

# Introduction to Strings: Part II

*Largely an open-ended story...*

Debashis Ghoshal

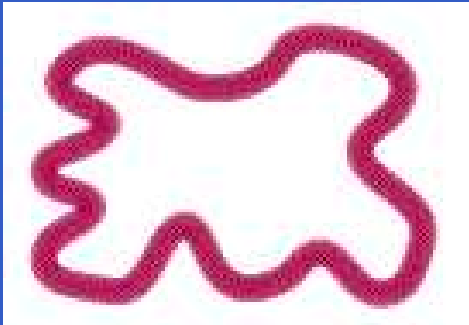
Harish-Chandra Research Institute

Allahabad

# Overview

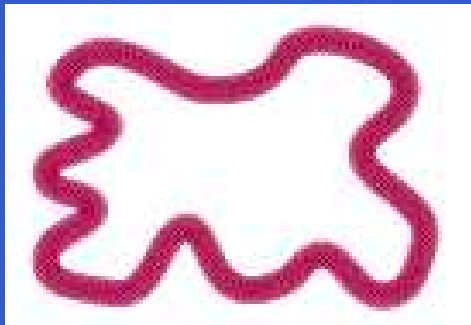
- Recap

[animated] Closed string



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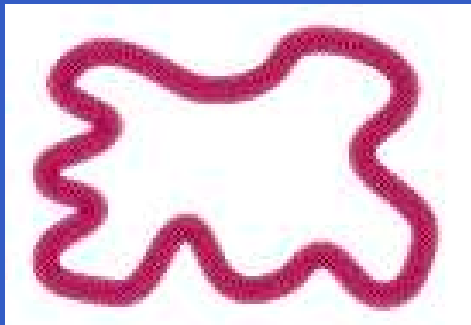


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- Open strings ... need boundary conditions
  - ◇ Neumann condition
  - ◇ Dirichlet condition

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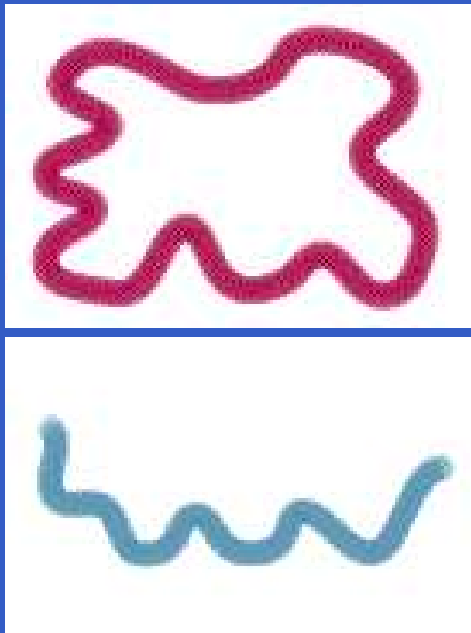


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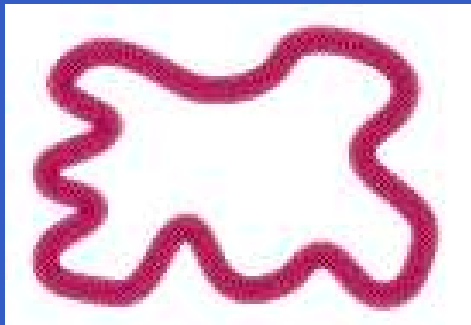


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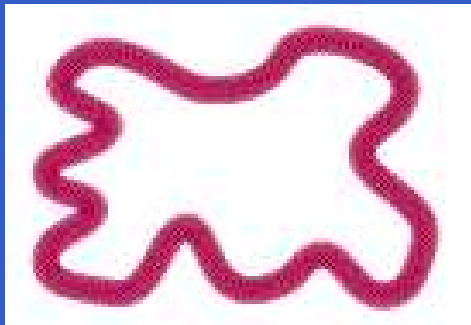


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- D-branes
- Open string interactions

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[animated] Open string

- Recap
- Open strings ... need boundary conditions
  - ◇ Neumann condition
  - ◇ Dirichlet condition
- Spectrum of open strings
- D-branes
- Open string interactions
- A larger picture?

# Recap

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strings of size  $\sim 10^{-33}$  cm.



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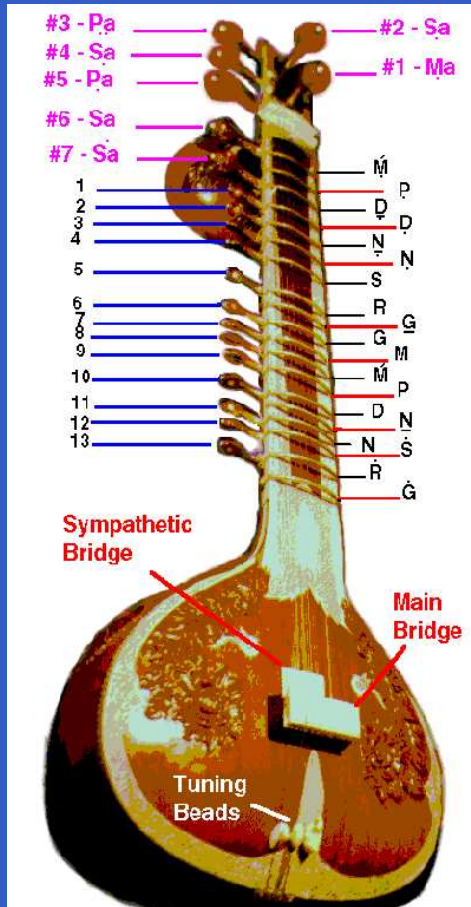
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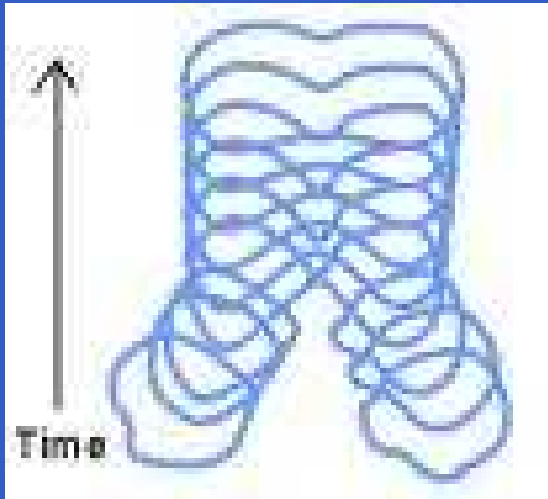


Fundamental constituent:  
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**Normal modes** of vibration of  
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gravitons, electrons, ...  
**The music of the strings!**

## Recap... (ii)

One string can **split** into two  
or two strings can **join** into  
one:

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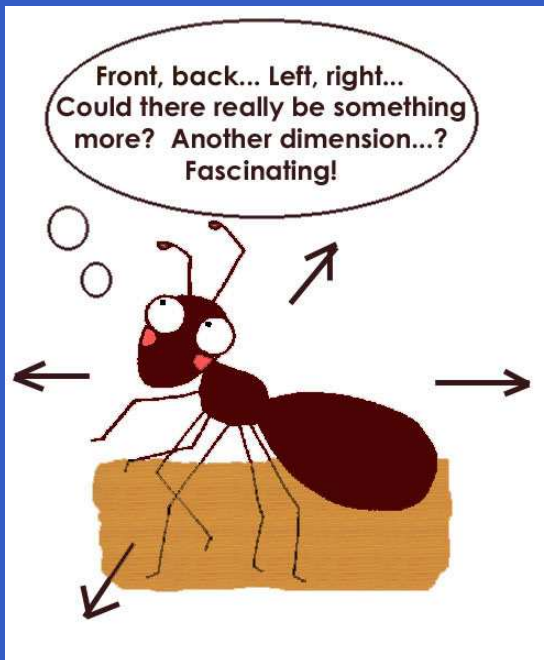


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**Interaction**



# Recap... (ii)



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String theory is consistent only  
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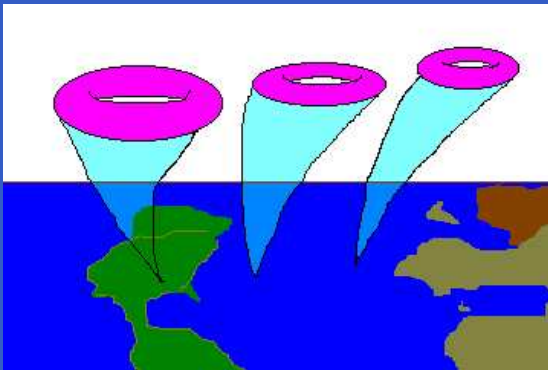
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**Compactification:** Six extra di-  
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# Everyday strings

The strings that we see around us are macroscopic nonrelativistic strings.

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These strings need boundary conditions.

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Fundamental open strings also need boundary conditions.

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Boundary conditions are independent at each end point.

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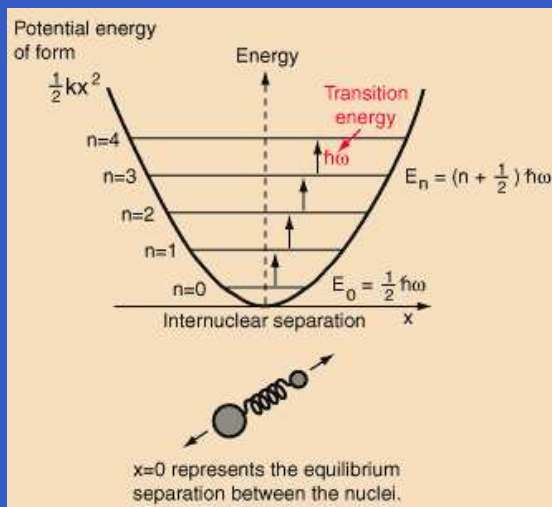
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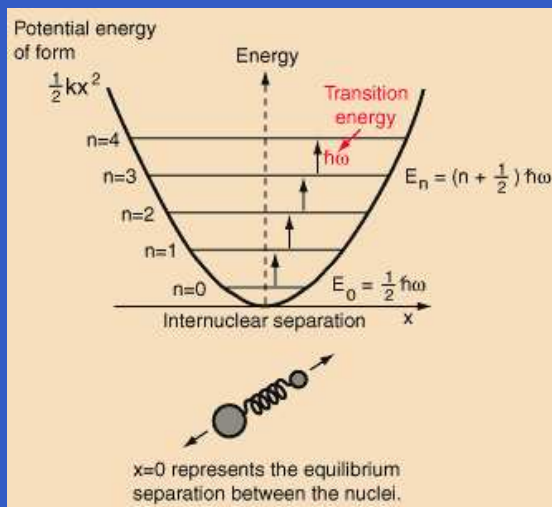
For Open strings: reflection at the end-points relate these modes. These are no longer independent:

$$x_n^i = \pm \tilde{x}_n^i$$

Each coordinate  $x^i$  is a collection of infinite number of simple harmonic oscillators

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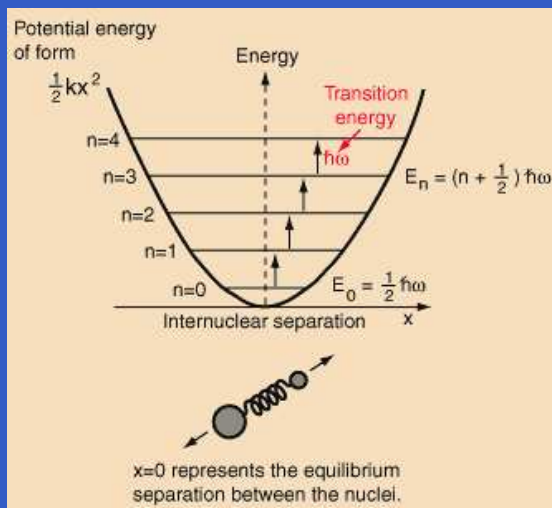
vacuum state  $|0\rangle$

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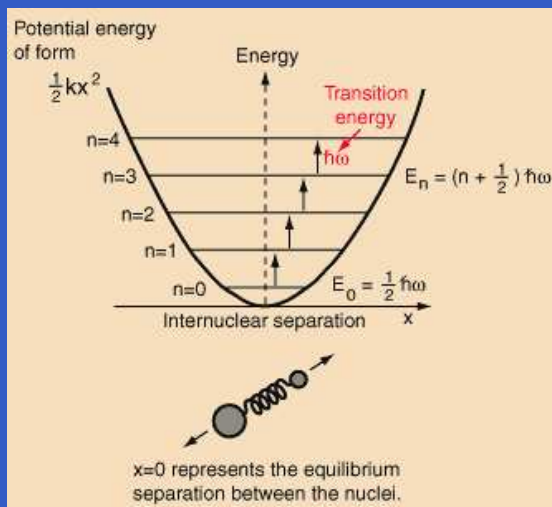
Each coordinate  $x^i$  is a collection of infinite number of simple harmonic oscillators

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vacuum state carrying momenta

$$|k\rangle = \prod (a_0^i)^\dagger |0\rangle$$

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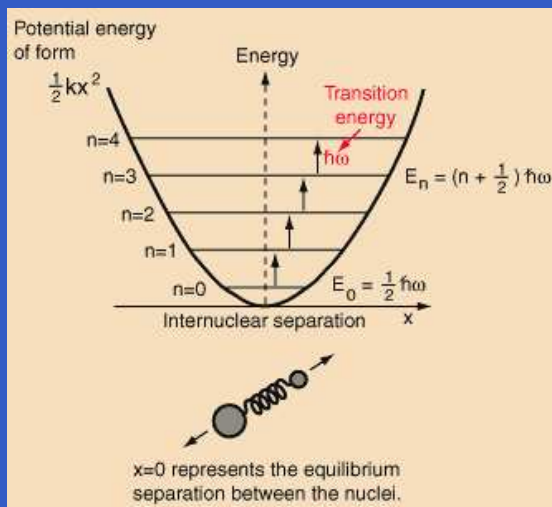
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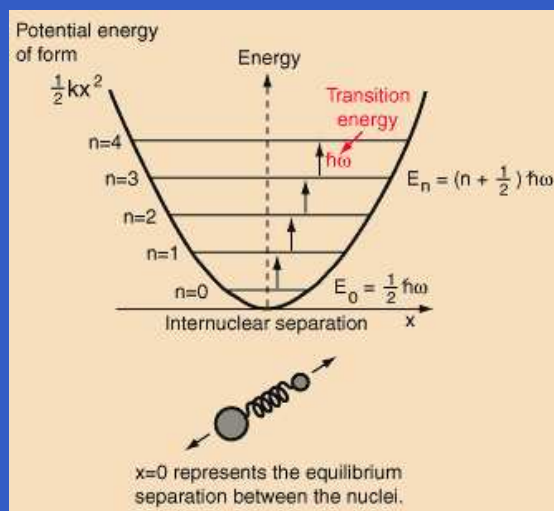
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$$(a_2^i)^\dagger |k\rangle, \quad (a_1^i)^\dagger (a_1^j)^\dagger |k\rangle$$

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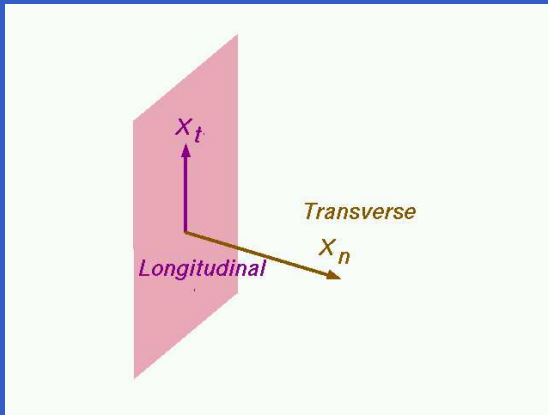
For Open strings: reflection at the end-points relate these modes. These are no longer independent:

$$x_n^i = \pm \tilde{x}_n^i$$

$$(a_1^i)_{IJ}^\dagger |k\rangle$$

The extra indices  $I, J$  refer to the end-points.

# Spectrum of open strings... (ii)



Open string can carry momentum only along the (hyper-)plane:

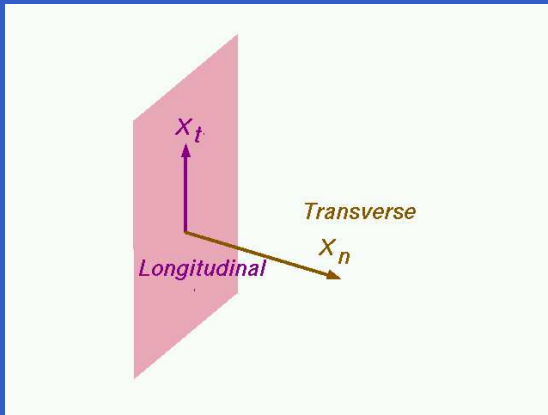
Neumann b.c. on  $X_{||}^i$

Dirichlet b.c. on  $X_{\perp}^i$

String end-points are on the hyperplane

$X_{\perp}^i = \text{constant}$

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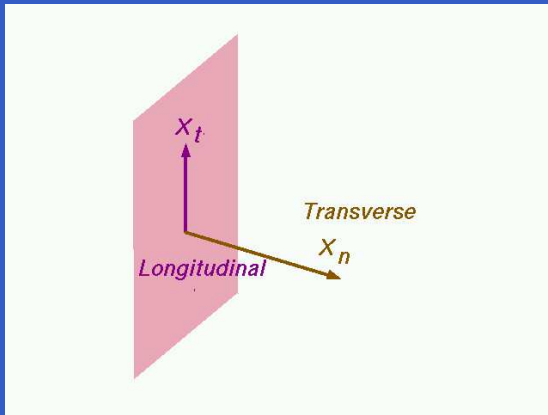
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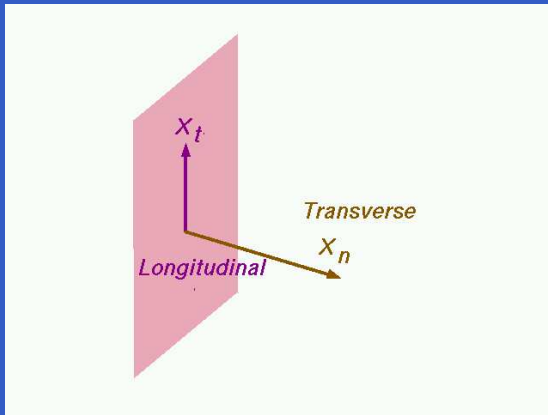
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$$\left( a_1^{\parallel i} \right)_{IJ}^{\dagger} |k_{\parallel}\rangle \quad \text{Gauge bosons.}$$

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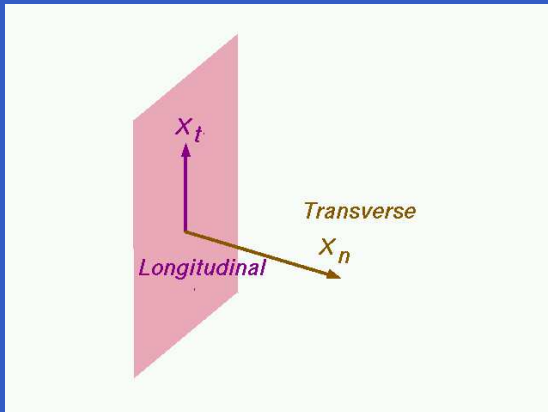
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Both live on the hyper-plane.



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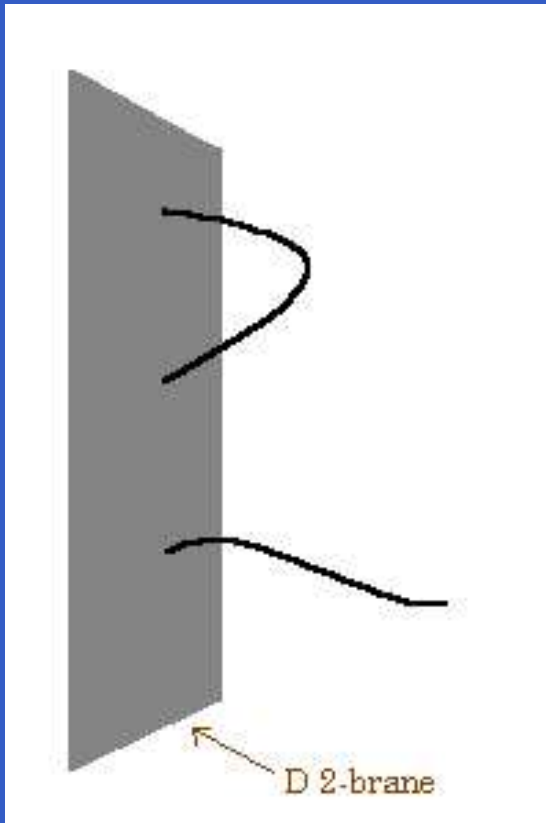
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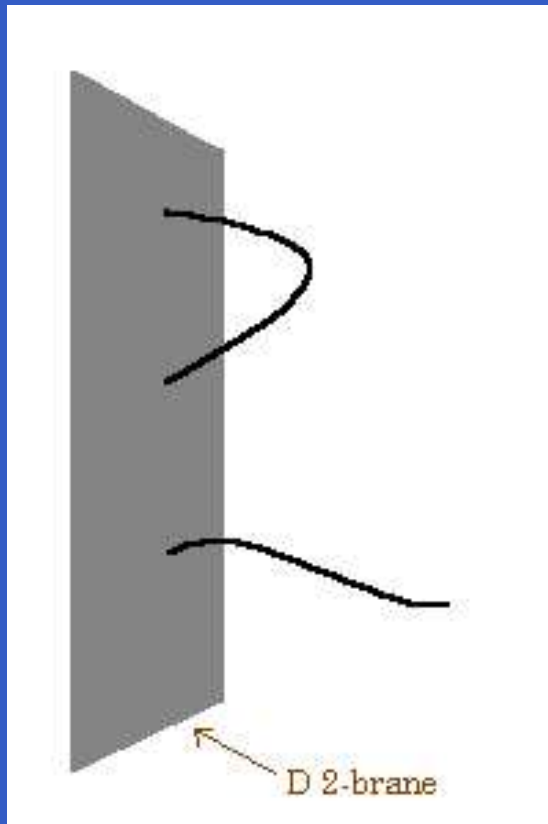
Both live on the hyper-plane. Position of the plane = Expectation value of these scalars.

# More on boundary conditions



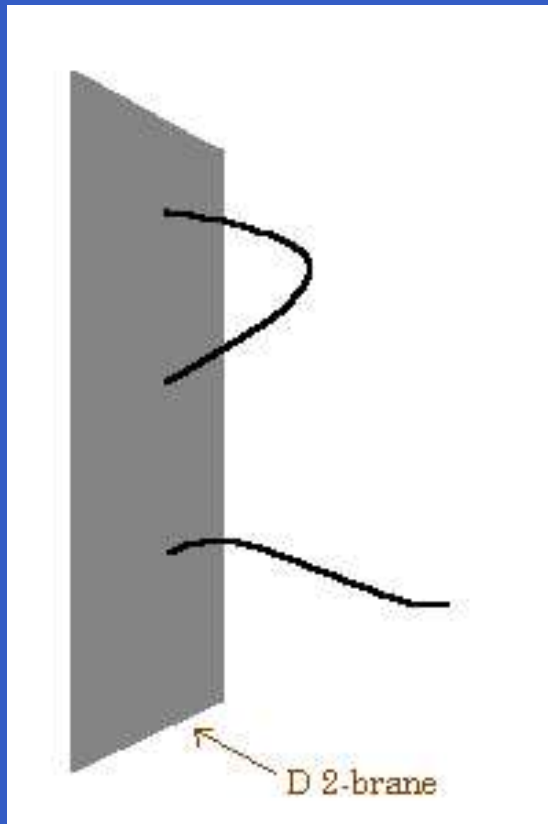
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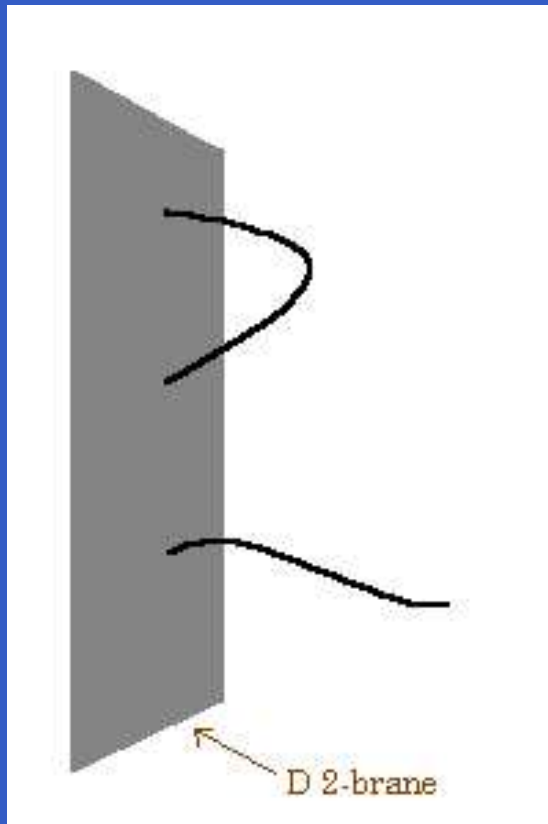
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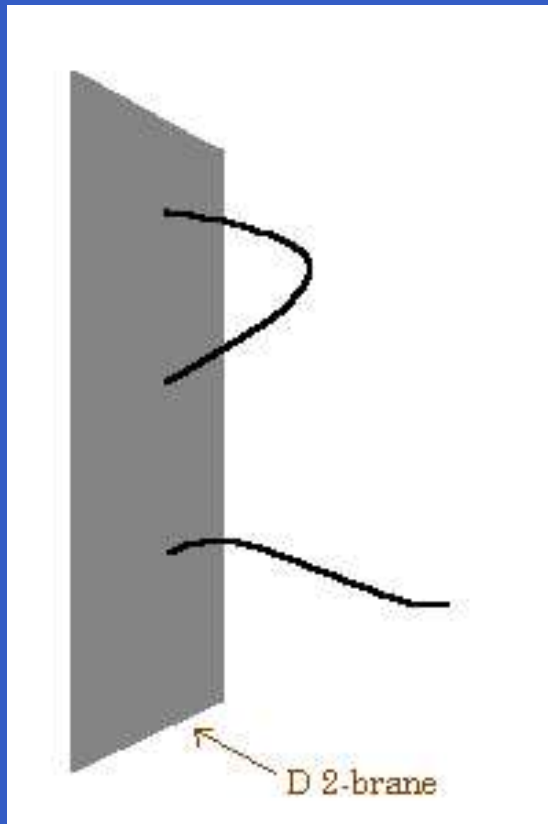


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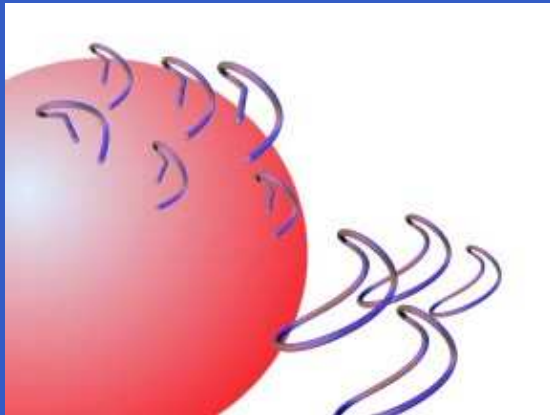
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Dirichlet 2-brane, or,

**D-2-brane**

# More on boundary conditions



Can be any  
(hyper-)surface.

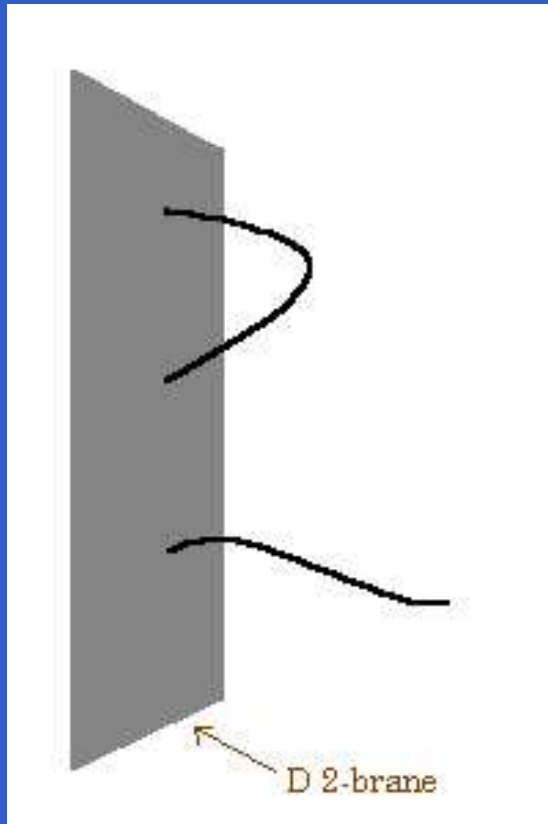
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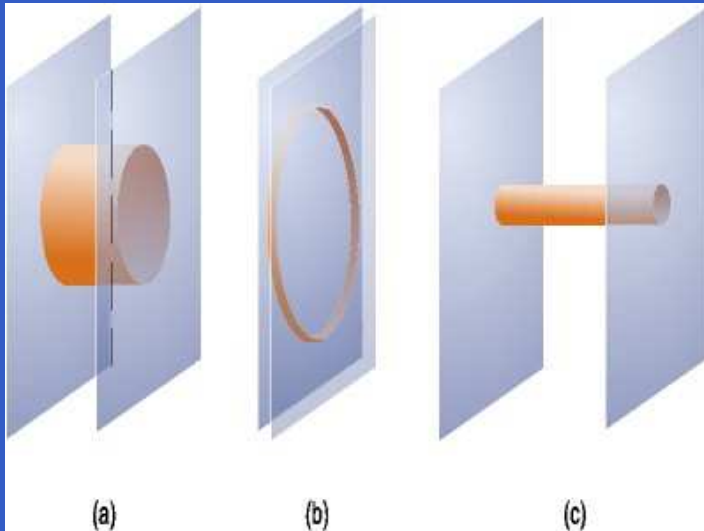
# More on boundary conditions ... (ii)



D-branes look like *defects/impurities* in space.

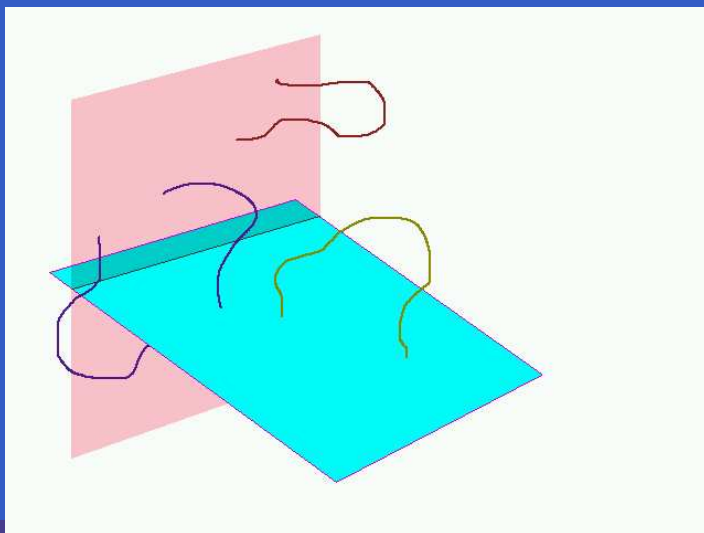
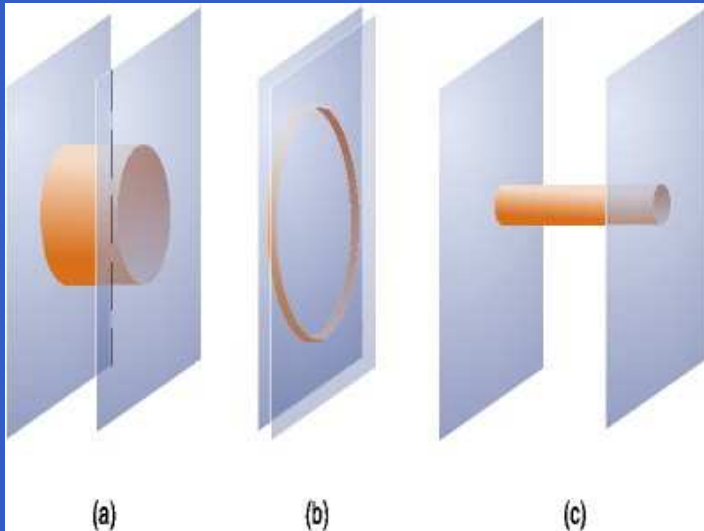


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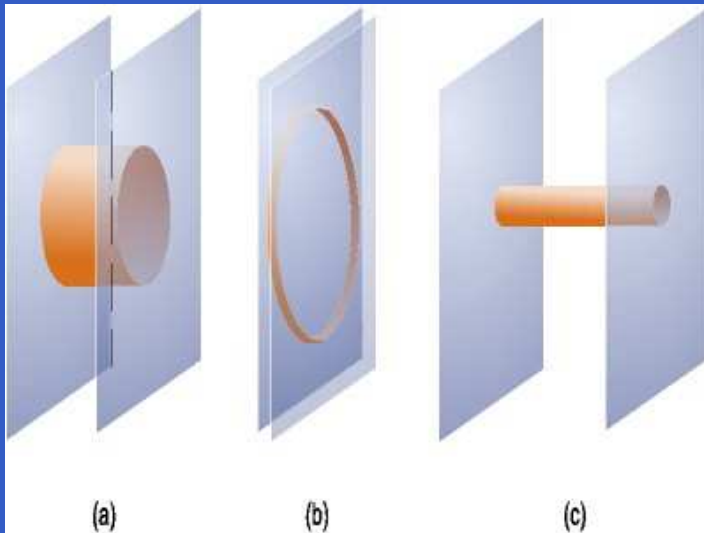
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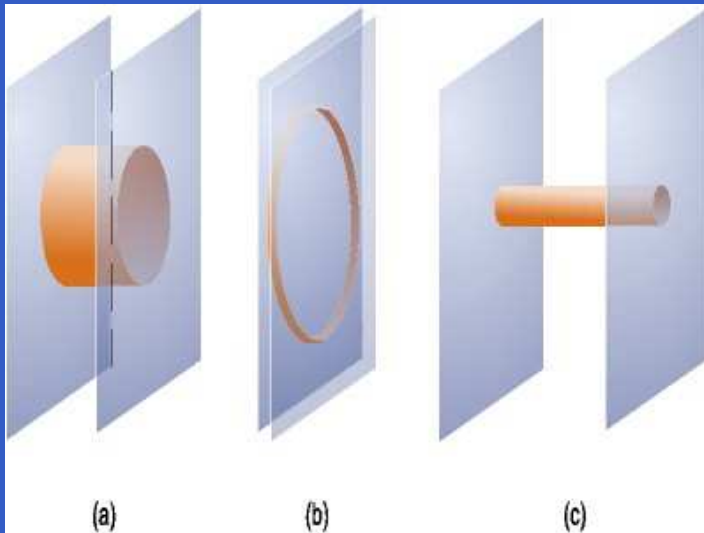


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Energy of the strings between branes depend on the stretching,

# More on boundary conditions ... (ii)



$$A_{IJ}^i(\mathbf{k}) = (a^i)_{IJ}^\dagger |\mathbf{k}\rangle$$

can become massless

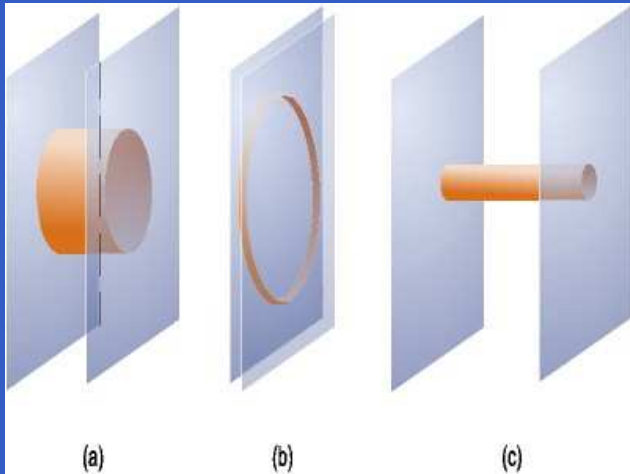
$$E^2 = |\mathbf{k}|^2$$

D-branes look like *defects/impurities* in space.

There may be more than one D-brane, aligned arbitrarily — can intersect each other.

Energy of the strings between branes depend on the stretching, *i.e.* on the **distance** between the branes.

# More on boundary conditions ... (ii)

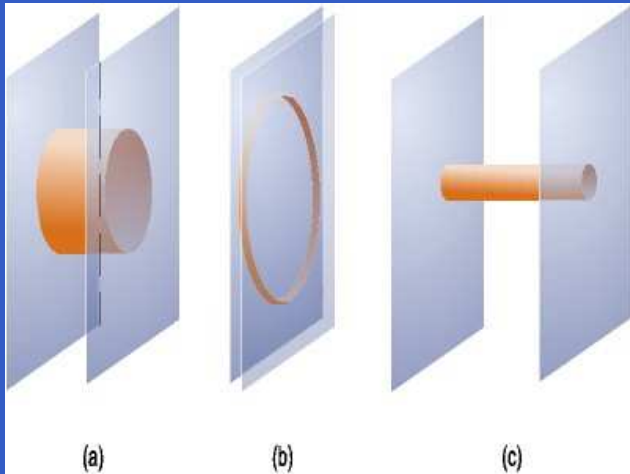


$A_{II}^i(\mathbf{k}) = (a^i)_{II}^\dagger |\mathbf{k}\rangle$  come from excitations of a string starting and ending on the *same* D-brane.

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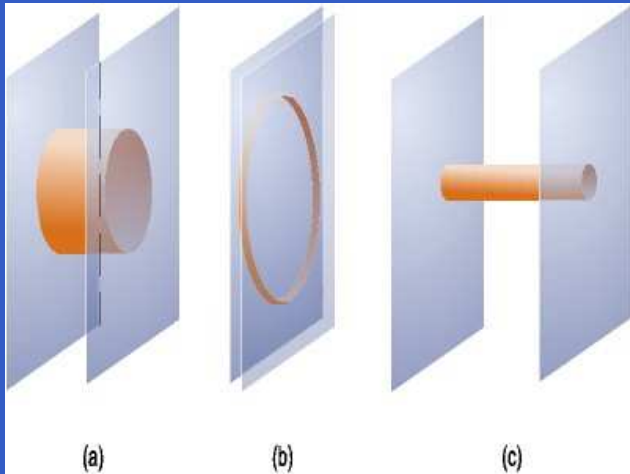
$A_{II}^i(\mathbf{k}) = (a^i)_{II}^\dagger |\mathbf{k}\rangle$  come from excitations of a string starting and ending on the *same* D-brane.

It is always massless.

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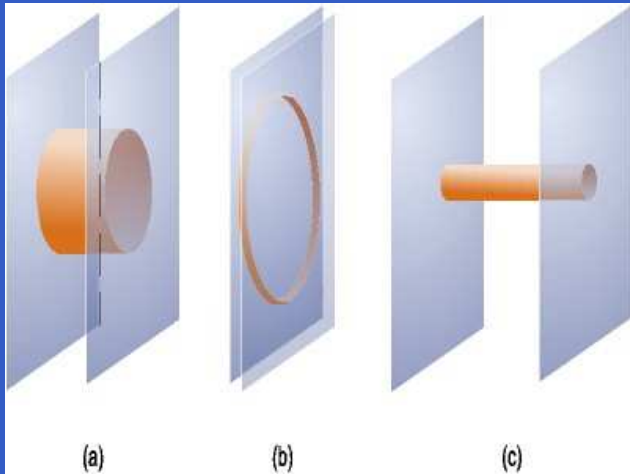
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$$E^2 = |\mathbf{k}|^2$$

These are like photons — satisfy **Maxwell eqns** plus corrections.

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If there are two D-branes, there are two types of photons  $A_{11}$  and  $A_{22}$ .

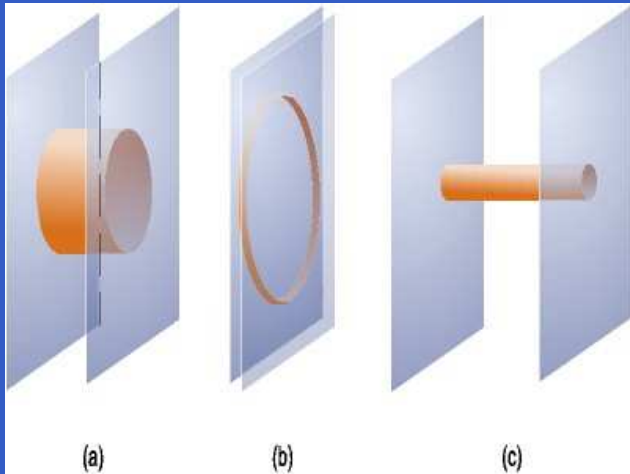
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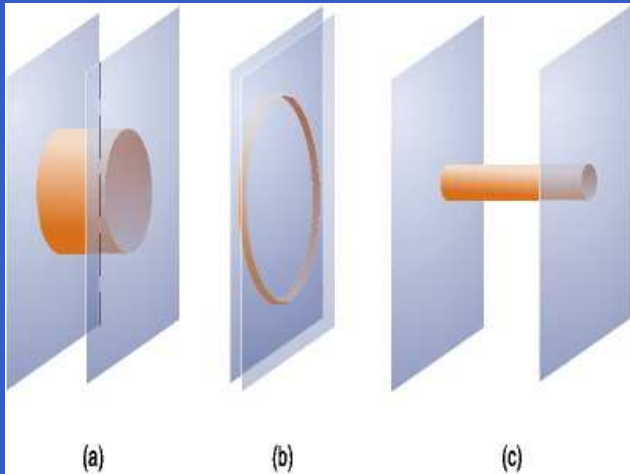
If there are two D-branes, there are two types of photons  $A_{11}$  and  $A_{22}$ .

If they coincide,  $A_{12}$  and  $A_{21}$  also become massless.

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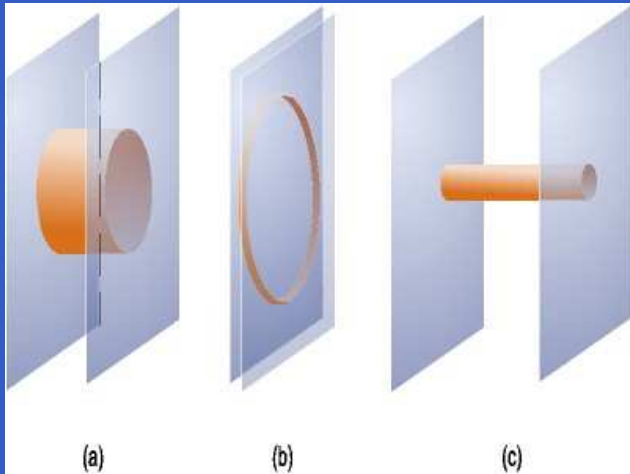
Now all *four* fields  $A_{11}$ ,  $A_{12}$ ,  $A_{21}$  and  $A_{22}$  are massless.

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Now all *four* fields  $A_{11}$ ,  $A_{12}$ ,  $A_{21}$  and  $A_{22}$  are massless.

Also massless are *transverse scalar* fields  $X_{IJ}^{\perp}$ .

$$A_{IJ}^i(\mathbf{k}) = (a^i)_{IJ}^{\dagger} |\mathbf{k}\rangle$$

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For superstrings one has

supersymmetric non-abelian  $U(N)$  gauge theory .

# The world of the end

The *dynamics* of open strings come from the end-points.

# The world of the end



The *dynamics* of open strings come from the end-points.



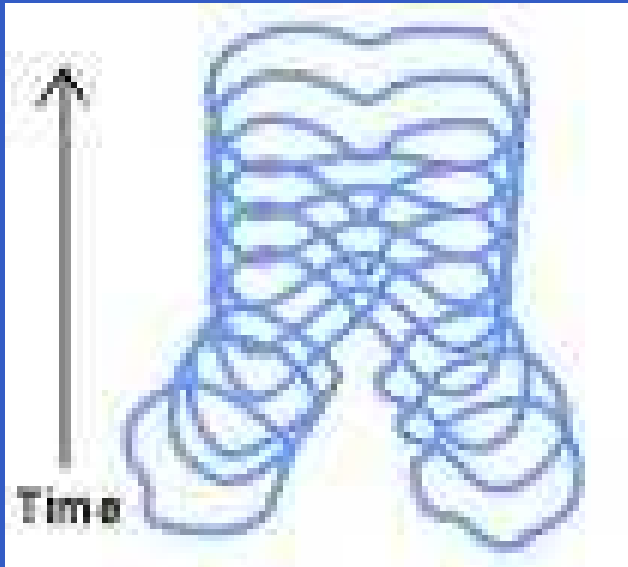
# The world of the end



The *dynamics* of open strings come from the endpoints.

This may be how we look to the big *green* giants living in *higher dimensions*!

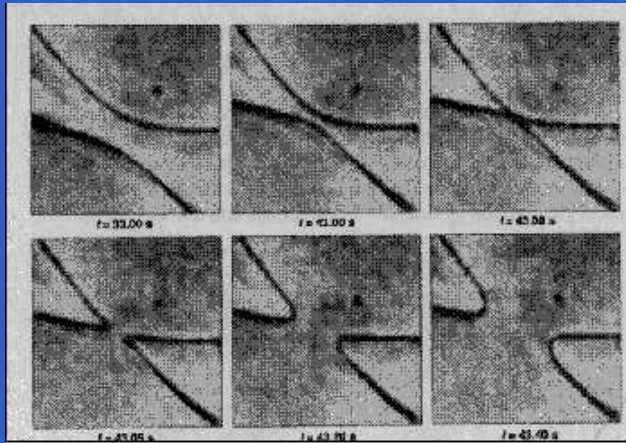
# String interactions



Closed strings interact by splitting and joining.

String scattering

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Interaction of strings in nematic liquid crystal

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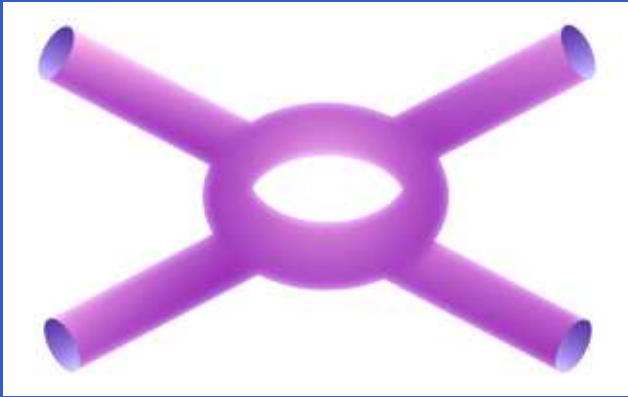
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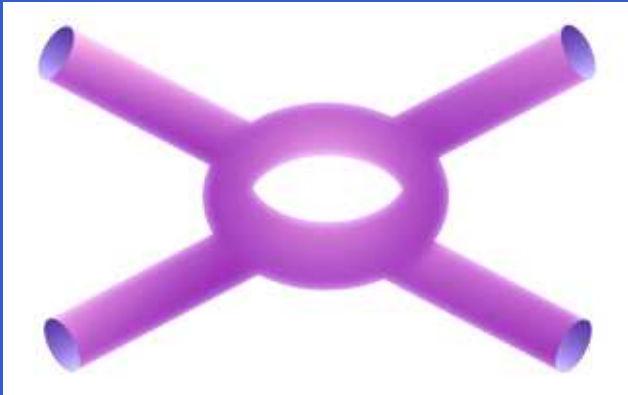


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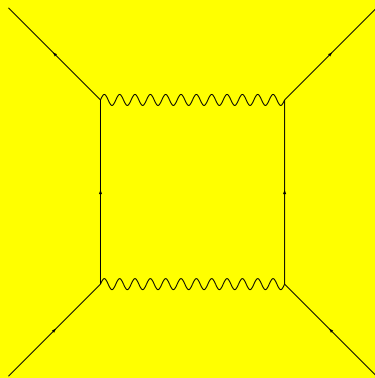


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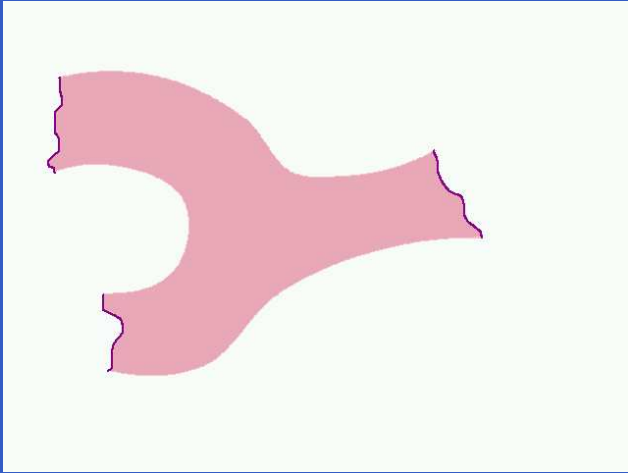
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Small string tension ( $\alpha' \rightarrow 0$ ) limit.



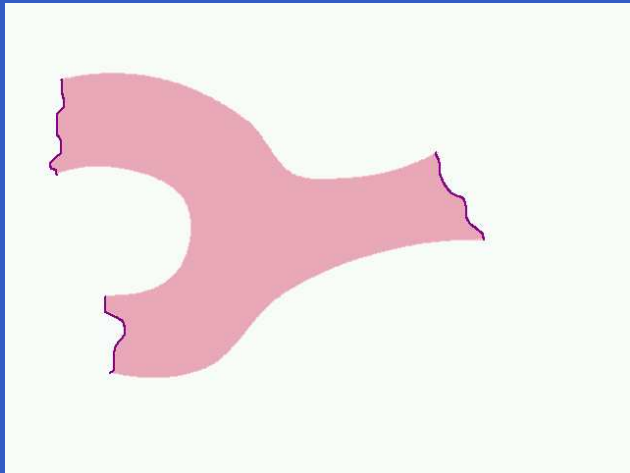


# Open string interactions



Open strings interact by splitting and joining **at the end points**.

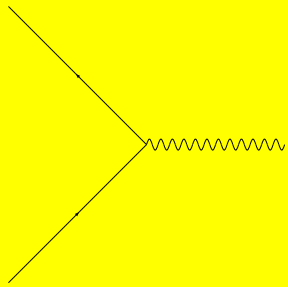
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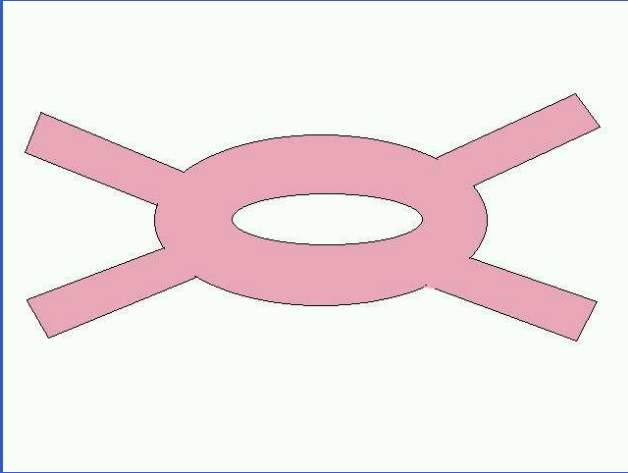
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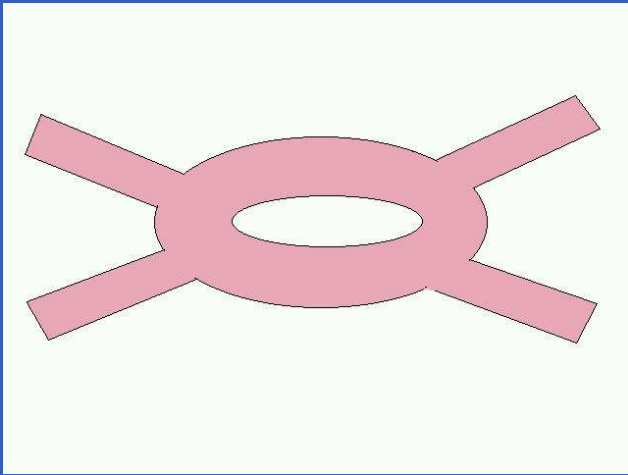
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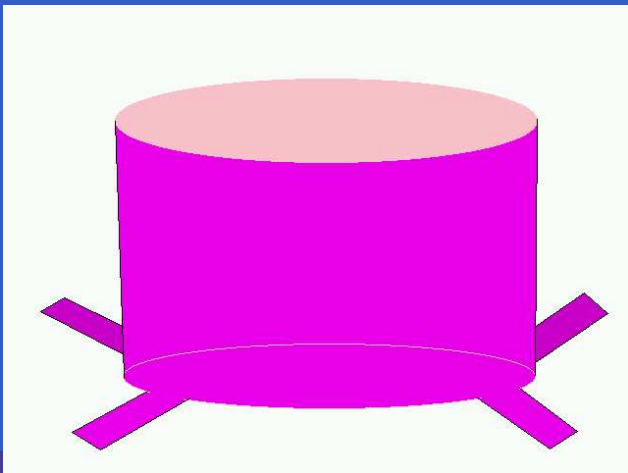
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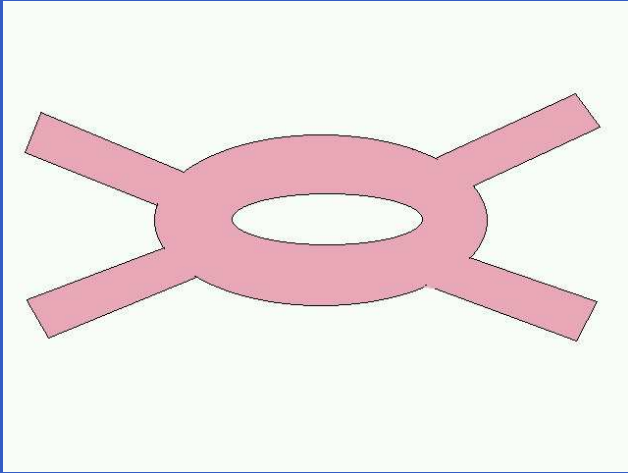
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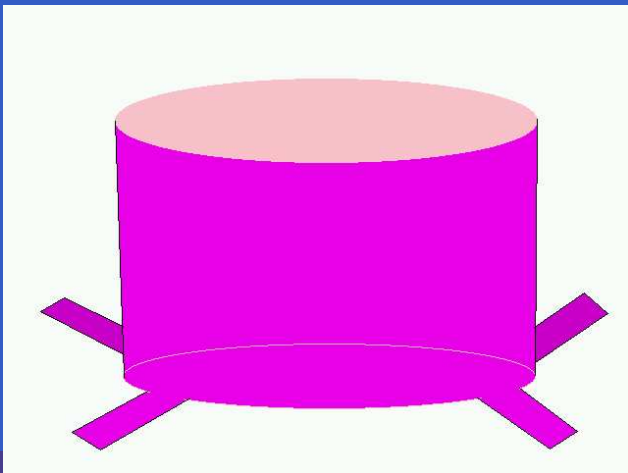
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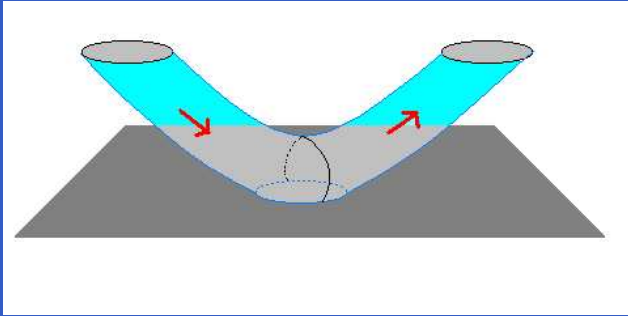
Now there is another way to see this:

**4 open strings  $\rightarrow$  1 closed string**

# Open string interactions ... (ii)

End points of an open strings  
can join to form a **closed string**.

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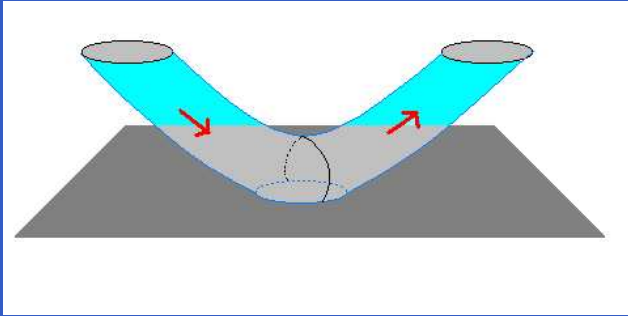


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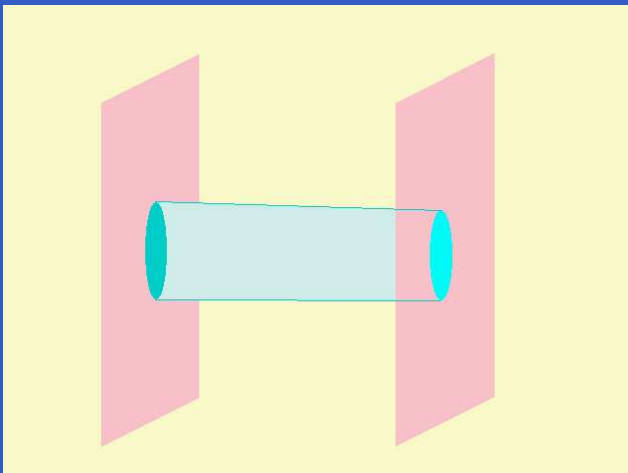
A **closed  $\rightarrow$  open  $\rightarrow$  closed** process involving a D-brane.



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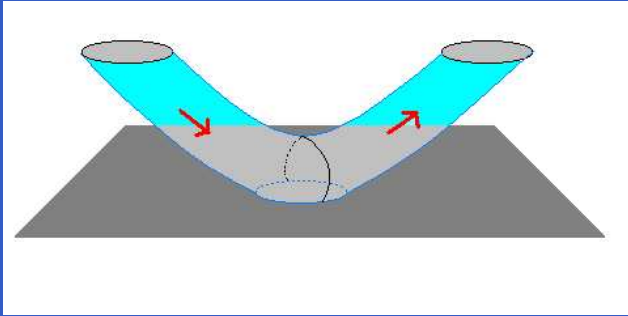


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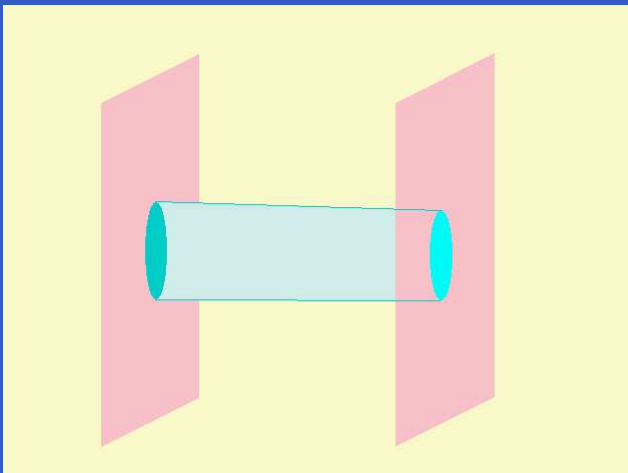
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**D-branes interact gravitationally.**

# Dynamical branes

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How does the D-brane tension compare with the excitation of the string?

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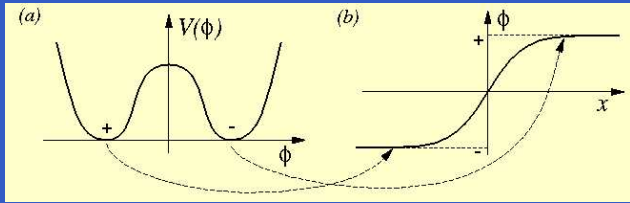
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D-branes are *not* rigid defects in spacetime.

# D-branes & closed strings

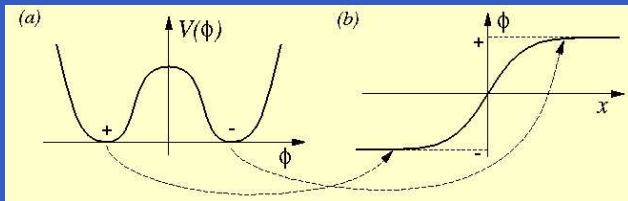
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D-branes are **black hole**-like solutions of *closed string* equations of motion.

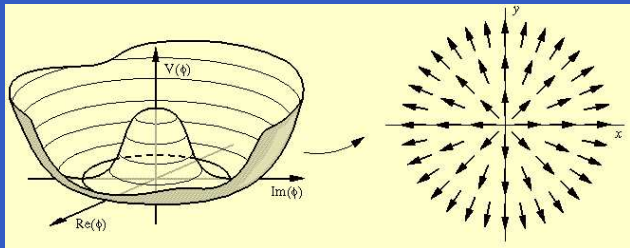
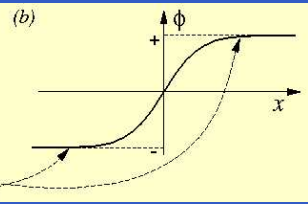
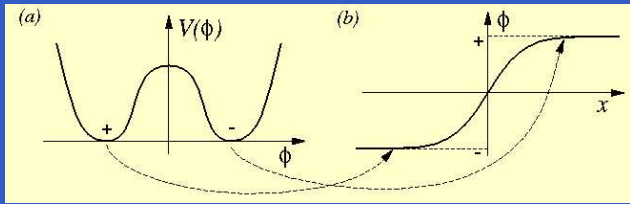
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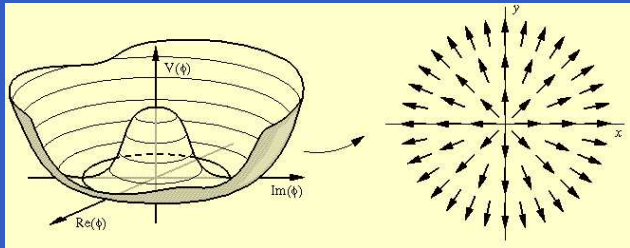
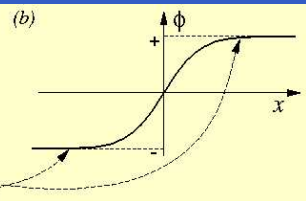
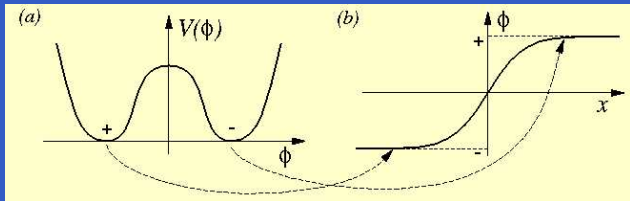
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D-branes are **black hole**-like solutions of *closed string* equations of motion. They carry **mass** and sometimes **charge**.

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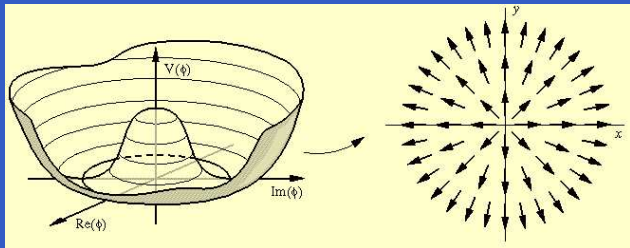
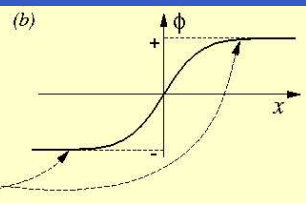
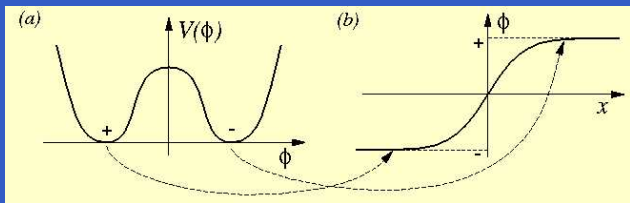


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# D-brane stability

A D-branes carrying *positive* charge is stable.

# D-brane stability

Another one with *negative* charge is also stable  
— an **anti-D-brane**.

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But . . . the two together:

a brane-antibrane pair is *unstable*.

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gravitational and Coulomb force cancel.

# D-brane spectrum

Type IIA

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D0

D2

D4

D6

D8

D1

D3

D5

D7

D9

S

U

S

U

S

U

S

U

S

U

# D-brane spectrum

## Type IIA

D0

D2

D4

D6

D8

D1

D3

D5

D7

D9

S

U

S

U

S

U

S

U

S

U

## Type IIB

# D-brane spectrum

## Type IIA

D0

D2

D4

D6

D8

D1

D3

D5

D7

D9

S

U

S

U

S

U

S

U

S

U

## Type IIB

D0

D2

D4

D6

D8

D1

D3

D5

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U

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# D-branes & T-duality

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$$\text{D-}m\text{-brane} = \begin{cases} \text{D-}(m-1)\text{-brane} \\ \text{D-}(m+1)\text{-brane} \end{cases}$$



# More symmetries

T-duality relates

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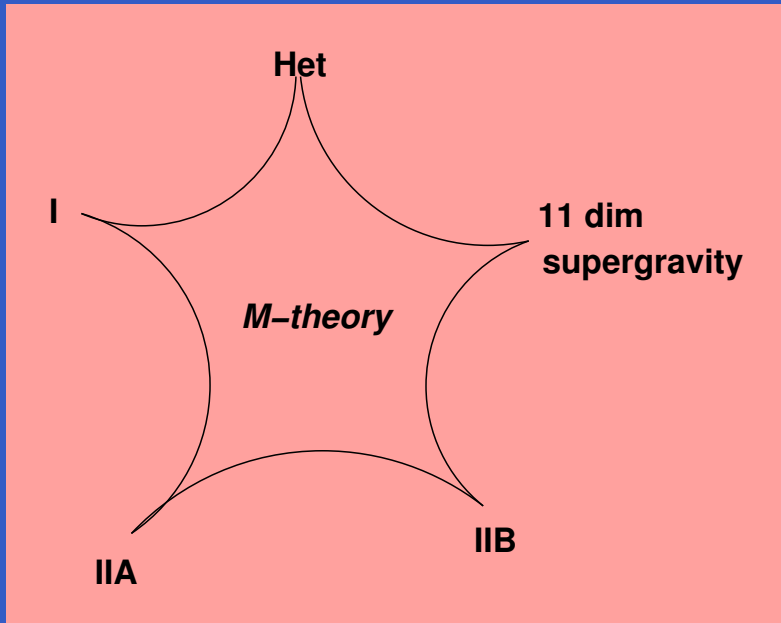
Sometimes Fundamental string  $\xleftrightarrow{g_s \rightarrow 1/g_s}$  D-string

Sometimes Graviton  $\xleftrightarrow{g_s \rightarrow 1/g_s}$  D-brane

# Towards a complete story...

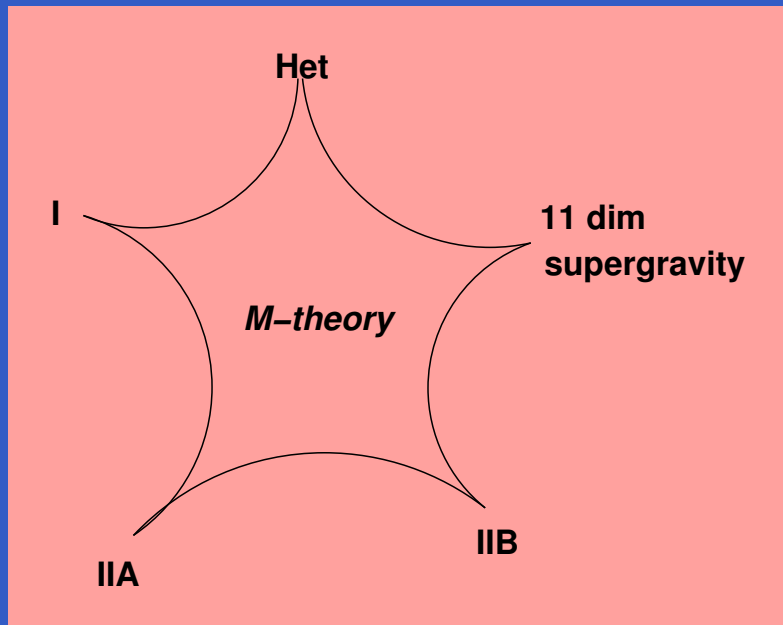
All these symmetries hint  
at a web of relations

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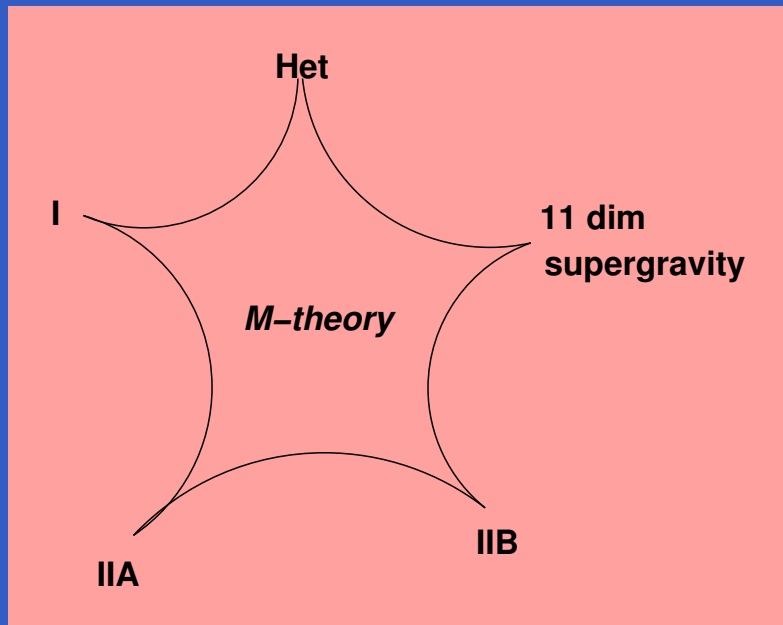


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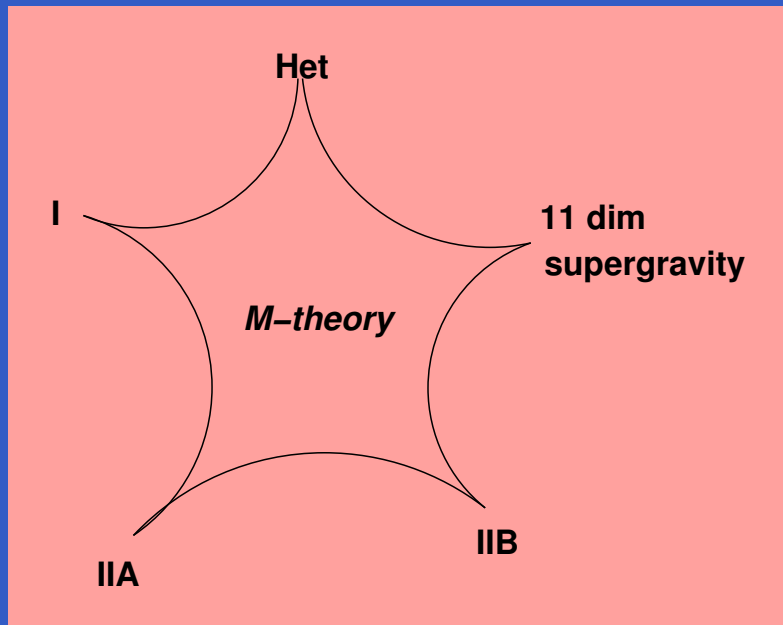
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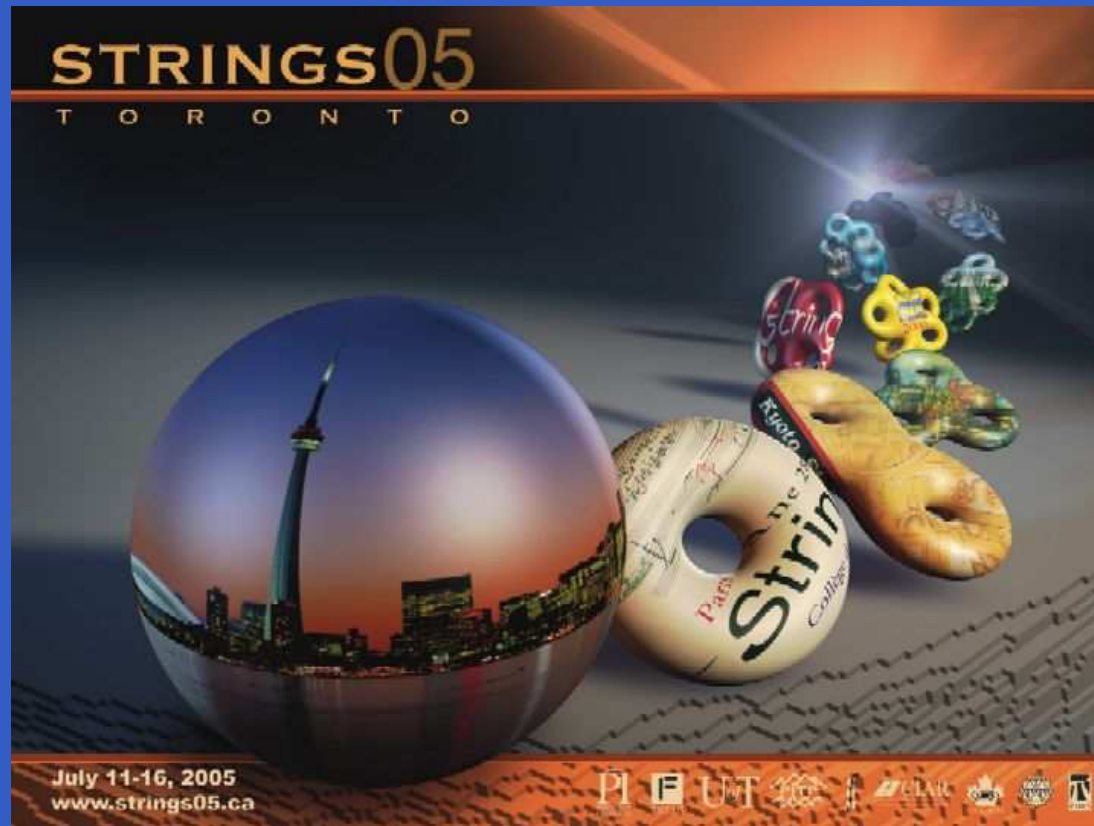
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# Towards a complete story...



एक *M* सद्  
विप्रा बहुधा  
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# Strings and poster art



# Strings and poster art



# Strings and poster art

