

Topological definition:

$$\omega^{(k,q)} \Rightarrow \oint_{\alpha} \langle \vec{z}, \vec{\bar{z}} \rangle$$

$$\equiv \frac{1}{p! q!} A_{\alpha_1 \dots \alpha_p} \bar{A}_{\beta_1 \dots \beta_q} (\vec{z}, \vec{\bar{z}}) d z^{\alpha_1} \wedge \dots \wedge d z^{\alpha_p} \wedge d \bar{z}^{\beta_1} \wedge \dots \wedge d \bar{z}^{\beta_q}$$

$$\bar{\partial} \omega^{(k,q)} = \frac{1}{p! q!} (-1)^k \bar{\partial}_{\beta_0} A_{\alpha_1 \dots \alpha_p} \bar{A}_{\beta_1 \dots \beta_q}$$

$$d z^{\alpha_1} \wedge \dots \wedge d z^{\alpha_p} \wedge \underbrace{d \bar{z}^{\beta_0} \wedge \dots \wedge d \bar{z}^{\beta_q}}_{\text{totally anti-symmetric}}$$

totally anti-symmetric

$$\bar{\partial}\text{-closed } (k,q) \text{ form: } \bar{\partial} \omega^{(k,q)} = 0$$

$$\bar{\partial}\text{-exact } (k,q) \text{ form } \omega^{(k,q)} = \bar{\partial} \omega^{(k,q-1)}$$

$$\bar{\partial}^2 = 0$$

Equivalence relation:

$$\omega^{(k,q)} \sim \tilde{\omega}^{(k,q)} \quad \text{iff}$$

$$\omega^{(k,q)} - \tilde{\omega}^{(k,q)} = \bar{\partial} \omega^{(k,q-1)} \quad \text{for some } \omega^{(k,q-1)}$$

→ Dolbeault cohomology.

Representative elements can be taken to satisfy:

$$(\bar{\partial} * \bar{\partial} * + * \bar{\partial} * \bar{\partial}) \omega^{(p,q)} = 0$$

Similar defn. can be made for cohomology using  $\bar{\partial}$ .

On a Kähler manifold:

$$\begin{aligned} d * d * + * d * d &= 2 (\bar{\partial} * \bar{\partial} * + * \bar{\partial} * \bar{\partial}) \\ &= 2 (d * \bar{\partial} * + * \bar{\partial} * d) \end{aligned}$$

$\Rightarrow$  Same representative elements  $\omega^{(p,q)}$  represent cohomology of  $d$ ,  $\bar{\partial}$  and  $\partial$ .

We can simply solve

$$(d * d * + * d * d) \omega^{(p,q)} = 0.$$

# of such solutions =  $h_{p,q}$ .

Hodge numbers.



Focus on Kähler manifolds from now on.

Consider a Kähler metric:

$$g_{\alpha\bar{\beta}} = \partial_{\alpha} \partial_{\bar{\beta}} K$$

$$J \equiv g_{\alpha\bar{\beta}} dz^{\alpha} \wedge d\bar{z}^{\beta}$$

$$\bar{\partial} J = 0. \quad (\text{also } \partial J = 0)$$

$$\Rightarrow \text{a } (1,1) \text{ form } \in \mathcal{H}^{1,1}.$$

↓  
Kähler class.

Conversely, given an element of the Kähler class we can define a Kähler metric.

$\Rightarrow$  Kähler moduli space is  $h_{1,1}$  dimensional.

(For  $\downarrow$  given complex structure)

Is there a similar understanding for the complex structure moduli space?

## Calabi's conjecture & Yau's theorem:

~~On~~ On a Kähler manifold we can introduce another  $(1,1)$  form

$$\Omega^{(1,1)} = R_{\alpha\bar{\beta}} dz^\alpha d\bar{z}^\beta.$$

$\bar{\partial}$  closed (also  $\partial$ -closed).

$\Rightarrow$  an element of  $\mathcal{H}^{1,1}$

$\downarrow$   
First Chern class.

If  $\Omega^{(1,1)}$  is a trivial element of  $\mathcal{H}^{(1,1)}$  i.e.  $\Omega^{(1,1)} = \bar{\partial} \wedge^{(1,0)}$  for some  $(1,0)$ -form  $\wedge^{(1,0)}$ , then we say that it has vanishing first Chern class.

Calabi-Yau: A Kähler manifold with vanishing first Chern class

admits a Ricci-flat metric (a metric of  $SU(n)$  holonomy).



## Complex structure moduli space:

Consider a Ricci flat Kähler metric

$$g_{\alpha\bar{\beta}} dz^\alpha d\bar{z}^\beta \text{ on a CY manifold.}$$

Now regard this as a real manifold

Keep the metric fixed, and deform the

Complex structure:  ~~$g_{\alpha\bar{\beta}} dz^\alpha d\bar{z}^\beta$~~

To first order the deformed metric is:

$$\underbrace{(g_{\alpha\bar{\beta}} + \delta g_{\alpha\bar{\beta}})}_{\text{deformation of Kähler structure}} dz^\alpha d\bar{z}^\beta$$

$$+ \delta g_{\alpha\beta} dz^\alpha dz^\beta + \delta g_{\bar{\alpha}\bar{\beta}} d\bar{z}^\alpha d\bar{z}^\beta$$

represent deformation of  
Complex structure.

On a  $2n$  dimensional Calabi-Yau manifold  $h_{n,0} = 1$ ,  $h_{0,n} = 1$ .

$$\omega_{\alpha_1 \dots \alpha_n} dz^{\alpha_1} \wedge dz^{\alpha_2} \wedge \dots \wedge dz^{\alpha_n} \quad \text{with}$$

$$\bar{\partial} \omega = 0 \Rightarrow \bar{\partial}_{\beta} \omega_{\alpha_1 \dots \alpha_n} = 0 \quad \text{harmonic}$$

$\omega_{\alpha_1 \dots \alpha_n}$ : Holomorphic  $(n, 0)$  form.

Similarly  $\bar{\omega}_{\bar{\beta}_1 \dots \bar{\beta}_n}$ : Anti-holomorphic  $(0, n)$  harmonic form.

Consider

$$\omega_{\alpha_1 \dots \alpha_{n-1}} g^{\alpha \bar{\beta}} \delta g_{\bar{\beta} \bar{\gamma}}$$

$\Rightarrow$  an  $(n-1, 1)$  form.

$\Rightarrow$  Can be shown to be harmonic  
 $\in \mathcal{H}^{(n-1, 1)}$ .

Conversely, given an element of  $\mathcal{H}^{(n-1, 1)}$   
 we can construct  $\delta g_{\bar{\beta} \bar{\gamma}}$ .



⇒ Deformation of complex structure

↔ elements of  $H^{n-1,1} \in H^{1,n-1}$ .

⇒ Dimension of complex structure moduli space =  $h_{n-1,1}$  complex.

Cahli-Yau 3-fold: (6-dimensional)  
real.

$$h_{3,0} = h_{0,3} = 1$$

$$h_{1,2} = h_{2,1}, h_{1,1} = h_{2,2} \text{ not fixed.}$$

$$h_{1,0} = h_{2,0} = h_{0,1} = h_{0,2} = h_{2,3} = h_{3,2} \\ = h_{1,3} = h_{3,1} = 0$$

Cahli-Yau manifolds admit  
covariantly constant spinor.

e.g. take 6-d CY manifold.

$SO(6)$ : tangent space group.  $\equiv SU(4)$

Vectors: 6 of  $SO(6)$

Spinors: 4 of  $SO(6)$

Spinor field  $\xi$ :

$$(\mathbb{D}_m \xi)^t = \left( \delta_{st} \partial_m + \frac{i}{4} \omega_m^{ab} (\Sigma^{ab})_{st} \right) \xi^t$$

$\downarrow$   
So(6) generators in  
spinor representation.

Homomorphism group  $SU(3)$ .

Splits in first  $3 \times 3$  block if we  
regard  $SO(6)$  as  $SU(4)$ .

$$4 = 3 + 1$$

$\downarrow$   
Singlet.

$\Rightarrow$  One component does not transform  
under  $SU(3)$ .

$$\Rightarrow \omega_m^{ab} (\Sigma^{ab})_{st} \xi^t = 0 \text{ if}$$

$\xi$  is taken to be a singlet.

$$\Rightarrow \partial_m \xi^t = 0 \Rightarrow \xi^t = \text{constant}$$

Covariantly constant  $\llcorner$  spinor.



Explicit construction:

1. Pick a point  $P \in CY_3$ .
2. Pick a basis in spinor representation s.t. holonomy group at  $P$  is a:

$$\left( \begin{array}{cc|c} SU(3) & 0 & \\ \hline 0 & 1 & \end{array} \right)$$

3. Take  $\mathbb{Z}(P) = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$

& define  $\mathbb{Z}(Q)$  at other point  $Q$  by // transport along a curve joining  $P$  and  $Q$ .

The result is independent of the choice of path since holonomy group leaves  $\mathbb{Z}(P)$  invariant.

Consequence: Take a 10-d Majorana-Weyl spinor.

$$16 \rightarrow \left( \underset{\text{so}(6)}{4}, \text{left-handed Weyl}_{\text{so}(3,1)} \right) + \left( \overline{\underset{\text{so}(6)}{4}}, \text{right-handed Weyl}_{\text{so}(3,1)} \right).$$

Existence of a covariantly constant spinor on  $CY_3$ .

$\Rightarrow$  Existence of a 16-d spinor satisfying:

$$D_m \psi = 0 \quad D_\mu \psi = 0$$

$\Downarrow$   
 $D_m$  if we

$\Downarrow$   
 $D_\mu$  if  $(3+1)$ -d spacetime is flat.

Choose the  $SU(3)$

singlet component of  $\psi$ .

$\Rightarrow$  4 ~~complex~~ fermion solns.

(left-handed Weyl + right-handed Weyl)

$\Downarrow$   
 2 complex = 4 real.



Heterotic or type I on  $CY_3$

$\Rightarrow N=1$  SUSY in  $D=4$ .

Type IIA / IIB on  $CY_3$

$\Rightarrow N=2$  SUSY in  $D=4$ .