

Neutrinos decouple at $T \sim \text{MeV}$.

After this $T_\nu = \frac{T_F \lambda_F}{\lambda}$

T_F : Freezeout temp.

λ_F : λ at freezeout.

e^+, e^-, γ remain in thermal equilibrium.

T : temperature of this system.

How does T evolve?

$$s \lambda^3 = \text{constant} \Rightarrow s_F \lambda_F^3 = s \lambda^3$$

F.M. interaction: $\frac{1}{m_W^2} \rightarrow \frac{1}{E^2} \sim \frac{1}{T^2}$

Apply this for $T \ll .5 \text{ MeV} \sim m_e \Rightarrow e^+ e^-$ non-relativistic

$$g_{\text{eff}}^{(1)} T_F^3 \lambda_F^3 = g_{\text{eff}}^{(2)} T^3 \lambda^3 = 2 \quad g_{\text{eff}}^{(1)} = 2 + \frac{7}{8} (2+2) = \frac{11}{2}$$

$$\frac{T_F \lambda_F}{\lambda} = T \lambda \left(\frac{4}{11} \right)^{1/3}$$

$$T_\nu = T \left(\frac{4}{11} \right)^{1/3}$$

$$\text{today } T_\nu = 2.73 \times \left(\frac{4}{11} \right)^{1/3}$$

For massless neutrinos we have black body spectrum at temperature T_ν .

For massive neutrinos the distribution is more complicated but follows what we derived earlier.

The range $\cdot 1 \text{ MeV} \lesssim T \lesssim 1 \text{ MeV}$ is important for another reason.

We have excess baryons $\rightarrow$ in thermal equilibrium with the bath at temperature T .

$\nearrow$ neutrino freezeout
Below T_F , π is the bath that the baryons feel.

At $T \gtrsim 1 \text{ MeV}$, $(n_B \lambda^3)$ is constant but $n \leftrightarrow p$ transition is possible via weak $n \rightarrow p + e^- + \bar{\nu}_e$, $n + \nu_e \leftrightarrow p + e^-$, ... | interaction.

Nuclear binding energy $\sim \text{MeV/nucleon}$

$T < \text{MeV}$, p, n will bind
into Nuclei $p, n \downarrow$

$\rightarrow$ By studying this process carefully
we'll be able to predict the relative
abundance of light elements ($d, \text{He}^3, \text{He}^4,$
 $\text{Li}^3, \dots$)

Heavy elements are formed in
Stars.

Broad overview of nucleosynthesis

① Above $T > \text{MeV}$, when $n + t \leftrightarrow p + t$ is in equilibrium, we expect

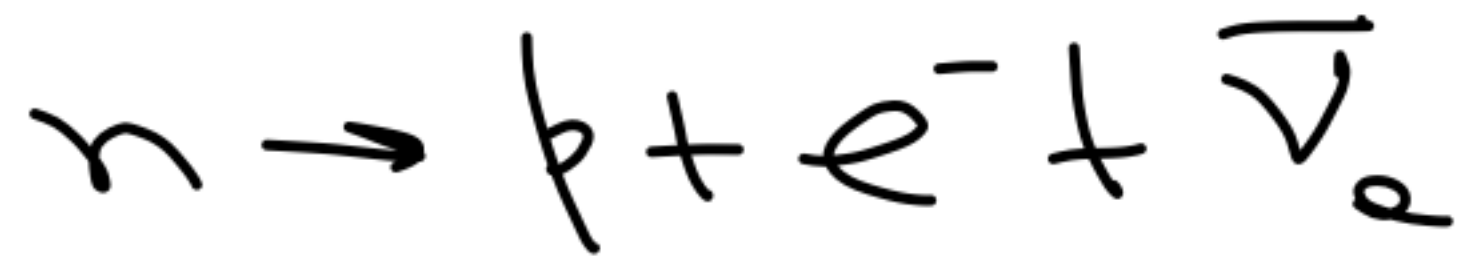
$$\frac{n_n}{n_p} \sim e^{-(m_n - m_p)/T} \quad m_n - m_p \sim \text{MeV}.$$

For $T \gg \text{MeV}$, $n_n \sim n_p$

For $T < \text{MeV}$, $\frac{n_n}{n_p}$ starts reducing.

At the same time, weak interaction falls out of equilibrium.

The reason that this number is not frozen is because n is not stable.



Half life of a free neutron ~ 13 minutes
 ~ 800 secs.

We need to compare this with the time scale that is involved as T varies between MeV and ~ 1 MeV.

Rough estimate: $\left(\frac{\dot{\lambda}}{\lambda}\right)^2 = \frac{8\pi G}{3} \rho \approx \frac{8\pi G}{3} \times T^4 \times \text{const.}$

$$\lambda \sim \frac{1}{T} \quad \left(\frac{\dot{\lambda}}{\lambda}\right)^2 \sim G T^4 \quad (G=1)$$

$$\sim T^4$$

$$\dot{T} = T^3$$

$$\frac{1}{T^2} = t$$

$$T \approx t^{-1/2}$$

$$T = 1 \text{ MeV} = \frac{1 \text{ MeV}}{10^{19} \text{ GeV}} \sim 10^{-22} \text{ in natural units.}$$

$$t \sim 10^{-44} t_{pl} \sim 1 \text{ sec.}$$

$$5 \times 10^{44} \text{ sec.}$$

$t \sim 100 \text{ sec}$ for
n. half life $T \sim 1 \text{ MeV}$

② Formation of bound state.

The beginning of the process requires some non-equilibrium process.

But once the formation starts, almost all n 's get bound inside He^4 (very few remain in other light nuclei like

$$d, \text{He}^3, \dots) \Rightarrow \frac{n_n}{n_p} = \frac{2n_{\text{He}^4} + \dots}{2n_{\text{He}^4} + n_H + \dots}$$

$$\Rightarrow \frac{n_{\text{He}^4}}{n_H} = \frac{n_n}{2(n_p - n_n)} \Rightarrow \text{If we know } n_n/n_p \text{ at the time of bound state formation, we know } n_{\text{He}^4}/n_H.$$

Details.

Step I. We need to find time dependence
of T and T_v
↓
Temperature of the bath
→ determines the number
of ambient $v, \bar{v}$
affects $n \leftrightarrow \frac{1}{2}$ transition.

$$\begin{aligned}
 \textcircled{i} \quad \mathcal{H}_{e^+e^- \gamma} &= \frac{2}{3} a_B T^3 \times 2 + \frac{p_{e^+e^-} + p_{e^-e^+}}{T} \\
 &= \frac{4}{3} a_B T^3 + \frac{1}{2\pi^2} * (2 + 2) \int_0^\infty k^2 dk \\
 &\quad \left[\ln(1 + e^{-\sqrt{k^2 + m^2}/T}) + \frac{1}{T} \sqrt{k^2 + m^2} \frac{e^{-\sqrt{k^2 + m^2}/T}}{1 - e^{-\sqrt{k^2 + m^2}/T}} \right]
 \end{aligned}$$

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