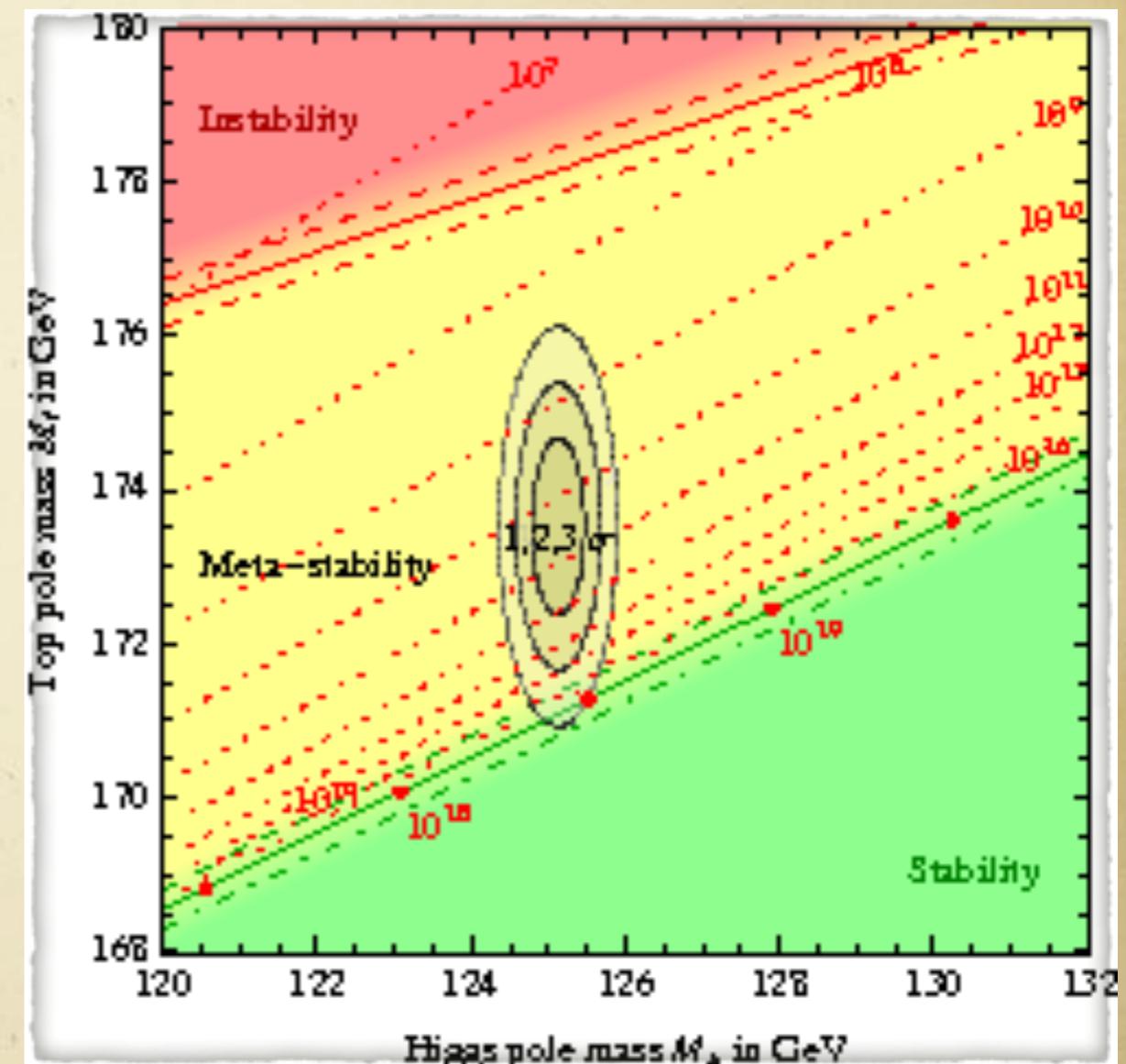
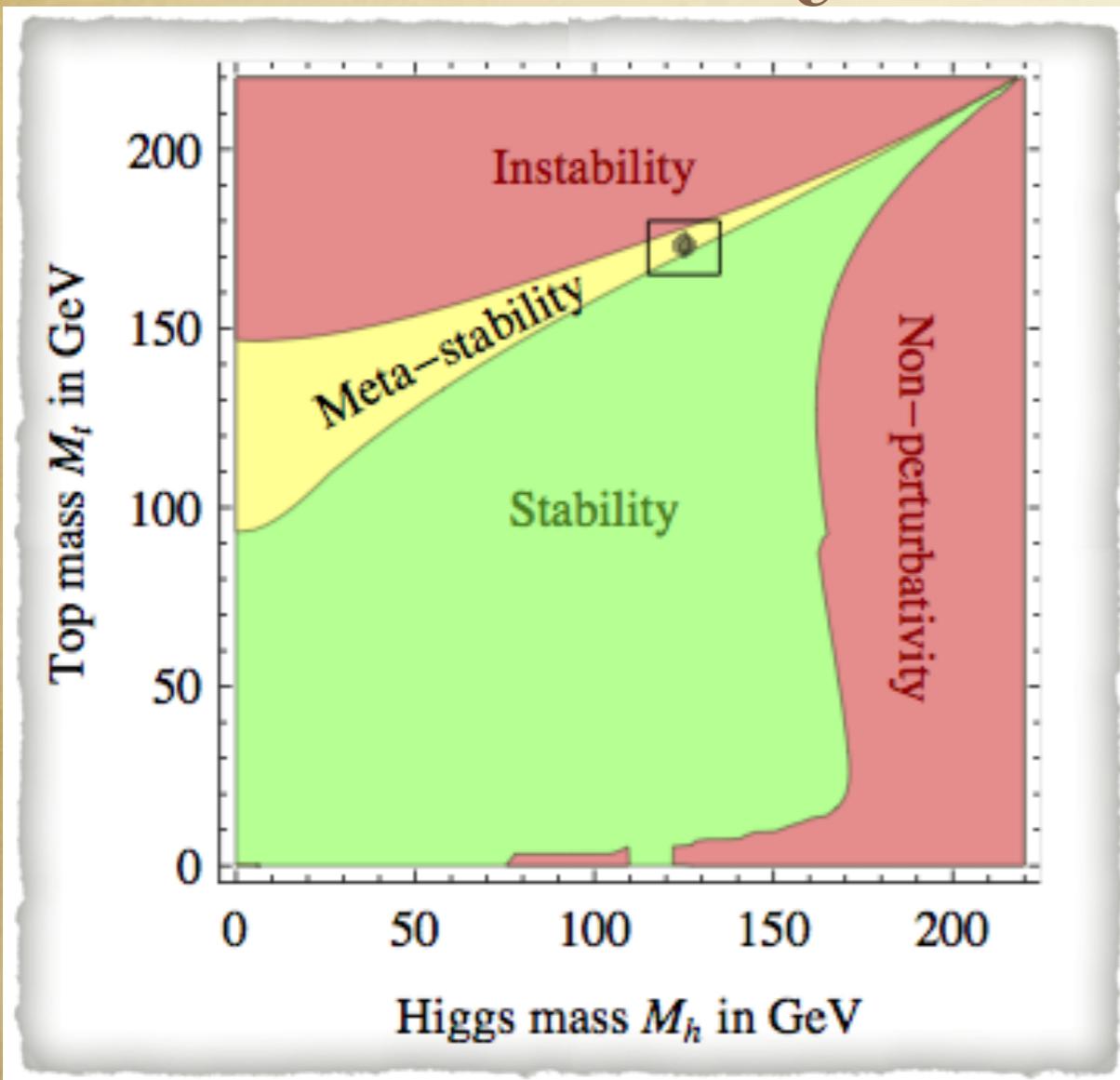
Reference: Coleman; PRD Vol 10 No15, 2929-2936

SM: RG Running of Higgs quartic



Reference: Di Vita et al; arXiv:1205.6497

Stability of SM Vacuum



Near Criticality of Higgs Mass

Reference: Di Vita et al; arXiv:1205.6497

What have we learned from the Higgs discovery: **Criticality**

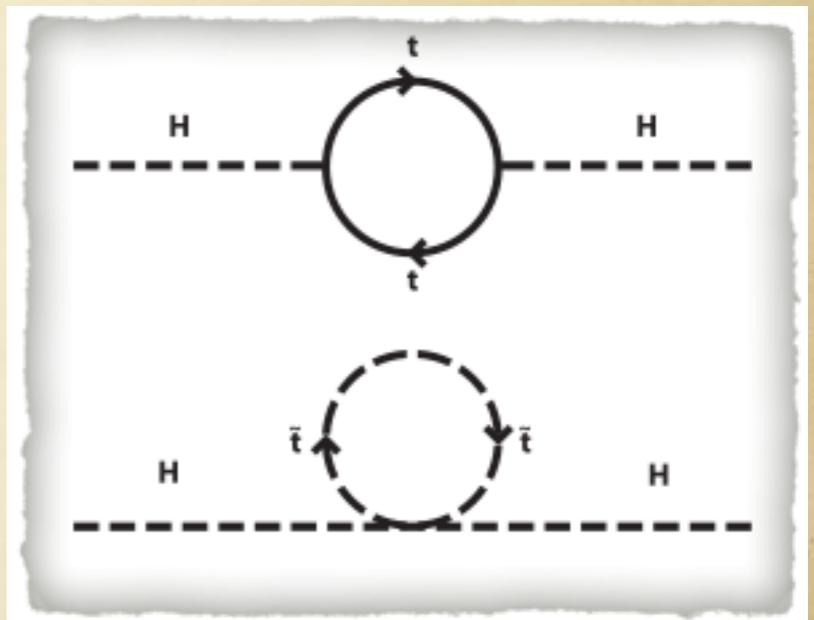
- ▷ Why is the Higgs vev $\ll M_{\text{Planck}}$? : Criticality of the EWSB
- ▷ Why is M_h close to 129 GeV? : Criticality of the Higgs mass

Why are we sitting on the edge?

BSM: Can do away with Criticality: SUSY, Composite Higgs
It can explain criticality dynamically: Relaxion

SUSY

- ▷ Imbibed chiral symmetry from super partner protects the Higgs mass



- ▷ The potential is positive definite. But one can have deeper charge/color breaking minima making the correct EW vacuum metastable

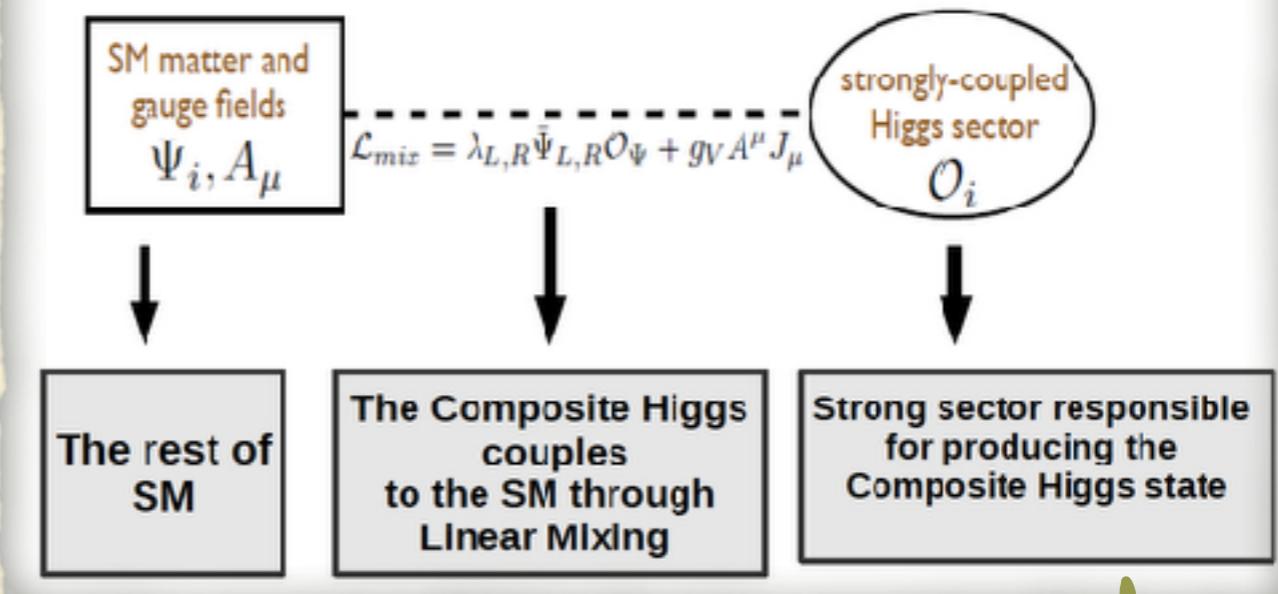
$$V(h) = F^\dagger F + D^\dagger D$$

$$\lambda_h \propto g_1^2 + g_2^2$$

Composite Higgs

- ▷ The inexact shift symmetry of the pNGB Higgs insulates it from the high scale
- ▷ The Higgs “melts” away into constituent state beyond the compositeness scale f where the potential becomes non-existent

The Composite Higgs Program



Dynamics in strong sector produces pNGB mesons out of preons

Above f

Below f

	Sp($2N_c$)	SU(4)
ψ	□	4
M	1	6

Relaxion: Criticality

Extend SM with QCD Axion $a(x)$: pNGB of PQ symmetry (protected from UV instability)

$$\mathcal{L}_{relaxion} \supset (-M^2 + qa(x))h^2 + V(a(x)) + \frac{1}{32\pi^2} \frac{a(x)}{f} \tilde{G}_{\mu\nu} G^{\mu\nu}$$

- At beginning $qa(x) \gg M^2$:
No Higgs vev
- As $q a(x) < M^2$: EWSB;
chiral symmetry broken
- QCD Axion $a(x)$ gets
potential due to feedback:

$$V(a(x))_{QCD} = yv f_\pi^3 \cos(a(x)/f)$$

