

Review

n_e : no. density of free electrons

n_{1s} : no density of 1s atoms.

$n_{n,l}$: no. " " " n, l "

$$n_{n,l} = n_{2s} (2l+1) \exp(-(B_2 - B_n)/T)$$

$$\Rightarrow \sum_{\substack{n,l \\ n \geq 2, l}} n_{n,l} = n_{2s} f(T) \rightarrow \sum_{\substack{n,l \\ n \geq 2}} (2l+1) \exp(-(B_2 - B_n)/T)$$

$$n_e + n_{1s} + n_{2s} f(T) = 0.73 \times n_B$$

$$\text{For } T \ll (B_2 - B_3), \quad f(T) \approx 4$$

$$x = \frac{n_e}{0.73 n_B}, \quad y = \frac{n_{1s}}{0.73 n_B}, \quad n_{2s} = \frac{0.73 n_B (1 - x - y)}{f(T)}.$$

$$\frac{dx}{dt} = -\alpha x^2 (0.73 n_B) + \frac{\beta}{f(\tau)} (1-x-y)$$

$$\frac{dy}{dt} = (\Gamma_{2s} + 3\Gamma_{2p} P) \frac{1}{f(\tau)} (1-x-y) - \varepsilon y$$

$\alpha = \sigma v$
 $\hookrightarrow e^+p$ capture x-section into nl states for $n \geq 2$.

$$\beta = \alpha \left(\frac{m_e T}{2\pi} \right)^{3/2} e^{-B_2/T}, \quad \varepsilon = (\Gamma_{2s} + 3\Gamma_{2p} \bar{P}) e^{-\frac{B_1 - B_2}{T}}$$

Γ_{2s} : decay rate of $2s \rightarrow 1s + \gamma + \gamma$

Γ_{2p} : " " " $2p \rightarrow 1s + \gamma$

$\alpha, \Gamma_{2s}, \Gamma_{2p}$ can be calculated using microscopic theory or measured.

$P(t)$: The probability that a Lyman α photon produced at time t is not captured by a $1s \rightarrow 2p$ transition.

$\bar{P}(t)$: Fraction of thermal photons (near Lyman α) that survive capture up to time t .

Preliminaries: $p(\omega) d\omega$ will denote the probability that a photon emitted during $2p \rightarrow 1s$ transition has energy in the range $(\omega, \omega + d\omega)$.

$$p(\omega) = \frac{\Gamma_{2p}}{2\pi} \frac{1}{(\omega - \omega_\alpha)^2 + (\Gamma_{2p}/2)^2} \quad \text{Ex. } \int_{-\infty}^{\infty} p(\omega) d\omega = 1$$

$\Gamma_p \ll \omega_\alpha$ | ω_α is the dist. of Γ_{2p} .

$\sigma(\omega)$: defined to be the absorption x-section of a photon of energy ω by 1s state to make transition to 2p state.

$$\sigma(\omega) = \frac{3\pi^2}{\omega_a^2} \Gamma_{2p} \cdot p(\omega).$$

Suppose a photon of energy ω is produced at time $\underline{t}$.

Total probability that it is absorbed between t' and $t' + dt'$ assuming that it survives till t' :

$$\sigma\left(\omega \frac{\lambda(t)}{\lambda(t')}\right) n_{1s}(t') \underset{=1}{\subset} dt'$$

Define $Q_\omega(t')$ to be the probability that a photon of energy ω emitted at time t survives till t' .

$$\frac{dQ_\omega(t')}{dt'} = -Q_\omega(t') \times \sigma\left(\omega \frac{\lambda(t)}{\lambda(t')}\right) n_{is}(t')$$

B. c. $Q_\omega(t') = 1$ at $t' = t$.

$$\Rightarrow Q_\omega(\infty) = \exp\left[-\int_t^\infty dt'\right]$$

$$\sigma\left(\omega \frac{\lambda(t)}{\lambda(t')}\right) n_{is}(t')$$

change variables: From t' to $\omega' = \omega \frac{\lambda(t)}{\lambda(t')}$
 $d\omega' = -\omega \frac{\lambda(t)}{\lambda(t')^2} \dot{\lambda}(t') = -\omega' \lambda(t')$

$$Q_{\omega}(\infty) = \exp \left[- \int_0^{\omega} d\omega' n_{1s}(t') \sigma(\omega') \frac{1}{\omega' \hbar(t')} \right]$$

$$\approx \exp \left[- n_{1s}(t) \frac{1}{\omega_{\alpha} \hbar(t)} x(t) \right]$$

$$\times \left[\frac{3\pi^2 \Gamma_{2p}}{\omega_{\alpha}^2} \int_0^{\omega} d\omega' p(\omega') \right]$$

$\frac{3\pi^2 \Gamma_{2p}}{\omega_{\alpha}^2} p(\omega')$

→ assuming that $p(\omega')$ is sharply peaked around ω_{α} .

$P(t)$: Probability that a photon emitted at t by $2p \rightarrow 1s$ transition is not absorbed by $1s \rightarrow 2p$ transition

$$= \int_{-\infty}^{\infty} p(\omega) d\omega Q_{\omega}(\infty)$$

$$P = \int_{-\infty}^{\infty} d\omega \, p(\omega) \exp \left[-\pi \int_0^{\omega} p(\omega') d\omega' \right]$$

$$p(\omega) = \frac{\Gamma_{2p}}{2\pi} \frac{1}{(\omega - \omega_{\alpha})^2 + (\Gamma_{2p}/2)^2}$$

Change variables to

$$u = \frac{2}{\Gamma_{2p}} (\omega - \omega_{\alpha}),$$

$$\vartheta = \frac{2}{\Gamma_{2p}} (\omega' - \omega_{\alpha})$$

$$P = \frac{1}{\pi} \int_{-\infty}^{\infty} du \, \frac{1}{u^2 + 1} \exp \left[-\frac{\pi}{\pi} \int_{-\infty}^u \frac{v}{v^2 + 1} d\vartheta \right]$$

Change variable
 $s = \tan^{-1} u$

$$\tan^{-1} u \neq \frac{\pi}{2} \equiv$$

$$\omega = \frac{2\omega_{\alpha}}{\Gamma_{2p}} \text{ large.}$$

Ex. $P = F(x(t))$, $F(x) = \frac{1}{x} (1 - e^{-x})$

$$x = \frac{3\pi^2 n_{15}(t)}{\omega_x^3 H(t)} \quad \Gamma_{2p}$$

$$\Downarrow F(x) \xrightarrow{x \rightarrow 0} \uparrow$$

① If $H(t) = 0 \Rightarrow x = \infty$, $F(\infty) = 0$

② x increases with n_{15}

$P = F(x)$ decreases with n_{15} .

$2p \rightarrow 1s + \gamma \Rightarrow$ dominates initially.

$2s \rightarrow 1s + \gamma + \gamma \Rightarrow$ dominates at late time.

Calculation of $\bar{P}$

$f(\omega) d\omega = \#$ density of thermal photons
in the range $(\omega, \omega + d\omega)$ assuming no
absorption by 1s atom.

Probability / time that a 1s atom makes
transition to 2p by absorbing a thermal photon
(assuming no prior absorption)

$$= \int \sigma(\omega) f(\omega) d\omega = \pi \pi_0.$$

Define π as the probability/unit time
for a thermal photon to cause $1s \rightarrow 2p$
transition, taking into account prior
absorption effects.

$$\bar{p} = \pi / \pi_0$$

$$\pi_0 = \int_{-\infty}^{\infty} \sigma(\omega) f(\omega) d\omega = \frac{3\pi^2 \Gamma_{2p}}{\omega_a^2} \int_{-\infty}^{\infty} p(\omega) f(\omega) d\omega$$
$$= \frac{3\pi^2 \Gamma_{2p}}{\omega_a^2} f(\omega_a) \int_{-\infty}^{\infty} p(\omega) d\omega = 1$$

Define: $q_\omega(t)$: The probability that a thermal photon of energy ω at time t has survived capture from $t' = -\infty$ to t .

$$q_\omega(t) = \exp \left[- \int_{-\infty}^t dt' n_{is}(t') \sigma \left(\omega \frac{\lambda(t)}{\lambda(t')} \right) \right]$$

Change variable $\omega' = \omega \frac{\lambda(t)}{\lambda(t')}$

$$q_\omega(t) = \exp \left[- \frac{3\pi^2 \Gamma_{2f}}{\omega_\alpha^2 \omega_\alpha H(t)} \int_{\omega}^{\infty} d\omega' f(\omega') \right]$$

x

π : Probability that a thermal photon
 (keeping track of absorption effects) cause
 transition from a $1S \rightarrow 2P$.

$$= \int_{-\infty}^{\infty} d\omega f(\omega) g_{\omega}(t) \sigma(\omega)$$

$$= \frac{3\pi^2 \Gamma_{2P}}{\omega_x^2} f(\omega_x) \int_{-\infty}^{\infty} p(\omega) d\omega \exp\left[-x \int_{\omega}^{\infty} p(\omega') d\omega'\right]$$

$$\pi_x = \pi_0 F(x(t)) \Rightarrow \bar{p} = \frac{\pi}{\pi_0} = F(x(t))$$